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# Quantum algorithms for energy management optimization problems 

## THÈSE

présentée et soutenue publiquement le 16 décembre 2022<br>pour l'obtention du<br>\section*{Doctorat de l'Université de Lorraine<br><br>(Informatique)}<br>par<br>Margarita Veshchezerova

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Mis en page avec la classe thesul.

## Remerciements

Three years on this thesis would be much less enjoyable without several people I would love to thank.

Firstly, I could not have undertaken this journey without my supervisors Emmanuel, Simon, and Marc. I'd like to thank them for the time they gave me, for their experience, and for all the advice that made this work better. I'm grateful for the large variety of topics they introduced to me, it was a pleasure to learn and manipulate many scientific concepts under their supervision.

I would like to extend my sincere thanks to Caroline and Verdran for reviewing my manuscript, and to Bernardetta, Philippe, and Ioan for being part of my defense committee.

Thanks should also go to my former office mate Vladimir and to other members of the (very big) MOCQUA team for the good time that I had in the lab. I gratefully acknowledge the help of my colleagues from EDF - Julien and Joseph.

Obviously, I can't forget to mention my parents and my closest friends - Vika, Sylvain, Charly, and Kiti. Their continuous support helped me to go through this thesis during its most satisfying and most challenging moments.

Finally, words cannot express my gratitude to Dinara and Pierre - for being at my side during these three years, for their encouragement in all moments of doubt, and for making me happy regardless of any difficulties unavoidable in a life of a Ph.D. student.

To my mother and grandmother,

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## Introduction

This thesis, co-funded and co-supervised by EDF R\&D and LORIA, is dedicated to the application of quantum algorithms to combinatorial optimization problems issued from the field of energy management. We consider several optimization problems related to the charge of electric vehicles from the perspective of quantum heuristics such as $Q A O A$ and Quantum Annealing. From a theoretical perspective, this work presents new results derived for the graphical framework for quantum computing called ZX-calculus.

## Smart charging

The field of energy management deals with a lot of complex optimization problems. Finding better solutions for these problems may significantly reduce costs, improve the quality of the service or solve such emergent operational problems as the integration of renewable energy sources.

Recent growth of the number of electric vehicles creates new challenges as well as additional opportunities for electricity management.

On one side, new problems emerge such as charging task attribution, scheduling, cost optimization (among many other). Indeed, even the most powerful charging terminals supplied by Tesld ${ }^{1}$ (with electric power going up to 250 kW ) require around 30 minutes to complete a full charge that assures approximately 400 kilometers of autonomy. We remark that 30 minutes per charge is an extremely optimistic estimation as only $6 \%$ of the charging points in France supply more than 50 kW of electric power. The operational challenges are high as the number of vehicles has dramatically increased in recent times. For instance, according to the French National Association of Electrical Mobility L'Avere-France, in August 2022 there were 69428 public charging points and more than 500000 electric vehicles in France ave, 2022.

The growing number of electric vehicles puts an important pressure on the electrical production and distribution. For instance, in 2019 an average annual electricity consumption was about 4792 kWH per housing Dupret et al., 2021. This statistic was evaluated on a selection of 100 housings one of which was equipped with an electric vehicle. In one year the vehicle in question consumed 2782 kWH , i.e. more than a half of the electrical consumption of an average housing.

On the other side, electric vehicles may positively contribute to the electricity management by improving existing solutions for the storage and the peak load management. Indeed, in a vehicle to grid (V2G) model electric vehicles act not only as consumers, but also as potential suppliers of the energy. For instance, the battery may be used to store the energy during the production peaks (particularly pertinent for renewable sources) and to return it to the grid afterwards. Such scenarios may improve the flexibility of the electrical system, reduce highpeaks, and thus generate significant energy savings, while offering customers with various services involving "earning money" and "saving costs" opportunities.

[^0]The field of optimization problems related to the charging of electric vehicles is known by the name of smart charging.

## Computational hardness of operational problems

A lot of problems of the electricity management in general and smart charging in particular are modeled as NP-hard combinatorial optimization problems. The NP-hardness implies that (at least in the worst case) the search for the optimal solution takes presumably exponential runtime. Therefore, in practice such NP-hard problems are solved only approximately. An approximate solution may be computed by two kinds of algorithms: the ones with a theoreticallyproven approximation ratio and the algorithms (called heuristics) that have demonstrated good empirical results.

In the context of energy management even a minor improvement of the solution's quality may significantly reduce the costs and increase profits, so the active search for better algorithms never ceases.

## Quantum approaches for combinatorial optimization

For computationally difficult problems the rapidly growing field of quantum computing inspires many hopes. Indeed, machines that exploit the principals of quantum mechanic such as entanglement and superposition are provably more powerful than classical Turing machines. Therefore, we expect quantum algorithms to provide significant speedups for some important problems. For the moment, two most impressive results are Grover's algorithm for a search in an unstructured database Grover, 1996 and the algorithm for integer factorization discovered by Shor Shor, 1995. While Grover's algorithm guarantees a quadratic speedup over the best possible classical algorithm, the speedup of Shor's algorithm is exponential over the best known classical counterpart.

The major question is, however, if quantum computing will ever become a reality. Recent progress of multiple hardware teams is very encouraging. For the moment, the biggest universal quantum computer is the 127-qubit machine IBM Eagle Dial, 2022. Google's Sycamore processor has 54 superconducting qubits Arute et al., 2019. The biggest analogous quantum annealer with 5000 qubits is the Advantage system constructed by DWave Systems, 2022. In France the quantum hardware is manufactured by such companies as Pasqa $\|^{2}$ that develops a platform on neutral atoms, Quandeld ${ }^{3}$ that is based on photonic qubits or Alice 83 Bob ${ }_{4}^{4}$ working on superconducting qubits. A comprehensive review of the quantum hardware and software ecosystems can be found in Ezratty, 2021.

On the down side, the experimental progress testifies how difficult it is to robustly scale quantum machines. Therefore, we are not expecting to see a big fault-tolerant computer in the close future. Instead, we can reasonably hope to access in next years Noisy Intermediate Scale devices (called NISQ devises) that will operate around $10^{2}-10^{3}$ qubits. Due to their limited size and high noise level, we don't expect NISQ machines to execute famous Grover's and Shor's algorithms on problems of practically interesting sizes Babbush et al., 2021

There exist, however, quantum algorithms that do not require deep circuits and heavy errorcorrection and, as a consequence, are well adapted for NISQ Preskill, 2018. For combinatorial optimization we are speaking about such heuristics as VQE, QAOA and Quantum Annealing. In

[^1]addition, we believe that the right way to exploit the power of NISQ is to use hybrid quantumclassical procedures that integrate quantum subroutines in some classical algorithm. Following this direction, in this thesis we introduce Quantum-assisted Branch \& Price - an algorithm for huge integer programs that integrate quantum heuristics in the classical Branch \& Price framework.

From the application side, the goal of this thesis is to explore the potential of NISQ-adapted algorithms as well as hybrid procedures on real-world optimization problems. We don't expect to provably solve NP-hard problems to optimality, but rather qualify quantum algorithms as heuristics. We remark that this ambition is limited by the hardware availability as well as its power. Indeed, both current quantum machines and most performant classical emulators such as the $Q L M$ developed by Atos QLM, 2022 can address only relatively simple models of very modest sizes (around 40 variables).

## Qualification of quantum algorithms for smart charging problems

Most conventional quantum heuristics for combinatorial optimization accept the input formulated as quadratic unconstrained binary optimization ( $Q U B O$ ) problems. This means that a special treatment is required to capture in a QUBO model such widely-spread things as constraints and non-binary variables. Proper modelization is indeed a fundamental challenge as many natural formulations of real-world optimization problems exceed the QUBO framework. It turns out that sometimes in order to get a reasonable QUBO formulation one has to significantly reduce of the complexity of the initial model.

We remark that there is an additional difficulty that is usually ignored in the qualification of quantum algorithms. This difficulty is the choice of a relevant classical competitor. Indeed, while an NP-hard optimization problem in general can't be efficiently solved to optimality, there may be very efficient methods that return suboptimal but still satisfying solutions. On the other hand, for many NP-hard problems there exist inapproximability results that bound the theoretically achievable approximation ratio.

In this thesis we consider two real-world problems issued from the domain of electric vehicles. The first problem relates to the scheduling: the goal is to assign a set of charges to different charging stations while minimizing the weighted completion time. In the second problem one has to select a subset of charge demands to satisfy while respecting conflict between demands.

In sections dedicated to numerical results we provide some insights about the experimental protocol for numerical simulations. Our algorithms are tested on realistic data generated from the Belib dataset Bel, 2017 and on a dataset produced by a simulator of charging demands developed by EDF. We compare the performance of our quantum routines to carefully selected classical counterparts.

The work dedicated to the application of quantum heuristics to smart charging problems was previously reported in our paper Dalyac et al., 2021 and is a part of the current thesis (chapters 4 and 5.

## ZX-calculus for variational algorithms

In parallel with experimental evaluation of variational algorithms we considered the ways to enhance the theoretical research in the field. For this purpose we decided to use the ZX calculus - a powerful framework that allows to graphically reason about quantum computing.

In ZX-calculus, originally introduced in Coecke and Duncan, 2011, quantum computations are represented by ZX-diagrams. Each ZX-diagram may be interpreted as a linear map. Conversely, for every linear map $M: \mathbb{C}^{2^{n}} \rightarrow \mathbb{C}^{2^{m}}$ there exists a corresponding ZX-diagram. The diagrammatic representation, however, is not unique, so the ZX-calculus comes with a set of rewrite rules - local graphical transformations that preserves the matrix interpretation of the diagram.

An introduction to the ZX-calculus is provided in the chapter 7.
The ZX-calculus was shown to be extremely helpful in many domains of quantum computing. It was, however, rarely applied in the analysis of variational algorithms such as QAOA. We believe that the reason why variational algorithms are still unexplored with the means of the ZX-calculus is the absence of a convenient way to add and differentiate diagrams. For this reason we introduced in our paper JJeandel et al., 2022] several original techniques that allow diagrammatic addition and the differentiation. We present these techniques in chapter 9 .

## Organization of the manuscript

This document contains three parts.
The first part (I) provides the necessary background in quantum computing and combinatorial optimization. For instance, in chapter (1) that introduces quantum computing we present the concepts of qubit and unitary transformation as well as the circuit notation for quantum programs.

In the following chapter (2) we introduce the complexity notions for combinatorial optimization problems that are relevant for comparisons of quantum and classical methods. In addition, in section 2.4 we provide a brief introduction to the semidefinite programming often used in the design of approximation algorithms.

The second part (II) is dedicated to the application of quantum heuristics to smart charging problems.

It begins with the chapter (3) that introduces quantum algorithms for combinatorial optimization. We present such algorithms as Adiabatic Algorithm (section 3.3), Quantum Annealing (section 3.4), QAQA (section 3.5) and RQAOA (section 3.6). We provide an extensive review of the bibliography on considered algorithms. Based on this bibliography, we compare the expected performance and applicability of the presented methods from the theoretical and practical points of view.

Concerning our contribution, in section 3.5.1 we present our analytical formula (3.37) for the energy expectation in $\mathrm{QAOA}_{1}$ on instances of the weighted MaxCut problem. We also provide our original proof for the bound on approximation ratio achievable by QAOA on the Ising chain model (claim 3.5.2).

The presentation of quantum algorithms for optimization is pursued with examples of their application to the smart scheduling problem (chapter 4) and the charge task selection problem (chapter 5). For both considered applications we present the original usecase formulations, the complexity analysis, the modelization process as well as numerical results that are compared to the performance of classical methods. For the modelization part we introduce in the section 4.3.1 an original binary encoding technique allowing the integration of bounded integer variables in the QUBO framework. Finally, we compare our quantum algorithms to carefully chosen classical counterparts in order to place our results in the context of the qualification of quantum approaches on real-world applications.

In the following chapter (6) we present an original hybrid quantum-classical algorithm for
huge integer programs that was developed during this thesis. The method is based on the column generation approach presented in section 6.2.4. We baptize our hybrid method under the name of Quantum-assisted Branch ES Price. Quantum-assisted Branch \& Price can be applied to integer programs with a large number of variables as ones obtained from Dantzig-Wolfe decomposition (see section 6.2.1). In section (6.4), we illustrate our approach on an example of graph coloring problem. We numerically evaluate the performance of our method on a version of the charge task selection problem.

The last part (III) presents the application of the ZX-calculus (introduced in chapter 7) to variational algorithms. In chapter (8), we show how QAOA routines can be represented with ZX-diagrams. We illustrate the process of graphical rewriting by showing how to derive an analytical expression for the loss function optimized by the classical loop of QAOA. This proof is presented in the current thesis for the first time.

In the final chapter (9) we describe our original techniques for addition (section 9.2) and differentiation (section 9.3) of ZX-diagrams. In section (9.4) we demonstrate how these techniques can be used to derive a ZX-diagram for an Ising Hamiltonians.

In order to make the work as comprehensible as possible, we include the appendix A that introduces the basic concepts of linear programming.

## Part I

## Background

## Chapter 1

## Quantum computing

The idea to use quantum effects to enhance the computations of properties of physical systems (and computation in general) was independently suggested by Richard Feynman [Feynman, 1986], Yuri Manin [Manin, 1980] and Paul Benioff [Benioff, 1982]. This idea is particularly relevant in the context of the simulation of quantum systems. Indeed, according to the laws of quantum mechanics such system behaves differently from what could be expected in classical physics. In particular, a quantum system can be "in many states at once" - a phenomenon known as quantum superposition. An explicit enumeration of a superposition state requires an exponential amount of complex coefficients. Therefore, the simulation rapidly becomes intractable when the size of the quantum system grows. In other words, the nature is quantum and classical computing models are not powerful enough to efficiently compute properties of quantum systems.

On the practical side, the miniaturization process that made the classical computer so powerful, can't continue at the previously-observed exponential rate reclaimed by the famous empirical Moore law Wikipedia, 2022b. The limitation is due, inter alia, to quantum effects that have a significant impact on the dynamics of small systems. In opposition, quantum computers are physical systems that, rather than considering quantum effects as a resource of undesirable noise, use them as an opportunity to compute differently. Somewhat tautologically, quantum computing can be defined as the domain of computer science that explores the computational power of such quantum computers. The quantum computational model was formalized in a definition of universal quantum Turing machines in Deutsch, 1985.

Early algorithms for quantum computers such as Deutsch-Jozsa Deutsch and Jozsa, 1992 and Simon's algorithms Simon, 1997] were mostly of theoretical interest. It was the procedure for an efficient integer factorization suggested by Shor [Shor, 1995] that attracted a lot of attention to quantum computing.

This chapter provides a brief introduction to the basic concepts of quantum computing. We consider, in particular, the digital model of quantum computing that is realized on universal machines. Analogue quantum machines supporting adiabatic quantum computing are presented in chapter 3 .

A seminal reference for an introductory course in the domain of quantum computing is Nielsen and Chuang, 2011. Alternatively, the lecture notes de Wolf, 2019 provides a comprehensive and accurate presentation of the field of quantum computing and related domains.

### 1.1 Dirac notations

In a nutshell, a quantum algorithm is a transformation of a vector from the space $\mathbb{C}^{2^{n}}$ by unitary matrices from $\mathbb{C}^{2^{n} \times 2^{n}}$ followed by a probabilistic sampling (we denote by $\mathbb{C}^{a \times b}$ a complex matrix with $a$ columns and $b$ rows). In order to reason about vectors and matrices, most modern works on quantum mechanics and quantum computing use the so-called Dirac notation.

In this notation we call by "ket" a column vector:

$$
|\psi\rangle=\left(\begin{array}{c}
\psi_{0}  \tag{1.1}\\
\vdots \\
\psi_{n}
\end{array}\right)
$$

and by "bra" a row vector:

$$
\begin{equation*}
\langle\psi|=\left(\psi_{0}^{*}, \ldots, \psi_{n}^{*}\right) \tag{1.2}
\end{equation*}
$$

with complex numbers as elements. We can get a bra-vector from the ket-vector by taking its conjugate transpose.

We denote by dagger $v^{\dagger}$ (alternatively called adjoint) the transposition and complex conjugation, i.e. for $v=\left(\begin{array}{c}v_{0} \\ \vdots \\ v_{n}\end{array}\right)$ we have $v^{\dagger}=\left(v_{1}^{*}, \ldots, v_{n}^{*}\right)$. Certainly, for the ket-state $|\psi\rangle$ the adjoint is the bra-state $\langle\psi|$. For a linear map $M: \mathbb{C}^{n \times m}$ it adjoint $M^{\dagger}: \mathbb{C}^{m \times n}$ is similarly defined as a combination of the transposition and the complex conjugation.

Using the "bra" and "ket" notations we can express the norm of a column vector $|\psi\rangle$ in the Hilbert space $\mathbb{C}^{n}$ as a bracket $\langle\psi \mid \psi\rangle$. We recall that in the Hilbert space the inner product of vectors $\langle\phi|$ and $|\psi\rangle$ is precisely $\sum_{i=1}^{n} \phi_{i}^{*} \psi_{i}$. We denote the inner product with the bracket $\langle\phi \mid \psi\rangle$.

Following the widely-adopted conventions, we denote by $|0\rangle$ and $|1\rangle$ the standard basis in a two-dimensional complex space $\mathbb{C}^{2}$ :

$$
\begin{equation*}
|0\rangle=\binom{1}{0}, \quad|1\rangle=\binom{0}{1} \tag{1.3}
\end{equation*}
$$

The basis in the space $\mathbb{C}^{2^{n}}$ of dimension $2^{n}$ is usually enumerated by bitstrings $x \in\{0,1\}^{n}$.

### 1.2 Qubits

In classical computing, the information is carried by bits. A bit is a classical system that is in one of two mutually exclusive states denoted by 0 and 1 . Crucially, at each moment a classical system is strictly in one state. We remark that in probabilistic algorithms some bits $b \in\{0,1\}$ can be sampled from a probability distribution $p \in \mathbb{R}^{n}$ given by a vector of real values with the 1 -norm $\sum_{i=1}^{n} p_{i}=1$ while in quantum computing the information if carried by a vector $|\psi\rangle$ of complex numbers with the Hermitian norm $\langle\psi \mid \psi\rangle=1$. In addition, once a classical bit is assigned a random value it remains in the same (potentially unknown) state unless it is explicitly reassigned. In quantum computing, a simple observation of the system actually modifies the state of the system.

A classical algorithm can interact with the memory in two ways. Firstly, we can assign a value $v \in\{0,1\}$ to the bit $b:=v$. An inverse procedure outputting the value of the bit $o:=b$ is called reading.

In quantum computing, the information is contained in a state of some quantum system. We consider systems that have two distinct elementary states $|0\rangle$ and $|1\rangle$. The simplest example of a such two-dimensional quantum system is the spin of a particle that can be turned up or down: $|0\rangle=|\uparrow\rangle$ and $|1\rangle=|\downarrow\rangle$. A two-level quantum system is called qubit.

In mathematical terms, state of a qubit $\left|\psi_{0}\right\rangle \in \mathbb{C}^{2}$ is a vector of two complex numbers $a$ and $b$ often called amplitudes:

$$
\begin{align*}
\left|\psi_{0}\right\rangle= & a|0\rangle+b|1\rangle  \tag{1.4}\\
& |a|^{2}+|b|^{2}=1 \tag{1.5}
\end{align*}
$$

We observe that the only restriction on amplitudes is the condition stating that the norm of the vector should be equal to one. Therefore, a state of a qubit can be any vector on the unit sphere $\mathbb{S}^{2}$ in the two-dimensional complex space. If both amplitudes $a$ and $b$ are non-zero we say that the qubit is in a superposition of basis states.

### 1.2.1 Preparation and measurement

Similarly to reading and writing for classical bits, there exist two ways to interact with a qubit called preparation and measurement.

Preparation in a trivial state is analogous to the writing for bits:

$$
\begin{equation*}
|\psi\rangle=|0\rangle \tag{1.6}
\end{equation*}
$$

The reading part, called measurement, is slightly less intuitive. The complexity is due to the fact that we are classical, so we can access only classical information about the quantum system. More precisely, the measurement of a qubit returns a classical output $o \in\{0,1\}$ which is a binary number. In such settings, the quantum state $|\psi\rangle=a|0\rangle+b|1\rangle$ defines the probability to measure a specific value:

$$
\begin{align*}
& P(o=0)=|a|^{2}  \tag{1.7}\\
& P(o=1)=|b|^{2} \tag{1.8}
\end{align*}
$$

Crucially, two subsequent measurements without modifications in-between should return exactly the same output $|o\rangle$. Therefore, the measurement "collapses" the superposition state $|\psi\rangle=$ $a|0\rangle+b|1\rangle$ to the basis state $|o\rangle$ corresponding to the measurement result. We highlight that after the measurement the initial superposition is irreversibly lost. Moreover, the measurement returns a binary value and the coefficients $a$ and $b$ are never directly accessible to the classical spectator.

We remark that the probability distribution is defined by the squared norms of amplitudes. This fact explains why the state should have a unit norm, i.e. why $|a|^{2}+|b|^{2}=1$. We also notice that the multiplication by a global complex phase $|\psi\rangle \rightarrow e^{i \alpha}|\psi\rangle$ has no influence on the distribution, so we can safely ignore it. Therefore, the qubit state has precisely two degrees of freedom: one for each angle $\beta$ and $\gamma$ in the expression $|\psi\rangle=\cos \beta|0\rangle+e^{i \gamma} \sin \beta|1\rangle$. As the state has two degrees of freedom it can be visualized on a Bloch sphere shown in figure 1.1

### 1.2.2 N -qubit systems

## Tensor product

For each pair of vectors $|v\rangle$ and $|u\rangle$ we introduce the shortcut notation:

$$
\begin{equation*}
|v\rangle|u\rangle=|v\rangle \otimes|u\rangle \tag{1.9}
\end{equation*}
$$



Figure 1.1: The qubit state $|\psi\rangle=\cos \beta|0\rangle+e^{i \gamma} \sin \beta|1\rangle$ on the Bloch sphere. The figure is plotted with the python package QuTIP [https://qutip.org]
where $\otimes$ is the tensor product of vectors also sometimes called Kronecker product. We recall that for vectors $|v\rangle=\left(\begin{array}{c}v_{1} \\ \vdots \\ v_{n}\end{array}\right) \in \mathbb{C}^{n}$ and $|u\rangle=\left(\begin{array}{c}u_{1} \\ \vdots \\ u_{m}\end{array}\right) \in \mathbb{C}^{m}$ the tensor product $|v\rangle \otimes|u\rangle$ is a vector in $\mathbb{C}^{n m}$ of the form:

$$
|v\rangle \otimes|u\rangle=\left(\begin{array}{c} 
 \tag{1.10}\\
v_{1}\left(\begin{array}{c}
u_{1} \\
\vdots \\
u_{m}
\end{array}\right) \\
\vdots \\
\\
v_{n}\left(\begin{array}{c}
u_{1} \\
\vdots \\
u_{m}
\end{array}\right)
\end{array}\right)
$$

For the matrices $A \in \mathbb{C}^{n \times m}$ and $B \in \mathbb{C}^{k \times l}$ the tensor product is similarly defined as

$$
A \otimes B=\left(\begin{array}{ccc}
a_{1}^{1} B & \ldots & a_{1}^{n} B  \tag{1.11}\\
& \ddots & \\
a_{m}^{1} B & \ldots & a_{m}^{n} B
\end{array}\right)
$$

## Entanglement

Trivially, a state of $n$ classical bits $\boldsymbol{b} \in\{0,1\}^{n}$ is in a one-to-one correspondence with $n$ states of $n$ bits $\left[b_{0}, \ldots, b_{n-1}\right]$. A peculiar fact about quantum computing is that in general a state of $n$ qubits can't be decomposed to $n$ independent states of qubits $\left|\psi_{0}\right\rangle \ldots\left|\psi_{n-1}\right\rangle$.

Indeed, a state of $n$ qubit system $|\boldsymbol{\psi}\rangle$ lays on the unit sphere in the Hilbert space $\mathbb{C}^{2^{n}}$ and not in the direct-sum of $n$ separated 2-dimensional Hilbert spaces $\mathbb{C}^{2} \oplus \cdots \oplus \mathbb{C}^{2}$. This implies in particular that some states of an n-qubit quantum register can't be decomposed on a tensor product of $n$ individual one-qubit states. Such phenomenon is called entanglement and it is believed to be a fundamental source of "quantumness" in quantum computing.

As in the one-qubit case, the state $|\boldsymbol{\psi}\rangle$ defines a probability distribution. Formally, the state in $\mathbb{C}^{2^{n}}$ can be written as a linear combination of basis states $|x\rangle, x \in\{0,1\}^{n}$ :

$$
\begin{gather*}
|\boldsymbol{\psi}\rangle=\sum_{x \in\{0,1\}^{n}} \alpha_{x}|x\rangle  \tag{1.12}\\
\sum_{x \in\{0,1\}^{n}}\left|\alpha_{x}\right|^{2}=1 \tag{1.13}
\end{gather*}
$$

where $\left|\alpha_{x}\right|^{2}$ corresponds to the probability to measure the output $|x\rangle$. As the basis states $|x\rangle$ are orthogonal, the coefficients $\alpha_{x}$ correspond to the inner product $\alpha_{x}=\langle x \mid \boldsymbol{\psi}\rangle$.

The equality 1.13 can be alternatively written as:

$$
\begin{equation*}
\langle\psi \mid \psi\rangle=1 \tag{1.14}
\end{equation*}
$$

The simplest example of an entangled state is the Bell state:

$$
\begin{equation*}
|\phi\rangle=\frac{|00\rangle+|11\rangle}{\sqrt{2}} \tag{1.15}
\end{equation*}
$$

We can verify that there don't exist $\left|\phi_{0}\right\rangle=\binom{a}{b}$ and $\left|\phi_{1}\right\rangle=\binom{c}{d}$ such that $|\phi\rangle=\left|\phi_{0}\right\rangle\left|\phi_{1}\right\rangle$. Indeed, $\left|\phi_{0}\right\rangle\left|\phi_{1}\right\rangle=a c|00\rangle+a d|01\rangle+b c|10\rangle+b d|11\rangle$ and the requirement $a d=b c \equiv 0$ contradicts to the fact that $a c=b d \equiv \frac{1}{\sqrt{2}}$.

### 1.3 Quantum evolutions

In quantum algorithm, qubit states are transformed by linear maps:

$$
\begin{equation*}
|\psi\rangle \xrightarrow{U}\left|\psi^{\prime}\right\rangle=U|\psi\rangle \tag{1.16}
\end{equation*}
$$

In order to obtain valid resulting states the transformation should map the unit sphere $\mathbb{S}^{2 n}$ to itself, i.e. be norm-preserving. This requirement is satisfied by so-called unitary matrices.

Unitary matrices can be equivalently defined as square matrices for which we have:

$$
\begin{equation*}
U^{\dagger} U=I \tag{1.17}
\end{equation*}
$$

In other words, the inverse of $U$ is precisely the adjoint $U^{\dagger}$ of $U$. Moreover, the inverse of a unitary matrix is also unitary i.e. quantum computing allows only invertible transformations. The only exception to this rule is the measurement process that irreversibly collapses the superposition $|\psi\rangle=\sum a_{x}|x\rangle$ to the measurement output $|\hat{x}\rangle$.

A transformation $U$ is fully defined by the matrix of the corresponding linear map:

$$
\begin{equation*}
M_{U} \in \mathbb{C}^{2^{n} \times 2^{n}} \tag{1.18}
\end{equation*}
$$

An explicit matrix representation for unitary transformations is extremely inefficient and resource-consuming. It is also impractical for programming, as we can't use the matrix to specify a set of instructions that have to be executed on the quantum computer. In classical computing of boolean functions, the analog of the matrix representation are the truth tables. We remark that truth tables and matrices are never used for large computations as there exist much more convenient representations that use circuits.

## Circuit representation

A classical circuit is a graphical scheme that has i) input wires i) output wires and iii) elementary blocks conventionally called gates. Crucially, a small set made out of elementary AND-gate, ORgate, and NOT-gate is powerful enough to express any boolean function as a circuit made out of them.

For example, we can represent the function $f(\boldsymbol{b})=b_{2} \vee\left(b_{0} \wedge b_{1}\right)$ with the circuit shown in figure 1.2


Figure 1.2: Classical circuit for $f(\boldsymbol{b})=b_{2} \vee\left(b_{0} \wedge b_{1}\right)$
The idea to use circuits was borrowed by quantum computer science. A wire in a quantum circuit corresponds to a qubit. A circuit for the unitary evolution $U$ is denoted as:


In a quantum circuit model, unitary transformations (just as boolean functions in classical circuits) are written as compositions of some elementary gates. The only fundamental difference with the classical case is that for purely quantum gates the number of inputs is always equal to the number of outputs. This condition directly follows from the reversibility of unitary quantum transformations.

A typical quantum circuit looks like the following:


In a circuit representation, we often omit the measurements as we usually assume that they are performed at the end of the circuit. If it is not the case (for instance in protocols that use ancilla qubits) we explicitly integrate the non-trivial measurement in the circuit.

### 1.3.1 Elementary gates

We briefly present some most commonly used gates.

## Hadamard gate

So-called Hadamard gate is probably the most famous one-qubit transformation.

$$
H-=\frac{1}{\sqrt{2}}\left(\begin{array}{cc}
1 & 1  \tag{1.21}\\
1 & -1
\end{array}\right)
$$

The Hadamard gate maps the computational basis states $|0\rangle$ and $|1\rangle$ to a pair of orthogonal states $|+\rangle$ and $|-\rangle$ :

$$
\begin{align*}
& \mathrm{H}|0\rangle=\frac{|0\rangle+|1\rangle}{\sqrt{2}}=|+\rangle  \tag{1.22}\\
& \mathrm{H}|1\rangle=\frac{|0\rangle-|1\rangle}{\sqrt{2}}=|-\rangle \tag{1.23}
\end{align*}
$$

The pair $\{|+\rangle,|-\rangle\}$ with $|+\rangle=\frac{1}{\sqrt{2}}\binom{1}{1}$ and $|-\rangle=\frac{1}{\sqrt{2}}\binom{1}{-1}$ is usually called the Hadamard basis. Alternatively, in some cases the bases $\{|0\rangle,|1\rangle\}$ and $\{|+\rangle,|-\rangle\}$ are respectively called $Z$-basis and $X$-basis.

## Pauli gates

For the reader with physical background Pauli gates should remind of Pauli matrices used to represent spin interactions in quantum mechanics.

The $X$-gate is equivalent to the classical NOT gate. For the states $|0\rangle$ and $|1\rangle$ it performs the bit flip: $|b\rangle \rightarrow|\bar{b}\rangle$. The matrix for $X$ is:

$$
-X-=\left(\begin{array}{ll}
0 & 1  \tag{1.24}\\
1 & 0
\end{array}\right)
$$

The $Z$-gate performs the so-called phase fip: $|b\rangle \rightarrow(-1)^{b}|b\rangle$. The matrix for the $Z$-gate is

$$
-Z-\left(\begin{array}{cc}
1 & 0  \tag{1.25}\\
0 & -1
\end{array}\right)
$$

We remark that vectors from computational basis $\{|0\rangle,|1\rangle\}$ are eigenvectors for the $Z$-gate while the vectors $|+\rangle$ and $|-\rangle$ are eigenvectors for the $X$-gate. This is by the way the reason why these bases are sometimes called $Z$-basis and $X$-basis.

Matrices $X$ and $Z$ have real-values eigenvectors. This is not the case for the last Pauli gate called $Y$-gate:

$$
-Y=\left(\begin{array}{cc}
0 & -i  \tag{1.26}\\
i & 0
\end{array}\right)
$$

Pauli gates $X, Y$ and $Z$ are Hermitian and involutary $\left(P^{2}=I\right.$ for $\left.P=X, Y, Z\right)$. Moreover, together with the identity $I$ they form a basis in the vector space of $2 \times 2$ Hermitian matrices.

## Rotation gates

A rotation gate is actually not a single matrix but rather a family of gates parameterized by a real value $\gamma$ called rotation angle.

We denote by $R_{X}(\gamma)$ the rotation around $X$-axis:

$$
-R_{X}(\gamma)=e^{i \gamma X / 2}=\left(\begin{array}{cc}
\cos \left(\frac{\gamma}{2}\right) & i \sin \left(\frac{\gamma}{2}\right)  \tag{1.27}\\
i \sin \left(\frac{\gamma}{2}\right) & \cos \left(\frac{\gamma}{2}\right)
\end{array}\right)
$$

The rotation around $Z$-axis is:

$$
R_{Z}(\gamma)=e^{i \gamma Z / 2}=e^{i \frac{\gamma}{2}}\left(\begin{array}{cc}
1 & 0  \tag{1.28}\\
0 & e^{-i \gamma}
\end{array}\right)
$$

We recall that the matrix exponential $e^{i \gamma A}$ is the power series:

$$
\begin{equation*}
e^{i \gamma A}=\sum_{k=0}^{\infty} \frac{(i \gamma)^{k} A^{k}}{k!} \tag{1.29}
\end{equation*}
$$

For involutary matrices $A^{k}=\left\{\begin{array}{ll}I, & \text { if } k=0 \bmod 2 \\ A, & \text { otherwise }\end{array}\right.$, so the power series 1.29 can be written as $e^{i \gamma A}=\cos \gamma I+i \sin \gamma A$.

We remark that for the rotation angle $\pi$ the gates $R_{X}(\pi)$ and $R_{Z}(\pi)$ are precisely the Pauli gates $X$ and $Z$ up to a global phase.

Another important rotation gate is the so-called T-gate $T=R_{Z}(\pi / 4)$. Together with Pauli gates, Hadamard gate, and a single two-qubit transformation the T-gate form an approximately universal set of gates for pure-qubit quantum computing.

## CNOT gate

Interesting quantum effects such as entanglement occur when qubits interact with each other. A basic multiqubit gate is the two-qubit controlled NOT gate called CNOT. Intuitively, the CNOT gate does nothing with the state if the first qubit is in the $|0\rangle$-state while it flips the value of the second qubit if the first one is in $|1\rangle$ state:

$$
\begin{align*}
& |0\rangle|b\rangle \xrightarrow{\mathrm{CNOT}}|0\rangle|b\rangle  \tag{1.30}\\
& |1\rangle|b\rangle \xrightarrow{\mathrm{CNOT}}|0\rangle|\bar{b}\rangle \tag{1.31}
\end{align*}
$$

The matrix representation for CNOT is:

$$
\underline{\Xi}=\left(\begin{array}{llll}
1 & 0 & 0 & 0  \tag{1.32}\\
0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 \\
0 & 0 & 1 & 0
\end{array}\right)
$$

In the circuit notation it is represented by a wire that connects two different dots - one corresponding to the control and the other to the target qubit.

### 1.3.2 Compositions

Two elementary gates can be applied in parallel. The underlying transformation of parallel composition corresponds to the tensor product of individual unitaries:

$$
\begin{array}{ll}
\mathscr{A} & =A \otimes B  \tag{1.33}\\
\mathcal{B} &
\end{array}
$$

A sequential composition of gates implements the unitary equal to the multiplication of individual blocks:


We remark that in circuits the time evolves from left to right.
The gate set H , CNOT, and $R_{Z}(\gamma)$ is universal for quantum computation. Strictly speaking, on can write an arbitrary unitary evolution with a circuit containing only these gates. We remark that for each matrix the decomposition is usually not unique, for example two circuits:

and $\qquad$
represent the same matrix $X=\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right)$.

### 1.4 Quantum program

In nutshell, a quantum program prepares a system in a trivial state $\left|\psi_{0}\right\rangle=|0 \ldots 0\rangle$ to which it applies a unitary transformation $U$ given by a circuit. In the end the state $U\left|\psi_{0}\right\rangle$ is measured and a classical bitstring is recovered:


Just before the measurement the quantum state is:

$$
\begin{equation*}
|\psi\rangle=U|0 \ldots 0\rangle=\sum_{x \in\{0,1\}^{n}} \alpha_{x}|x\rangle \tag{1.37}
\end{equation*}
$$

Measurement returns a sample from probability distribution:

$$
\begin{equation*}
\mathbb{P}(x)=\left|\alpha_{x}\right|^{2}, \quad x \in\{0,1\}^{n} \tag{1.38}
\end{equation*}
$$

More complicated distributions can be prepared using auxiliary ancilla qubits and postselection.

Indeed, it is possible to consider additional ancilla inputs initialized in states $|0\rangle$. Ancillas allow us to perform the linear quantum evolution in a bigger space. One can also use a postselection on the measurement output, for example, keep only outputs where the first qubit takes value the 1 . We remark that the post-selection, while possible in computing, is not a physical process.

Ancillas and postselection lead to non-unitary transformations. Nevertheless, there exist blocks permitting to integrate them inside the circuit notations:


From a very general point of view programming for a quantum computer consists in combining small bricks to prepare a probability distribution that is relevant for the problem in question. Certainly, it is a stiff exercise to design algorithms in such a framework. The challenge gets even more difficult when in addition we aim to find quantum algorithms that have a provable advantage over classical counterparts.

In part due to the difficulty of designing quantum algorithms the field of quantum computing transferred a big switch of paradigm moving to variational algorithms. In a nutshell, variational algorithms are not fixed programs solving a problem but rather trainable models that are optimized to find good solutions. We come back to the variational algorithm in chapter 3.

## Chapter 2

## Combinatorial Optimization

Combinatorial optimization is the field of mathematics that considers problems over discrete variables. Models with discrete variables are ubiquitous in many real-life applications such as scheduling, resource allocation, logistics, and others. From the formal side, combinatorial optimization incorporates such important models as constraint satisfaction problems, problems on graphs and other discrete structures as well as general optimization problems over integer variables.

Many problems of combinatorial optimization are computationally difficult. In the first part of the current chapter we formally define what is the computational difficulty of an optimization problem. In addition, we briefly introduce another important concept of approximation complexity. Most of our definitions are formulated as in the book Ausiello et al., 1999.

The second part of the chapter is dedicated to semidefinite programs (Section 2.4). Strictly speaking, semidefinite programs are optimization problems over continuous variables. However, they are extremely important in the design of approximation algorithms for combinatorial problems.

### 2.1 Decision, search and optimization problems

Optimization problems are closely related to decision problems and search problems. The relation to decision problems is particularly meaningful for the complexity analysis of the optimization problems. Indeed, $N P$-completeness which is the fundamental notion in the complexity theory, is defined in the framework of decision problems.

An optimization problem as well as decision and search problems are provided with a set of instances $I$ and solutions $S$. A predicate $\pi: I \times S \rightarrow\{0,1\}$ indicates if $s \in S$ is a solution for an instance $i \in I$.

- Decision problems. The solution set contains only two elements $S=\{0,1\}$. For every instance $i \in I$ we have $\pi(i, 0)$ or $\pi(i, 1)$. In other words, the predicate $\pi$ splits the set $I=I_{p} \sqcup I_{n}$ on positive and negative instances. The problem is to decide for a given instance $i$ if $i \in I_{p}$ or $i \in I_{n}$.
- Search problems. To each instance $i$ corresponds a set of feasible solutions $F(x)=\{y \in$ $S \mid \pi(x, y)=1\}$. The goal is to find for a given input instance a feasible solution $y^{*}$ such that $\pi\left(x, y^{*}\right)=1$.
- Optimization problems. For each instance feasible solutions are assigned to a score. The score is typically computed with so-called measure function $m: i \times F(i) \rightarrow m_{i} \in \mathbb{Z}$. We are looking for a feasible solution with the highest (or lowest) score value.

Formally, an optimization problem is defined by a tuple ( $I, F, m$ ) and a goal which can be maximization or minimization. The set $I$ contains instances of the problem. The map $F$ associates to each valid instance $i$ a set of feasible solutions $F(i)$. The function $m: I \times F(I) \rightarrow \mathbb{Z}$ is called measure function. The measure function associates to each pair $(i, y)$ where $i \in I$ is an instance of the problem and $y \in F(i)$ is a feasible solution for the instance $i$ a positive integer value. If the goal is to maximize the score, the measure function can be referred to as profit function. For minimization problems terms cost function and loss function are employed in the literature.

An optimal solution $v^{*}$ to a maximization problem for a given instance $i \in I$ is a feasible solution $y \in F(i)$ of highest measure:

$$
\begin{equation*}
y^{*} \in F(i) \text { s.t. } \forall y \in F(i): m(i, y) \leq m\left(i, y^{*}\right) \tag{2.1}
\end{equation*}
$$

The value $m^{*}(i)=m\left(i, y^{*}\right)$ is called the optimum value of the problem. For each instance the optimum value $m^{*}(i)$ is unique. Certainly, the set of optimal solutions $F^{*}(i)=\{y \in F(i) \mid m(i, y)=$ $\left.m^{*}(x)\right\}$ may contain several elements.

We illustrate the definition of an example of graph coloring problem. In graph coloring the set of instances $I$ is the set of all graphs $G=(V, E)$. For each graph feasible solutions are proper colorings, i.e. functions $c: V \rightarrow \mathbb{N}$ such that connected vertices $(u, v) \in E$ are assigned to different colors: $c(u) \neq c(v)$. The measure or cost of coloring is the number of different colors, i.e. $m(i, c)=\#\{c \in \mathbb{N} \mid \exists u \in V: c(u)=c\}$.

Each optimization problem $\mathcal{P}$ straightforwardly generates a decision problem $\mathcal{P}_{D}$. For instance, in the graph coloring example we may ask if there exists a coloring $c$ that use less than $k$ colors for a fixed number $k \in \mathbb{N}$. We remark that in the search version we would ask for such a coloring. Importantly, an algorithm solving the optimization problem directly solves the decision problem. Indeed, given $m^{*}(i)$ it is simple to verify if $m^{*}(i) \leq k$.

Most of optimization problems of practical relevance belong to so-called class NPO:
Definition 2.1.1 An optimization problem $(I, F, m)$ is in class NPO if it satisfied the following criteria:

- the set I is recognizable in polynomial time
- feasible solutions have polynomially-bounded size, i.e. there exists a polynomial $p: \mathbb{N} \rightarrow \mathbb{R}$ such that for each $y \in F(i)$ we have $|y| \leq p(|i|)$. Moreover, for all $y$ such that $|y| \leq p(|i|)$ we can decide in polynomial time if the solution $y$ is feasible for the instance $i$.
- the measure function $m$ is computable in polynomial time


### 2.2 Complexity of optimization problems

The complexity analysis of an optimization problem essentially indicates the scaling of the amount of resources required to solve the problem on the size of the input instance. The size of the instance is usually defined as the number of bits that are necessary to encode it. We remark
that different encodings lead to different size evaluations. However, all natural encoding schemes that do not introduce unnecessary redundancy are polynomially related Ausiello et al., 1999.

An algorithm solving an optimization problem defines an upper bound on its runtime complexity. We notice that an algorithm may have very different behavior on instances of the same size. For example, the greedy approach easily finds an optimal coloring for any bipartite graph $G_{b}=\left(V_{b}, E_{b}\right)$ while it may have very poor performance on another graph with the same number of nodes and edges. Therefore, the computational complexity considers the runtime scaling in the worst-case.

For decision problem we say that it is solvable in $O(g(n))$ if there exists an algorithm $\mathcal{A}$ such that for any instance $i$ of size $|i| \leq n$ the algorithm $\mathcal{A}$ decides if the instance is positive in the runtime bounded by $g(n)$.

A lower bound, if there is one, indicates how much time is necessary for any algorithm that finds an optimal solution. Surely, lower bounds are much more difficult to establish than upper bounds. The lower-bound analysis is particularly challenging for problems with an exponential upper bound. Therefore, we typically analyze the complexity of difficult optimization problems in relation to other problems. In other words, we establish the complexity by proving that a problem is at least as difficult as another problem that is believed to be hard. Formally, this relation is captured in the classification of the problems on complexity classes.

### 2.2.1 NP-hardness

The framework of the complexity analysis was historically formulated in the terms of decision problems. As we have seen before, each optimization problem is at least as difficult as its decision version for a fixed $k$. We remark that the inverse is in general not true. For instance, for a constraint satisfaction problem over $m$ clauses each containing at most 2 binary variables, it is easy to decide if we can simultaneously satisfy all $k=m$ constraints. However, finding an assignment that maximizes the number of satisfied clauses is a much more difficult task de Beaudrap et al., 2021.

## NP-completeness of decision problems

The mathematical logic operates the concept of decidability. Decidability indicates if a decision problem can be solved at all by a computer program. Roughly speaking, a problem is decidable if there exists a program that terminates on every input with a correct positive or negative answer. We remark that by definition such a program may take an arbitrary finite time until termination.

In reality, we are not satisfied with a result that takes too long to be computed. Moreover, arguably all optimization problems of practical relevance are already decidable with a trivial exhaustive search. Therefore, rather than separating problems by their decidability we distinguish tractable and intractable problems. Conventionally, we consider a problem to be tractable if it can be solved in polynomial time. Tractable problems form the well-known complexity class P.

In quantum computing an analog to P is BQP - a class of problems solvable in polynomial time on quantum computers with an error probability of at most $1 / 3$ on all instances (positives as negatives). Problems that are solvable with high probability in polynomial time by a probabilistic classical algorithm form so-called class BPP. Classes PostBPP and PostBQP correspond to the extensions of BPP and BQP respectively where we allow the algorithms to perform post-selection on the outputs of probabilistic samplings.

We remark that in real-world applications the division on tractable and intractable problems is still too coarse, as computational algorithms scaling as $O\left(n^{d}\right)$ for high $d$ may take weeks even
on an input of a rather modest size.
In a typical situation, a decision problem admits a constructive solution. For instance, a proper $k$-coloring is a constructive solution for the decision version of the graph coloring problem that asks if there exists a $k$-coloring. A constructive solution may be interpreted as a certificate for $i \in I_{P}$. A certificate can be used in a procedure that verifies if the constructive solution indeed proves that the instance is positive. If the verification can be done in polynomial time, we say that the problem belongs to the class NP:

Definition 2.2.1 $A$ decision problem $\mathcal{P}_{D}$ is in the class NP if there exists a polynomial-time computable function $V$ called verifier such that

- for each positive instance $i \in I_{p}$ there exists a certificate $c(i) \in C$ of size $O(\operatorname{poly}(n))$.
- Given a instance $i$ and a certificate $c$ the verifier $V: I \times C \rightarrow\{0,1\}$ accepts a valid certificate for a positive instance $i \in I_{p}$, i.e.

$$
\begin{equation*}
\forall i \in I_{p} \exists c_{i},\left|c_{i}\right| \leq O(\operatorname{poly}(n)) \text { such that } V\left(i, c_{i}\right)=1 \tag{2.2}
\end{equation*}
$$

- For all negative instance $i \in I_{n}$ and any certificate $c$ the verifies rejects the proof:

$$
\begin{equation*}
\forall i \in I_{n} \forall c,|c| \leq O(\operatorname{poly}(n)) \text { we have } V(i, c)=0 \tag{2.3}
\end{equation*}
$$

We remark that the class NP can be equivalently defined via non-deterministic algorithms. More difficult counting problems that ask the number of accepting certificates for some $\mathcal{P}_{D} \in \mathrm{NP}$ (if any) form the class \# P.

A quantum version of the NP class is the class QMA. QMA is similarly defined as a class of problems for which positive instances can be certified in polynomial time. The crucial difference with NP is that in QMA the certificate can be a quantum state and the verifier has an access to a quantum computer. In addition, we remark that as the quantum program returns a probabilistic output the conditions $(2.2)$ and $(2.3)$ have to be satisfied with high probability rather than exactly.

In classical computing, the probabilistic verification of the conditions 2.2 and $(2.3)$ appears in the Arthur-Merlin complexity class.

A decision problem $\mathcal{P}_{D} \in \mathrm{NP}$ is called NP-complete if any other problem $\mathcal{P}_{D}^{\prime}$ from the class NP there is a polynomial procedure $f: I \rightarrow I^{\prime}$ transforming instances of $\mathcal{P}_{D}^{\prime}$ to instances of $\mathcal{P}_{D}$ in a way that preserves positivity:

$$
\begin{equation*}
i \in I_{p} \Longleftrightarrow f(i) \in I_{p}^{\prime} \tag{2.4}
\end{equation*}
$$

Informally speaking, NP-completeness of $\mathcal{P}_{D}$ implies that given a polynomial algorithm solving $\mathcal{P}_{D}$, we can solve any other problem from NP in polynomial time. The famous CookLevin theorem was the first proof of NP-completeness for the Boolean Satisfiability problem Cook, 1971. Since then many problems were proven to be NP-complete, mainly by reduction to the problems for which NP-completeness was previously established. An extensive list of NP-completeness results may be found in Garey and Johnson, 1979.

It turns out that among NP-complete problems some are more difficult than others. In particular, this is the case for decision problems that, while being untractable in polynomial time on the size of the instance $|i|$, can however be solved in $O(\operatorname{poly}(|i|, \max (i)))$ where $\max (i)$ is the number of the largest magnitude in the instance. Such problems are called pseudopolynomial. For instance, the decision version of the Knapsack problem is a pseudopolynomial problem.

On the opposite side, NP-complete problems that remain intractable even when the runtime is allowed to scale as $O($ poly $(|i|, \max (i)))$ are called strongly NP-complete Garey and Johnson, 1978.

## Complexity of optimization problems

An important connection between an optimization problem $\mathcal{P}$ and the corresponding decision problem $\mathcal{P}_{D}$ is that if $\mathcal{P} \in N P O$ then $\mathcal{P}_{D} \in \mathrm{NP}$ Ausiello et al., 1999. Optimization problems solvable in polynomial time belong to so-called class PO. The analogue to the $N P$-completeness in the framework of optimization problems is captured by the notion of $N P$-hardness.

Definition 2.2.2 An optimization problem $\mathcal{P}$ is NP-hard if every decision problem $\mathcal{P}_{D}^{\prime} \in \mathrm{NP}$ can be solved in polynomial time given an oracle for $\mathcal{P}$.

In other words, the NP-hardness of the optimization problem $\mathcal{P}$ implies that if there were a polynomial algorithm $\mathcal{A}$ that finds an optimal solution $x^{*}(i)$ for each instance $i$, then any decision problem in the class NP can be solved in polynomial time. Notably, if the decision version $\mathcal{P}_{D}$ of $\mathcal{P}$ is NP-complete, then $\mathcal{P}$ is NP-hard.

Many combinatorial optimization problems of practical importance are known to be NP-hard. Therefore, under the conjecture $\mathrm{P} \neq \mathrm{NP}$ they are not expected to be solved by polynomial-time algorithms. However, more careful analysis is needed to evaluate the practical difficulty of an NP-hard problem.

Firstly, NP-hardness reflects the worst-case behavior. It may happen that the average runtime performance of an algorithm is satisfactory even if it has an exponential scaling in the worst-case. Moreover, in practical applications a specific structure of real-world instances may render the computation easier.

From the theoretical point of view variations in the difficulty of NP-hard optimization problems are captured by the notions of pseudopolynomial problems and strongly NP-hard problems. An optimization problem $\mathcal{P} \in \mathrm{NPO}$ is pseudopolynomial if it can be solved by an algorithm polynomial on $|i|$ and $\max (i)$ where $\max (i)$ is the magnitude of the largest number occurring in the instance $i$. For example, in the Knapsack problem the value $\max (i)$ corresponds to the largest weight or the largest profit of an item in the instance. Notably, instances of a pseudopolynomial problem with polynomially bounded magnitudes $\max (i) \leq O(\operatorname{poly}(|i|))$ can be handled in polynomial time.

Strong NP-hardness for optimization problems is similar in spirit to the strong $N P$-completeness for decision problems. Formally, for each optimization problem $\mathcal{P}$ and a polynomial $p$ we can define a problem $\mathcal{P}^{\text {max, } p}$ that is the restriction of $\mathcal{P}$ on instances such that $\max (i) \leq p(|i|)$. A problem $\mathcal{P} \in \mathrm{NPO}$ problem is strongly $N P$-hard if there exists a polynomial $p$ such that $\mathcal{P}^{\text {max }, p}$ is NP-hard. For instance, many seminal optimization problems such as MaxCut, Maximum Independent Set, and graph coloring are strongly NP-hard Garey and Johnson, 1978.

As follows from the definition, strongly NP-hard problems cannot be pseudopolynomial.
Finally, in the context when the exact solution is unreachable in reasonable time we may instead consider near-optimal solutions. A near-optimal solution is usually referred to as approximate solution or approximate optimum. The complexity to reach an approximate solution of certain quality leads to a new classification of NP-hard optimization problems in terms of their approximability.

### 2.3 Approximation algorithms

NP-hardness usually implies that it is impossible to compute the exact optimal solution for the optimization problem on instances of relevant sizes in an acceptable time. A way to address
this inconvenience is to consider near-optimal solutions that can be potentially easier to find. Polynomial-time algorithms returning near-optimal solutions of NP-hard optimization problems are called approximation algorithms. Approximation algorithms can be divided into algorithms with guaranteed performance and heuristics.

### 2.3.1 Approximation ratio

We denote by $\mathcal{A}$ an approximation algorithm for an optimization problem $\mathcal{P}=(I, F, m)$. We distinguish several types of errors that are due to the non-optimality of the solution returned by $\mathcal{A}$ :

- absolute error. For an instance $i$ with optimal solution $m^{*}(i)$ the absolute error corresponds to the difference between the optimal solution and the solution returned by the approximation algorithm:

$$
\begin{equation*}
E(i)=\left|m^{*}(i)-m(x, \mathcal{A}(i))\right| \tag{2.5}
\end{equation*}
$$

- relative error. This metric captures a relative deviation from the optimum value:

$$
\begin{equation*}
r(i)=\frac{\left|m^{*}(i)-m(x, \mathcal{A}(i))\right|}{\max \left\{m^{*}(i), m(i, \mathcal{A}(i))\right\}} \tag{2.6}
\end{equation*}
$$

In the current thesis we extensively use another performance metric called approximation ratio. For a minimization problem the approximation ratio of an algorithm $\mathcal{A}$ on an instance $i \in I$ is $\alpha(i)=\frac{m^{*}(i)}{m(i, \mathcal{A}(i))}$. Alternatively, for the maximization problem the approximation ratio is $\alpha(i)=\frac{m(i, \mathcal{A}(i))}{m^{*}(i)}$. We remark that the approximation ratio is always lower or equal to one. In the case of equality $\alpha(i)=1$ we can safely claim that the approximation algorithm found an optimal solution.

In the formal analysis of the performance of approximation algorithms we are typically interested in the worst-case behavior. We call an algorithm an $\alpha$-approximate algorithm if it has the approximation ratio of at least $\alpha$ on any instance of the optimization problem, i.e.

$$
\begin{equation*}
\forall i \in I: \alpha(i) \geq \alpha \tag{2.7}
\end{equation*}
$$

We say that the bound 2.7 is tight if there exists at least one instance $i^{*} \in I$ such that the approximation ratio $\alpha\left(i^{*}\right)$ is precisely $\alpha$.

An optimization problem $\mathcal{P}$ is called $\alpha$-approximable if there exists a polynomial-time $\alpha$ approximate algorithm solving it.

NP-hard problems have different approximability properties. Some can be approximated to any desired error rate while for others the relative quality of the best possible algorithm is always poor. In the latter case, the approximation ratio usually decreases with the size of the instance.

In this section, we consider classes of problems with different achievable approximation guarantees.

### 2.3.2 APX

Definition 2.3.1 APX is a class of optimization problems such that there exists a polynomialtime approximate algorithm with approximation ratio $\alpha$ bounded by a constant $\alpha_{l} \leq \alpha$.

For instance, the MaxCut problem belongs to the class APX Goemans and Williamson, 1995.
To show that an optimization problem is in APX it suffices to provide an algorithm that has a constant approximation ratio $\alpha_{l}$. Usually, we are interested in known if the achieved ratio is tight, i.e. if there can be a (probably still undiscovered) algorithm that has a better performance. For some problems it can be shown that if it were possible to approximate the solution with a factor $\alpha>\alpha_{u}$, then $\mathrm{P}=\mathrm{NP}$. If the approximation ratio of the algorithms is $\alpha_{l}=\alpha_{u}$ then in a sense it provides the best achievable approximation performance.

However, for many problems the strong inapproximability results leave a gap $\alpha_{u}-\alpha_{l}>0$ with the best-known algorithms. If this is the case, we can check if a stronger bound $\alpha<\alpha_{u}^{\prime}<\alpha_{u}$ can be derived assuming complexity conjectures that are weaker than $\mathrm{P} \neq \mathrm{NP}$. Such inapproximability result was derived, for instance for the MaxCut problem assuming the Unique Game Conjecture (UGC) Khot et al., 2007. The UGC-bound for the MaxCut implies, in particular, the tightness of the Goemans-Williamson approximation ratio $\alpha=0.878 \ldots$ Khot et al., 2007.

In general, there exist problems in NPO that are even not in APX. In other words, for such problems there is no algorithm returning a solution of any guaranteed constant ratio unless $P=$ NP. A seminal example of a problem $\mathcal{P} \notin$ APX is the Maximum Independent Set problem Håstad, 1999.

The proofs $\mathcal{P} \notin \mathrm{APX}$ can be derived using various techniques. For instance, we can show that there exists an NP-complete decision problem that can be reduced to the approximation problem with arbitrary fixed ratio $\alpha$. Such approach is called gap technique. An illustration of the gap technique for the inapproximability bound of the Traveling Salesperson problem can be found in Ausiello et al., 1999]. More involved proofs are usually derived from a fundamental result of the complexity theory called PCP-theorem Arora and Safra, 1998.

### 2.3.3 PTAS

There exist optimization problems for which an exact solution is difficult to compute but it can nevertheless be approximated to any desired ratio. Formally speaking, for such a problem there exists a family of $\alpha$-approximation algorithms for all $\alpha<1$. Such a family is called polynomialtime approximation scheme or PTAS for short. A slightly weaker notion is that of asymptotic approximation scheme or, equivalently, a family of polynomial-time algorithms such that the approximation ratios $\alpha<1$ are achieved in the limit of large instances.

We remark that usually an approximation scheme takes more time to find solutions for better approximation ratios. In fact, for a polynomial-time approximation scheme, the dependence on the desired solution quality can still be exponential. This is the case, for instance, for the PTAS solving the smart scheduling problem considered in the current thesis Skutella and Woeginger, 2000.

We say that an optimization problem $\mathcal{P} \in$ NPO belongs to the class PTAS if it admits a polynomial-time approximation scheme. An example of a problem in PTAS is the Maximum Independent Set (MIS) in a planar graph Ausiello et al., 1999. We remark that the MIS problem on non-planar graphs is much harder, to such an extent that it is not even in APX.

Importantly, if $P \neq$ NP the class PTAS is strictly contained in APX, i.e. PTAS $\subsetneq$ APX.

### 2.3.4 FPTAS

The runtime of PTAS can be significantly increased for more demanding performance guarantees. However, for some problems there exists a family of $\alpha$-approximation algorithms with the
complexity that scales polynomially both on the size of the instance $|i|$ and on the desired ratio $\alpha$. Such family is called a fully-polynomial approximation scheme (FPTAS). For an NP-hard optimization problem this is in a sense the strongest possible approximability result. Roughly speaking, the existence of FPTAS for $\mathcal{P} \in$ NPO implies that even if the exact solution is difficult to compute, it can be efficiently approximated to any desired ratio. The class of the optimization problem that admits such a scheme is denoted as FPTAS. An example of the problem from FPTAS is the Maximum Knapsack problem.

The book Garey and Johnson, 1978 identifies certain conditions on optimization problems that guarantee that the problem does not possess an FPTAS. For instance, this is the case for polynomially bounded strongly NP-hard problems. We recall that a problem is called polynomially bounded if its measure function is bounded, i.e. $\forall i \in I$ and $\forall y \in F(i): m(x, y) \leq p(|x|)$ for some polynomial $p$.

On the opposite sense, an existence of FPTAS for a problem $\mathcal{P}$ such that $m^{*}(x)<p(|x|, \max (x))$ implies that $\mathcal{P}$ is pseudo-polynomial.

In most cases an FPTAS for a problem $\mathcal{P}$, if it exists, is derived from dynamic programming formulation. For instance, if the dynamic programming formulation has a special property introduced in Woeginger, 2001 under the name DP-benelovence, then the problem admits an FPTAS. DP-benevolence is a particularly interesting sufficient condition for the existence of FPTAS as $i$ ) it is easy to verify $i i$ ) it captures a relatively large family of optimization problems. Typically, FPTAS is derived from a dynamic programming formulation with rounding-the-inputdata technique Sahni, 1976. Another approach originally introduced in Ibarra and Kim, 1975 is called trimming-the-state-space technique.

We remark that the smart scheduling problem considered in chapter 4 admits an FPTAS.

### 2.3.5 Heuristics

Approximation algorithms with guaranteed ratios are crucial in the evaluation of the theoretical complexity of optimization problems. In practice, however, they are often outperformed by heuristics. Heuristics are polynomial-time programs without performance guarantees that, nevertheless, demonstrate good behavior in experimental evaluations on instances of practical interest. On the other side, heuristics may have very poor performance in the worst case. Some works include in the definition of heuristics methods with an exponential scaling in the worst case if they are relatively rapid on relevant instances.

Heuristics are usually inspired by one of two traditional approaches. The first one consists of an iterative construction of the solution from a partial solution. For instance, greedy algorithms implement such a constructive approach.

Alternatively, local search heuristics start from a complete feasible solution and try to improve it. For this purpose, they search for an improving solution inside the neighborhood of the actual solution. Contrary to greedy algorithms that are usually tailored for a specific application, local search is a very general framework. It is implemented by various metaheuristics such as simulated annealing Kirkpatrick et al., 1983 and tabu search Battiti and Protasi, 2001, Kochenberger et al., 2013. The performance of local search heuristics heavily depends on the quality of the initialization point. Usually, the initialization point is found by a constructive heuristic. In addition, local-search metaheuristics have a large number of parameters that should be efficiently fine-tuned to assure good practical performance.

In the last decades emerged heuristics inspired by the natural selection process. For instance, an imitation of natural selection was realized in genetic algorithms Marchiori, 1998 and the differential evolution algorithm.

The domain of heuristic algorithms is large and continuously developing. A comprehensive overview of the most important heuristic methods for combinatorial optimization can be found in Burke and Kendall, 2006.

### 2.4 Semidefinite programming

Semidefinite programming is a subfield of convex optimization over continuous variables. In the field of combinatorial optimization semidefinite programs were found to be extremely useful in the design of approximation algorithms with guaranteed approximation ratios. We refer to Gärtner and Matoušek, 2013 for a general introduction to the field.

A semidefinite program is a problem of maximizing a linear function over $n^{2}$ continuous variables $x_{i, j}$. Variables may be arranged in a matrix $X=\left(\begin{array}{lll}x_{1,1} & \ldots & x_{1, n} \\ & \ddots & \\ x_{n, 1} & \ldots & x_{n, n}\end{array}\right)$.

Semidefinite program (SDP) is formulated as follows:

$$
\begin{array}{ll}
\min & \operatorname{Tr}\left(A^{T} X\right) \\
& \operatorname{Tr}\left(C_{k}^{T} X\right) \leq b_{k}, \\
X \succeq 0 & k \in[1, \ldots, m] \tag{2.10}
\end{array}
$$

where $\operatorname{Tr}(Y)$ denotes the trace of the matrix $Y$ and $X \succeq 0$ means that the matrix $X$ is positive semidefinite. We remark that the objective function and the constraints are linear functions on variables $x_{i, j}$ as $\operatorname{Tr}\left(A^{T} X\right)=\sum_{i, j} a_{i, j} x_{i, j}$.

Under certain relatively mild conditions, SDP can be solved up to any desired accuracy in polynomial time. In particular, trace-bounded instances can be efficiently approximated with an algorithm implementing multiplicative weight update Arora and Kale, 2016.

## Goemans-Williamson algorithm

Many optimal approximation algorithms for combinatorial problems are based on semidefinite programming. Usually, an SDP is formulated as a continuous relaxation of a problem over integer variables. For instance, this is the case for the Goemans-Williamson algorithm for the MaxCut problem Goemans and Williamson, 1995.

MaxCut problem on a graph $G=(V, E)$ is a search of a partition $V=V_{1} \sqcup V_{2}$ such that the number of crossing edges $(u, v): u \in V_{1}, v \in V_{2}$ is maximized. We assign one variable $z_{v} \in\{-1,1\}$ to each node. The partition corresponding to a variable assignment $\boldsymbol{z} \in\{-1,1\}^{n}$ is exactly $V_{1}=\left\{v \mid z_{v}=1\right\}$ and $V_{2}=\left\{v \mid z_{v}=-1\right\}$. The MaxCut problem may be written as:

$$
\begin{equation*}
\max \sum_{(u, v) \in E} \frac{1-z_{u} z_{v}}{2} \tag{2.11}
\end{equation*}
$$

The Goemans-Williamson algorithm is based on a semidefinite relaxation of the MaxCut problem. The value of the relaxation provides an upper bound on the size of the optimal cut $c^{*}(G)$.

In the Goemans-Williamson algorithm we associate to each node a vector $y_{v}$ on the unit sphere $\mathbb{S}^{n}$ in $\mathbb{R}^{n}$, i.e. $S^{n}=\left\{x \in \mathbb{R}^{n}| | x \mid=1\right\}$.

With vectors from the unit sphere, a "continuous" version of the MaxCut can be written as:

$$
\begin{align*}
& \max \sum_{(u, v) \in E} \frac{1-y_{u}^{T} y_{v}}{2}  \tag{2.12}\\
& y_{v} \in S^{n}, \quad \forall v \in V \tag{2.13}
\end{align*}
$$

We remark that the formulation over $y_{v} \in \mathbb{S}^{n}$ is the relaxation of the integer problem 2.11. Indeed, an optimal solution for the integer problem leads to a feasible solution $y_{v}=$ $\left[0, \ldots, 0_{v-1}, z_{v}, 0_{v+1} \ldots, 0\right]$ for the vector problem with the same objective value. Therefore, the optimum solution for (2.12) has the value at least $c^{*}(G)$.

From the formulation over $y_{v}$ we can get an equivalent SDP program. We denote by $x_{u, v}=$ $y_{u}^{T} y_{v}$. An equivalent SDP program is:

$$
\begin{array}{ll}
\max & \sum_{(u, v) \in E} \frac{1-x_{u, v}}{2} \\
& \\
& x_{u, u}=1,  \tag{2.16}\\
& X \succeq 0
\end{array} \quad u \in[1, \ldots, n]
$$

In the SDP formulation the condition $y_{v} \in S^{n}$ was translated to $x_{u, u}=1$. Proof of equivalence of both problems can be found in Gärtner and Matoušek, 2013. We remark that this proof is constructive, i.e. it explicitly specifies how to recover the vectors $y_{v}$ from the matrix $X$. Therefore, an optimal solution $X^{*}$ of the SDP relaxation associates to each node a vector $y_{v}^{*}$ from the unit sphere $S^{n}$.

An integer solution for the MaxCut is obtained from vectors $y_{v}^{*}$ with a help of a random rounding plane. To be exact, the algorithm chose a vector $\rho \in \mathbb{R}^{n}$ at random. The integer solution is recovered from the separation of $\mathbb{R}^{n}$ on two subspaces generated by $\rho$ :

$$
z_{v}^{*}= \begin{cases}1, & \text { if } \rho^{T} y_{v}^{*} \geq 0  \tag{2.17}\\ -1, & \text { otherwise }\end{cases}
$$

It was proven in Goemans and Williamson, 1995 that such an algorithm achieves an approximation ratio of $0.878 \ldots$. Moreover, this ratio is tight Karloff, 1999 and optimal under the Unique Game Conjecture Khot et al., 2007.

## Part II

## Quantum algorithms for Smart Charging problems

## Chapter 3

## Quantum algorithms for combinatorial optimization

### 3.1 Motivation

A lot of combinatorial optimization problems such as 3-SAT, Maximum Cut, Maximum Independent Set, Graph Coloring, Traveling Salesman, Integer Linear Program and many others are known to be NP-hard. Therefore, unless the widely believed conjecture $P \neq N P$ is false, these problems don't admit polynomial-time algorithms solving them to optimality Garey and Johnson, 1979]. There exists, however, exact algorithms that explore the solution space either explicitly (a brute-force approach) or implicitly by using relaxations that in some cases allow to efficiently eliminate suboptimal solutions (Branch \& Bound). Obviously, exact algorithms have an exponential runtime in the worst-case.

Another way to practically address NP-hard problems consists in relaxing the requirement of optimality. Indeed, in a restricted-time regime, it is preferable to have a suboptimal solution (approximate optimum) rather than wait too long for the exact optimum to be computed. Two types of polynomial algorithms implement this idea: approximation algorithms and heuristics Ausiello et al., 1999. While the approximation algorithms provide theoretical guarantees on the quality of the obtained solution, the heuristics are usually validated by experimental demonstrations. In this setting, the performance is determined by the resource consumption (runtime and memory) but also by the quality of the solution (the approximation ratio). The approximation ratio indicates how far the approximate solution is from the exact optimum.

Hence the question arises if quantum computers can provide an advantage for combinatorial optimization applications, either by speeding up the execution of exact or approximate routines or via new heuristics that improve the quality of the approximate solution over state-of-the-art classical methods either in theory or in practice.

### 3.1.1 Quantum acceleration of classical routines

Indeed, the integration of quantum routines can accelerate both approximate and exact optimization algorithms.

The exact optimization typically profits from Grover's algorithms for the search in an unstructured database Grover, 1996. Grover's original algorithm assures a quadratic speedup over the best-possible classical counterpart [Grover, 1996]. This technique can provide a polynomial speedup for routines solving the SAT problem Ambainis, 2005, for the Branch \& Bound

Montanaro, 2020 or for the nested search for structured problems Cerf et al., 2000.
Furthermore, a quantum algorithm with a quadratic speedup was suggested for the semidefinite programming problem (SDP) van Apeldoorn et al., 2020. SDP followed by a randomized rounding often appears as a subroutine in the approximate algorithms for binary optimization with the best-established performance guarantees Gärtner and Matoušek, 2013. The quantum speedup for SDP implies a proportional acceleration for these approximation routines for such problems as MaxCut G.S L. Brandão et al., 2022.

However, even if the asymptotic time complexity of the algorithms mentioned above improves over their classical counterparts the exponential scaling is preserved and no significant improvement over a random guess is expected in a constrained time budget McClean et al., 2021. Moreover, for asymptotic polynomial speedups a practical runtime advantage is not expected to be observed on the NISQ machines Babbush et al., 2021. This limitation is due to heavy errorcorrection schemes that are necessary to ensure the correctness of the circuit execution. Error correction seems to be the genuine bottleneck for practical quantum computing: even for quantum heuristics (believed to be NISQ-compatible) fault-tolerant realizations induce a significant resource consumption Sanders et al., 2020.

### 3.1.2 Quantum heuristics

On the other hand, quantum heuristics can potentially improve over the classical ones in terms of the quality of the approximate solution. In addition, these heuristics (among them Quantum Annealing (QA), Quantum Approximate Optimization Algorithm (QAOA) Farhi et al., 2014a and Variational Quantum Eigensolver (VQE) Peruzzo et al., 2014) are expected to be wellsuited for the NISQ devices Preskill, 2018. Indeed, these algorithms can be implemented in a low-depth regime and because of their approximate nature, one can bypass the explicit error correction.

Unfortunately, the actual error rate leaves serious doubts on the ability of quantum heuristics to demonstrate quantum advantage Albash et al., 2017, França and García-Patrón, 2021, Lotshaw et al., 2022]. The limiting result [França and García-Patrón, 2021] is due to the existence of a thermal Gibbs state $G_{T}=e^{-H / T} / \operatorname{Tr}\left(e^{-H}\right)$ that has the same mean energy $\langle H\rangle$ as the quantum state returned by a noisy implementation of a quantum heuristic. The corresponding temperature $T$ of the state depends on the noise level. Crucially, for sufficiently high noise, the temperature $T$ is high enough to allow an efficient classical sampling from $G_{T}$. Therefore, the quantum advantage is unlikely unless at least one of the two conditions is fulfilled: the problem matches the hardware-native architecture or the noise level is significantly reduced (by two orders of magnitude from the actual values).

This chapter introduces quantum algorithms for the combinatorial optimization that are believed to be adapted for NISQ: Quantum Annealing (QA), Quantum Approximate Optimization Algorithm ( $Q A O A$ ) and Recursive $Q A O A$ ( $R Q A O A$ ). We concentrate on the analysis in the noisefree regime and only briefly mention the results about the impact of decoherence on physical realizations.

## Ground state problem

Traditionally, a quantum combinatorial optimization algorithm receives as input an observable $H$ called Ising Hamiltonian. From the physical side, Ising Hamiltonians are used to describe the energy levels of quantum many-body systems with pair interactions such as spin-glass models Zarinelli, 2012. Mathematically, an Ising Hamiltonian is a Hermitian operator diagonal in
computational basis that can be expressed as a polynomial of degree 2. The original framework was further expended to deal with higher-degree interactions Hadfield, 2018, Glos et al., 2022 and to incorporate hard restrictions on the solution space Hadfield et al., 2019, Hadfield, 2018.

The problem we are interested in consists in finding the ground state (eigenstate with the smallest eigenvalue) for a given Hamiltonian. Quantum heuristics will typically return not exactly the ground state but a low-energy state that is considered as an approximate solution.

In order to apply quantum algorithms to combinatorial optimization problems we first need to map the solution space $S$ to a subset $S_{b}$ of computational basis states of an $n$-qubit system $|x\rangle, x \in\{0,1\}^{n}$. Then we need to define a corresponding diagonal Hamiltonian $H$ such that for every feasible solution its eigenvalue corresponds to the value of the objective function:

$$
\begin{equation*}
H|x\rangle=f(x)|x\rangle, \quad|x\rangle \in S_{b} \tag{3.1}
\end{equation*}
$$

The translation is straightforward for Quantum Binary Optimization Problems (QUBO) and relatively easy for a lot of NP-hard optimization problems Lucas, 2014. Section 3.2 contains a detailed description of the Ising Hamiltonian ground state problem and related complexity results.

## Quantum evolution in algorithms for optimization

In section 3.3 we start the exploration of quantum optimization algorithms with the Adiabatic Algorithm (AA) originally introduced in Farhi et al., 2000. By definition, the Adiabatic Algorithm returns an exact solution for a given optimization problem. The AA initializes the system in an easy-to-prepare ground state of some Hamiltonian $H_{\text {init }}$. The state evolves under the time-dependent Hamiltonian which gradually changes from $H_{\text {init }}$ to the $H_{\text {final }}$ where the ground state of $H_{\text {final }}$ corresponds to the solution of the optimization problem.

According to Schrödinger's equation, the time evolution of any quantum system is governed by the Hamiltonian acting on it:

$$
\begin{equation*}
i \hbar \frac{\partial_{t}|\psi(t)\rangle}{\partial t}=H(t)|\psi(t)\rangle \tag{3.2}
\end{equation*}
$$

where $\hbar=1.054 \cdots \times 10^{-34} \mathrm{Js}$ is the reduced Plank constant.
The adiabatic theorem states that if the evolution is sufficiently slow (with respect to the difference between the ground state energy and the energy of the first excited state usually called spectral gap), the system remains close to the instantaneous ground state during the whole process Jansen et al., 2007, Amin, 2009a. Therefore, with high probability the measure output of the final state will be the ground state of $H_{\text {final }}$ - precisely the exact solution of the optimization problem. The downside, although consistent with the conjecture NP $\nsubseteq \mathrm{BQP}$, is that "sufficiently slow" is generally difficult to compute and sometimes the best bounds imply an exponential evolution runtime for difficult instances of NP-hard problems van Dam et al., 2001 or even for some classically tractable problems Reichardt, 2004, Albash and Lidar, 2018. Even worse scaling when the gap vanishes faster than exponentially was observed for random instances of the NP-hard Exact 3 Cover problem Altshuler et al., 2010.

In addition, even in the case of a polynomial gap, the adiabatic runtime condition is difficult to satisfy on machines with a realistic coherence time limit. Therefore the Adiabatic Algorithm remains mostly a theoretical construction. It has, however, inspired several polynomial-time heuristics such as Quantum Annealing and QAOA Farhi et al., 2014a.


Figure 3.1: Hamiltonians governing the evolution of the quantum system at different moments of the execution of optimization algorithms.

The Quantum Annealing changes the Hamiltonian in a similar way that the Adiabatic Algorithm but much faster than allowed by the adiabatic condition (see figure 3.1). The continuous evolution is performed on special-purpose devices called quantum annealers, and the runtime is typically limited by the hardware characteristics. In the DWave Advantage 6.1 the available annealing time ranges from $0.5 \mu s$ to $2000 \mu s$ Systems, 2022. The Quantum Annealing returns a solution with no guarantees on its quality and an experimental benchmark is required to conclude about the potential of this approach. An experimental study reported in King et al., 2015 claims that, under a fixed time budget, Quantum Annealing outperforms the state-of-the-art classical solvers on some optimization problems. On the other hand, the work Martin-Mayor and Hen, 2015 demonstrates that QA scales worse with the hardness of the instance than the thermal classical algorithm. In section 3.4 we report some experimental studies concerning the application of Quantum Annealing to optimization problems.

In contrast to the Quantum Annealing that requires a special-purpose analog device, the Quantum Approximate Optimization Algorithm (see section 3.5) adapts the ideas of the Adiabatic Algorithm to the universal gate-based quantum computers Farhi et al., 2014a. Beyond the usual considerations about the relative power of two types of hardware Barends et al., 2016 the gate-based computers, contrary to existing annealers that allow only pair interactions, can directly deal with polynomials of higher degrees Glos et al., 2022. This property is useful in practice. For instance, the E3LIN2 problem for which the quantum advantage was firstly obtained and then beaten by a new classical algorithm, has cubic terms Farhi et al., 2014b. In addition, the ability to directly treat higher degree interactions sometimes leads to implementations with lower resource complexity Glos et al., 2022, Fuchs et al., 2020.

The QAOA dynamics is ruled not by a continuous interpolation of two Hamiltonians but by a sequence of pulses ("bangs") that jumps between $H_{\text {init }}$ and $H_{\text {final }}$ (see figure 3.1). The total amount of pulses $p \in \mathbb{N}$, called depth, is predefined by the user while the durations of bangs are optimized directly inside the algorithm. The depth of the quantum circuit that performs the evolution scales proportionally with the QAOA depth, so due to the limited coherence time the near-future implementations are likely to stay in the regime $p=O(1)$.

Hence, the runtime of QAOA is determined by two factors: the depth $p$ and the total sum of pulse durations. It was observed that if the latter scales as the inverse polynomial on the problem's size then no quantum advantage is possible Hadfield et al., 2021.

[^2]
## Which quantum heuristic is better?

Both Quantum Annealing and QAOA alternate between $H_{\text {init }}$ and $H_{\text {final }}$ :

$$
H(t)=(1-g(t)) H_{\text {init }}+g(t) H_{\text {final }}
$$

where for the annealing the function $g(t)$ is continuous and for QAOA $g(t)$ has the pulse form. The work Yang et al., 2017 suggested that for a fixed time limit a "bang-bang" alternation protocol (as in QAOA) is optimal in terms of the achievable approximation ratio. More recently, another study reported that one of the assumptions used to prove the "bang-bang" optimality doesn't hold in general case Brady et al., 2021. This study shows that for a fixed evolution time, the bang-anneal-bang sometimes is better, in particular for some instances encoded by Ising Hamiltonians. However, in a more realistic regime where the depth is given but the sum of pulse durations is not fixed QAOA's performance rapidly approaches the one of the bang-annealbang protocol Brady et al., 2021. Moreover, QAOA has a practical advantage: the pulse control is expected to be easier to optimize and execute than a bang-anneal-bang protocol.

The deep connection between QAOA and Quantum Annealing is also reflected by the fact that the discretized annealing schedule, when used to initialize the search of QAOA pulse durations, generally leads to better QAOA schedules compared to the random initialization Sack and Serbyn, 2021.

Recently, there has been remarkable progress in the physical implementation of annealers: 5000 qubits are accessible on DWave Advantage system 6.1 Systems, 2022 compared to relatively modest 127 qubits on the universal IBM Eagle quantum computer Dial, 2022. Hence, QA is more amenable to experimental analysis. At the same time, due to the lack of tools, the theoretical analysis of the QA is significantly harder and only non-tight guarantee bounds Braida et al., 2022 and no-go results were established Moosavian et al., 2022.

On the other side, for QAOA a lot of theoretical results were obtained. Between these results, the particularly important ones are supremacy Farhi and Harrow, 2016 and universality Lloyd, 2018. Some tight performance guarantees for special instances of MaxCut Farhi et al., 2014a, Hadfield, 2018, Wurtz and Love, 2021, Marwaha, 2021, Basso et al., 2022 and E3LIN2 Farhi et al., 2014b were established, although no quantum advantage was proven. The works Bravyi et al., 2020, Farhi et al., 2020, Chou et al., 2021 report limiting results on the approximation ratio for constantdepth QAOA on MaxCut and Maximum Independent Set problems. Finally, a general analytical framework for QAOA was suggested in Hadfield et al., 2021.

Both QA and QAOA are expected to approach the exact optimum at the infinite time limit, but the convergence speed is unknown. Several results were established for QAOA that demonstrate that at least a logarithmic depth is required in order to outperform the classical worst-case approximation ratio guarantees for MaxCut Bravyi et al., 2020, Max-k-XORSAT Chou et al., 2021 and MIS Farhi et al., 2020. Most of these results were then generalized to the short-time Quantum Annealing using Lieb-Robinson bound Moosavian et al., 2022. A question if a bounded-depth QAOA can show an advantage over local classical algorithms is studied in Hastings, 2019, Barak and Marwaha, 2022.

The last algorithm presented in this chapter is Recursive QAOA (RQAOA) (see section 3.6). RQAOA was originally designed to deal with locality that was often mentioned as the limiting condition for the QAOA performance Bravyi et al., 2020, Farhi et al., 2020, Hastings, 2019, Chou et al., 2021. RQAOA proceeds by a recursive elimination of variables, where at each step the variable to eliminate is selected using the optimal QAOA state for the restricted problem. The promise behind this is that, contrary to constant-depth QAOA, RQAOA ends up seeing the global connectivity structure. As a demonstration, the work Bravyi et al., 2020 shows that at
depth $p=1$ RQAOA finds the optimal solution on 1D Ising chains given by the Hamiltonian $H=\sum_{i=1}^{n} h_{i} Z_{i} Z_{i+1}$, while QAOA's performance is bounded for any fixed $p$ Mbeng et al., 2019a.

From the experimental point of view, a particularly interesting fact about RQAOA is that for $p=1$ it can be efficiently simulated on a classical computer Bravyi et al., 2020, Egger et al., 2021.

In brief, quantum algorithms for combinatorial optimization can be classified by their runtime and the achieved approximation ratio. Thanks to the adiabatic theorem, Adiabatic Algorithm is guaranteed to find an exact solution but potentially in exponential time. QAOA and RQAOA by construction scale polynomially and, while there exist bounds on the approximation ratio in special cases, their performance is mostly evaluated experimentally. In actual implementations Quantum Annealing has a constant runtime and almost only experimental performance demonstrations.

### 3.2 Ising Hamiltonian

Ising Hamiltonian was originally introduced to model spin interactions in crystals. It was actively used to study such important phenomena as phase transitions or the dependency of thermodynamic functions on different parameters Ising, 1925, Wikipedia, 2022a.

## 2-local Hamiltonians

As most physical interactions are pairwise, quantum computes are expected to natively support only 2-local connections Rieffel et al., 2014. The Ising model is a restriction of 2-local Hamiltonians. A general 2-local Hamiltonian can be written as

$$
\begin{equation*}
H=\sum_{(u, v) \in E}\left[J_{u, v}^{x} X_{u} X_{v}+J_{u, v}^{y} Y_{u} Y_{v}+J_{u, v}^{x} Z_{u} Z_{v}\right]+\sum_{u \in V}\left[h_{u}^{x} X_{u}+h_{u}^{y} Y_{u}+h_{u}^{z} Z_{u}\right] \tag{3.3}
\end{equation*}
$$

where $E$ defines the set of non-zero interactions and $X, Y$ and $Z$ are Pauli matrices. Negative coefficients $J_{u, v}<0$ correspond to ferromagnetic couplings and positive ones - to antiferromagnetic. If the graph $E$ allows only near-neighbor interaction in $k$-dimensional lattice the corresponding Hamiltonians are called $k D$.

The ground state decision problem for such a Hamiltonian is known to be a QMA-complete Kempe et al., 2006. QMA is a quantum analog of NP, and QMA-completeness is in a sense the strongest hardness result achievable for decidable problems in quantum computing.

Several special cases may be defined by imposing restrictions on the coefficients $J$ and $h$. We list below the most famous examples.

- Heisenberg model where $J_{u, v}^{x}=J_{u, v}^{y}=J_{u, v}^{z}=J_{u, v}$ :

$$
\begin{equation*}
H_{\text {Heisenberg }}=\sum_{(u, v) \in E} J_{u, v}\left(X_{u} X_{v}+Y_{u} Y_{v}+Z_{u} Z_{v}\right) \tag{3.4}
\end{equation*}
$$

- $X Y$ model where $J_{u, v}^{x}=J_{u, v}^{y}=J_{u, v}$ and $J_{u, v}^{z}=0$ :

$$
\begin{equation*}
H_{X Y}=\sum_{(u, v) \in E} J_{u, v}\left(X_{u} X_{v}+Y_{u} Y_{v}\right) \tag{3.5}
\end{equation*}
$$

- Ising model where all coefficients are set to zero except $h_{u}^{z}$ and $J_{u, v}^{z}$ :

$$
\begin{equation*}
H_{I}=\sum_{(u, v) \in E} J_{u, v} Z_{u} Z_{v}+\sum_{u \in V} h_{u}^{z} Z_{u} \tag{3.6}
\end{equation*}
$$

A typical example of an Ising Hamiltonian is the Sherington-Kirkpatrik mean-field model for a spin glass $H=\frac{1}{\sqrt{N}} \sum_{(u, v) \in G} J_{u, v} Z_{u} Z_{v}$ where the coefficients $J_{u, v} \sim \mathcal{N}(0,1)$ are independent identically distributed standard Gaussian random variables Sherrington and Kirkpatrick, 1975.

Many combinatorial optimization problems may be encoded in a ground-state search of an Ising Hamiltonian Lucas, 2014. Therefore, the problem is NP-hard and exact algorithms (such as Branch \& Cut Simone et al., 1995) take an exponential time to find the solution.

## Transverse-field Ising Hamiltonian

A transverse-field Ising Hamiltonian:

$$
\begin{equation*}
H=\sum_{(u, v) \in E} J_{u, v} Z_{u} Z_{v}+\sum_{u \in V} h_{u}^{z} Z_{u}+\sum_{u \in V} \sigma_{u} X_{u} \tag{3.7}
\end{equation*}
$$

has supplementary mixing component $H_{\sigma}=\sum_{u \in V} \sigma_{u} X_{u}$. The original Ising model is diagonal in the computational basis: the evolution $|x\rangle \rightarrow e^{i t H_{I}}|x\rangle$ under $H_{I}$ preserves the probability distribution over basis states $|x\rangle, x \in\{0,1\}^{n}$. Therefore, the mixing component $H_{\sigma}$ is essential in quantum algorithms for optimization in order to "balance" the probability weight to better solutions. Mixing is necessary to explore the solution space in a search of an optimal solution. On the other hand, the physical interpretation of mixing term as the kinetic energy allows us to demonstrate that sometimes it can drive the evolution out of the optimum McClean et al., 2021.

In some special cases, polynomial algorithms are known for ground-state search of transversefield Ising Hamiltonians. For instance, a fully-polynomial time approximation scheme for a Hamiltonian with fully ferromagnetic interactions is given in Bravyi and Gosset, 2017. Another simple case is that of 1D transverse-field Ising model with $h_{u}^{z}=0$ Pfeuty, 1970.

An interesting question is the quantum complexity of the ground-state search problem for the transverse-field Ising model. Notably, such Hamiltonians are stoquastic, i.e. their off-diagonal elements are non-positive in the computational basis Albash and Lidar, 2018. Stoquastic Hamiltonian ground state problem belongs to classical Artur-Merlin complexity class and, contrary to the case of general 2-local Hamiltonians, it is not believed to be QMA-complete Bravyi et al., 2006.

From the practical point of view it is interesting to know if there exist instances such that quantum computer can efficiently find their ground states while the best-known classical algorithm takes exponential time. For the moment, only limited quantum speedup was observed: the work Hastings and Freedman, 2013 presents instances that are polynomially solvable on quantum computers but for which some special classical methods like Simulated Annealing take exponential time.

### 3.2.1 Quadratic Unconstrained Binary Optimization (QUBO)

A combinatorial optimization problem that is naturally equivalent to the ground state search of an Ising Hamiltonian is Quadratic Unconstrained Optimization Problem (QUBO):

$$
\begin{equation*}
\min _{x \in\{0,1\}^{n}} x^{T} Q x \tag{3.8}
\end{equation*}
$$

where $Q$ is a symmetric $n \times n$ matrix. The connectivity pattern of the function generates a graph $G=(V, E)$ where each node $G=(V, E)$ corresponds to a vertex and each edge corresponds to a quadratic term with non-zero coefficient.

The objective goal can be freely changed to maximization by multiplying the function by -1 .
A lot of combinatorial optimization problems such as various forms of assignment problems, sequencing, and ordering Rieffel et al., 2014, Alidaee et al., 1994, MaxCut Kochenberger et al., 2013, set covering Alidaee et al., 2008 and many others may be represented in the QUBO form (often in multiple ways). Moreover, it was shown that any unconstrained binary optimization problem can be reduced to QUBO Anthony et al., 2017. A list of applications and solution methods may be found in the relatively recent survey Kochenberger et al., 2014.

The curious thing about QUBO is that sometimes generic QUBO-solvers (such as Tabu search) rivals the performance of the specially-designed algorithms for different applications, for example graph coloring Kochenberger et al., 2005, MaxCut Kochenberger et al., 2013 and set packing Alidaee et al., 2008.

For each QUBO the corresponding Ising Hamiltonian $H$ may be found by replacing each binary variable $x_{i}$ by an operator $\frac{I-Z_{i}}{2}$ where $I$ is identity and $Z$ is the Pauli-Z operator:

$$
\begin{equation*}
Z|x\rangle=(-1)^{x}|x\rangle, \quad x \in\{0,1\} \tag{3.9}
\end{equation*}
$$

Indeed, for a vector $|x\rangle, x \in\{0,1\}^{n}$ from the computational basis we have $\frac{I-Z_{i}}{2}|x\rangle=\frac{1}{2}\left(1-(-1)^{x_{i}}\right)|x\rangle=$ $x_{i}|x\rangle$.

## QUBO and Ising Hamiltonian for MaxCut

The simplest example is that of Maximum Cut. In this problem the goal is, given a graph $G=(V, E)$, to split the vertex set $V=V_{1} \cup\left(V \backslash V_{1}\right)$ in a way that the number of edges $(u, v)$ that are cut by the partition (i.e. the edges that have precisely one node is in $V_{1}$ ) is maximal. The QUBO formulation associates one variable per node:

$$
\begin{equation*}
\max _{x \in\{0,1\}^{n}} \sum_{(u, v) \in E}\left(x_{u}+x_{v}-2 x_{u} x_{v}\right) \tag{3.10}
\end{equation*}
$$

and $V_{1}=\left\{i: x_{i}=1\right\}$. The corresponding Ising Hamiltonian is

$$
\begin{equation*}
H=-\sum_{(u, v) \in E} \frac{1}{2}\left(I-Z_{u} Z_{v}\right) \tag{3.11}
\end{equation*}
$$

As expected, QUBO is NP-hard Pardalos and Jha, 1992. Nevertheless, in particular cases there exist polynomial-time algorithms. For example, a linear algorithm is presented in Barahona, 1986 for instances with the connectivity $E$ given by a series-parallel graph. We recall that series-parallel graphs are recursively defined by series and parallel compositions Eppstein, 1992. The base case is a graph with a single edge and two distinguished nodes called source and sink.

QUBO instances with about $\approx 100$ variables can be solved exactly by such algorithms as Branch \& Bound or Column Generation Mauri and Lorena, 2012. Numerous metaheuristics such that tabu search, greedy, genetic, and evolutionary approaches as well as local search methods were adapted to QUBO Kochenberger et al., 2014. The most rapid of them can tackle about $\approx 10^{4}$ variables. It was also observed that Simulated Annealing metaheuristic, which is close in spirit to quantum heuristics behaves efficiently compared to other methods Alkhamis et al., 1998.

### 3.2.2 Beyond QUBO

QUBO is a very powerful model for optimization problems. However, most natural formulations in combinatorial optimization go beyond QUBO. Indeed, the native formulations of problems of interest may have non-binary integer variables (ex. graph coloring), higher-degree interactions (ex. E3LIN2, Max-Q-XORSAT) or hard constraints that restrict the solution space.

Non-binary bounded integers can be represented by unary encoding Hadfield et al., 2019 that associated to an integer $x \in[0, \ldots, k-1]$ a sequence of bits $x_{b} \in\{0,1\}^{k}$ where the value 1 is placed at the position $k$. This encoding induces a new hard constraint on the solution space: it is possible to recover the corresponding integer from a bitstring $x_{b} \in\{0,1\}^{k}$ only if it satisfies the condition $\sum_{i=1}^{k} x_{b}^{i}=1$. Another option is to use qudits - quantum systems that by definition can be in $k$ different states. This approach goes beyond the traditional framework for quantum computing. Despite this, in Bravyi et al., 2022 qutrits were used to solve the Max-3-Cut problem.

In quantum computing, there are two approaches to address hard constraints. The first one is penalization, i.e. the modification of the QUBO objective function $f(x)$ in a way that makes unfeasible solutions highly suboptimal. For a given constraint that restricts the solution space to $S$ we define a penalty function $g(x):\{0,1\}^{n} \rightarrow \mathbb{R}^{+}$such that:

$$
g(x)= \begin{cases}0, & x \in S  \tag{3.12}\\ \geq 1, & x \notin S\end{cases}
$$

If $f(x)$ is the objective in the initial constrained formulation, the modified objective is $f_{S}(x)=$ $f(x)+P g(x)$ where $P$ is a huge number compared to the possible values of $f$.

Widely-used penalty functions are $g(x)=x_{1} x_{2}$ for the constraint $x_{1}+x_{2} \leq 1$ and $g(x)=$ $\left(\sum_{i=1}^{n} a_{i} x_{i}-b\right)^{2}$ for linear constraints $\sum_{i=1}^{n} a_{i} x_{i}=b$ Kochenberger et al., 2014. For instance, with these penalties it is possible to transform an integer linear program for graph coloring in a QUBO Kochenberger et al., 2005.

Example: The Maximum Independent Set on a graph $G=(V, E)$ is a subset of pairwise disconnected nodes of the biggest cardinality. It is naturally formulated as an integer linear program:

$$
\begin{align*}
\max & \sum_{i \in V} x_{i}  \tag{3.13}\\
& x_{i}+x_{j} \leq 1, \quad \forall(i, j) \in E  \tag{3.14}\\
& x_{i} \in\{0,1\}, \quad \forall i \in V \tag{3.15}
\end{align*}
$$

After applying the transformation tricks the problem becomes

$$
\begin{equation*}
\max \sum_{i \in V} x_{i}-P \sum_{(i, j) \in E} x_{i} x_{j} \tag{3.16}
\end{equation*}
$$

where it suffices to take $P=n$ Bomze et al., 1999. In Farhi et al., 2020 it was suggested to fix a relatively small $P=1$ and to extract independent sets from infeasible solutions with a specially-designed postprocessing.

The second technique for constraint management suggests replacing the transverse field $\sum_{v \in V} \sigma_{v} X_{v}$ used by quantum algorithms with another mixing term that preserves the evolution
in the feasible subspace Hen and Spedalieri, 2016, Hen and Sarandy, 2016, Hadfield et al., 2019. This approach can significantly reduce the amount of required pair interaction compared to the penalization, which is important for limited-connectivity hardware. Moreover, it is particularly promising for the constraints arising from the unary encoding as the corresponding feasibility condition is preserved by so-called XY-mixers Hadfield, 2018, Wang et al., 2020. We remind that XY-Hamiltonian on $k$ qubits is

$$
\begin{equation*}
H_{X Y}=\sum_{(i, j) \in E} X_{i} X_{j}+Y_{i} Y_{j} \tag{3.17}
\end{equation*}
$$

where the connectivity $E$ usually is either a chain or a complete graph.
Some hardware-specific constraint embedding strategies were studied into Vyskocil and Djidjev, 2019 for linear equality constraints. Linear inequalities can be transformed in equalities by introducing slack variables that, in turn, have to be represented with a binary expansion. A hybrid quantumclassical method that doesn't require slack variables was introduced in Yonaga et al., 2020. This method solves QUBOs as a part of the iterative procedure.

Finally, natural objective functions for some important problems such as E3LIN2 or Max-K-XORSAT involve higher-degree interactions. Using additional variables they can always be reduced to QUBO Anthony et al., 2017 but such reduction leads to a significant increase in the amount of needed resources. As the number of qubits is strictly limited on NISQ machines, the approaches that deal directly with higher-degree functions may be more practical Glos et al., 2022.

Therefore, in general settings quantum algorithms deal with any pseudo-boolean cost function $C:\{0,1\}^{n} \rightarrow \mathbb{R}$ encoded by some diagonal Hamiltonian:

$$
\begin{equation*}
\hat{C}: \hat{C}|x\rangle=C(x)|x\rangle \tag{3.18}
\end{equation*}
$$

The work Hadfield, 2021 provides multiple examples of Hamiltonians encoding boolean and real-valued functions. In the following we will assume that instances of optimization problems are given by Hamiltonian operators $C$ satisfying the condition 3.18.

### 3.3 Quantum Adiabatic Algorithm

Adiabatic Algorithm for combinatorial optimization was originally introduced in Farhi et al., 2000. In the Adiabatic Algorithm the quantum system evolves under the weighted sum of two Hamiltonians $H_{M}$ and $H_{C}$. For a problem with $n$ binary variables these are a transverse-field $H_{M}=$ $\sum_{i=1}^{n} X_{i}$ and a problem Hamiltonian $H_{C}: H_{C}|x\rangle=f(x)|x\rangle$ where $f(x)$ is the function that we want to optimize.

The quantum evolution of total duration $t_{f}$ is governed by the Hamiltonian $H(t)$ :

$$
\begin{equation*}
H(t)=(1-A(t)) H_{C}+A(t) H_{M} \tag{3.19}
\end{equation*}
$$

The Hamiltonian $H(t)$ with diagonal $H_{C}$ and $H_{M}$ being transverse field is stoquastic, i.e. its off-diagonal elements are non-positive.

We assume that the time-dependent Hamiltonian may be written in a form:

$$
\begin{equation*}
H(t)=\widetilde{H}\left(\frac{t}{t_{f}}\right)=\widetilde{H}(s) \tag{3.20}
\end{equation*}
$$

where $t_{f}$ is the total time of evolution.
A system quantum system initialized in the state $\left|\psi_{0}\right\rangle$ evolves according to the Schroedinger equation:

$$
\begin{align*}
i \frac{\partial}{\partial s}|\psi(s)\rangle & =t_{f} \tilde{H}(s)|\psi(s)\rangle  \tag{3.21}\\
|\psi(0)\rangle & =\left|\psi_{0}\right\rangle \tag{3.22}
\end{align*}
$$

where $\widetilde{H}(s)=(1-\tilde{A}(s)) H_{C}+\tilde{A}(s) H_{M}$ (for simplicity we took $\hbar=1$ ).
In traditional Adiabatic Algorithm the initial state is taken to be the ground state of $H_{M}$ : $\left|\psi_{0}\right\rangle=\frac{1}{\sqrt{2^{n}}} \sum_{x \in\{0,1\}^{n}}|x\rangle$, and the control function $\bar{A}(s)$ gradually changes from 1 at $s=0$ to 0 at $s=1$. Another version suggests to gradually turn off the transverse field:

$$
\begin{equation*}
\tilde{H}(s)=H_{C}+\Gamma(s) H_{M} \tag{3.23}
\end{equation*}
$$

where $\Gamma(s)$ decreases from a big value $G_{0}$ to 0 at $s=1$. Sometimes the adiabatic evolution under (3.23) is called quantum annealing that we distinguish from the heuristic with the same name presented in section 3.4 .

## Adiabatic theorem

According to the adiabatic theorem, if the evolution is slow enough the final state of the system is close to the ground state of $\tilde{H}(1)=H_{C}$ Jansen et al., 2007, Amin, 2009a.

There exists several formulation of the adiabatic theorem with different conditions on the Hamiltonian's smoothness. We refer to the survey Albash and Lidar, 2018 for an comprehensive review of different formulations.

The work Amin, 2009b gives a formulation that doesn't impose any specific conditions on $\tilde{H}(s)$. This formulation claims that the system remains in a ground state if the following condition is satisfied:

$$
\begin{equation*}
\frac{1}{t_{f}} \max _{s \in[0,1]} \frac{\left.\left|\left\langle\epsilon_{i}(s)\right| \partial_{s} \tilde{H}(s)\right| \epsilon_{j}(s)\right\rangle \mid}{\left|\epsilon_{i}(s)-\epsilon_{j}(s)\right|^{2}} \ll 1, \quad \forall i \neq j \tag{3.24}
\end{equation*}
$$

where $\left|\epsilon_{j}(s)\right\rangle$ is an eigenstate of $\tilde{H}(s)$ with the eigenvalue $\epsilon_{j}(s)$. The index $j$ corresponds to the relative order of the energy level in the spectrum: $\epsilon_{1}(s)<\epsilon_{1}(s)<\ldots$. The difference $\Delta(s)=\epsilon_{2}(s)-\epsilon_{1}(s)$ is often referred as spectral gap. A informal version of the adiabatic theorem tells that the total evolution time $t_{f}$ should be big with respect to the inverse of the gap square $\frac{1}{\Delta^{2}}$ in order to preserve the evolution in the instantaneous ground state.

## Universality of adiabatic evolution

It was proven that the adiabatic evolution with local non-stoquastic Hamiltonians is universal Aharonov et al., 2007. It means in particular that any state reachable by the gate-based quantum computer in time $T$ is the ground state of some local $H_{C}^{\prime}$ and the gap $\Delta(s)$ of the transverse-field Hamiltonian $H(s)$ is polynomial at any moment $s$.

However, as Hamiltonians involved in the Adiabatic Algorithm for optimization are stoquastic it is unlikely that the algorithm allows a universal quantum computation Albash and Lidar, 2018. From the complexity point of view, it was shown that for any stoquastic Hamiltonian $H_{C}$ such that the spectral gap of $\tilde{H}(s)$ is polynomial the ground state problem belongs to PostBPP Bravyi et al., 2006.

It is an open question if there are some instances $H_{C}$ such that the gap $\Delta(s)$ is polynomial while classical algorithms take exponential time to find the solution. As quantum computers are not believed to solve NP-complete problems in polynomial time, the Adiabatic Algorithm is usually tested for limited quantum advantage Albash and Lidar, 2018. It means that the algorithm is not challenged by the best possible classical algorithm or even by the best already known algorithm but, instead, it is compared to classical metaheuristics inspired by similar ideas.

### 3.3.1 Classical annealing

The Adiabatic Algorithm is often compared to two classical algorithms: Simulated Annealing and Simulated Quantum Annealing.

## Simulated Annealing

Simulated Annealing is an acknowledged metaheuristic inspired by the physical procedure of the heating and slow cooling of a metal to a uniform crystalline state Kirkpatrick et al., 1983. In Simulated Annealing the cost function $f(x)$ is identified with the energy of a statisticalmechanical system. The role of the transverse field $H_{M}$ is played by the kinetic energy that is proportional to the system's temperature $T(t)$. The evolution starts at high temperature $T_{0}$ and is slowly cooled to $T\left(t_{f}\right)=0$. Thermal fluctuations, just as quantum tunneling in the quantum annealing process, help the system to explore the entire energy landscape. In particular, they allow escaping from suboptimal local minima by changing the configuration to one with higher potential energy. As the temperature gets lower, such transitions become less likely and at the end the system is in the lowest-energy state.

Just as the adiabatic condition guarantees the convergence to the optimum for the Adiabatic Algorithm, the Simulated Annealing is promised to terminate in a ground state if the system is maintained in the thermal equilibrium at every moment. It was shown in Geman and Geman, 1984 that if the temperature decreases as

$$
\begin{equation*}
T(t) \propto \frac{n}{\log (\alpha t+1)} \tag{3.25}
\end{equation*}
$$

where $n$ is the amount of variables the thermal equilibrium condition is satisfied. A similar result was derived for the adiabatic condition Morita and Nishimori, 2008: the condition is satisfied if the transverse field decreases as

$$
\begin{equation*}
\Gamma(t) \propto \frac{1}{(\delta t)^{1 /(2 n-1)}} \tag{3.26}
\end{equation*}
$$

where the parameter $\delta$ depends on the already mentioned spectral gap $\Delta$.
The decrease of quantum fluctuations (3.26) obeys a power law which is asymptotically faster than the logarithmic decrease (3.25) for the Simulated Annealing. This inspires a certain optimism in quantum annealing potentially improving over Simulated Annealing Morita and Nishimori, 2008. The work Farhi et al., 2002 demonstrates instances where the Adiabatic Algorithm doesn't provide any improvement over the Simulated Annealing as well as instances showing limited quantum advantage.

## Simulated Quantum Annealing

The programmed implementations of simulated annealing imitate the physical cooling process. Likewise, Simulated Quantum Annealing (SQA) was designed to emulate the quantum evolution
that obeys the Schroedinger equation (3.22) on a classical computer Martoňák et al., 2002. The idea is realized by Quantum Monte-Carlo Simulation that modifies a set of configurations imitating the quantum tunneling phenomena. A peculiar result Morita and Nishimori, 2008 claiming for the SQA the same convergence condition 3.26 as for the actual quantum annealing revitalizes the debate of their relative power. A family of instances that are polynomially solved by the Adiabatic Algorithm but not with SQA were presented in Hastings and Freedman, 2013.

Regarding the relation between SQA and Simulated Annealing, SQA was shown to perform better on random spin-glass Martoňák et al., 2002, traveling salesman problem Martoňák et al., 2004. However, on random instances of 3-SAT Problem the vanilla SQA obtained worse results than the thermal annealing Battaglia et al., 2005. However, the modification of the annealing schedule can inverse the situation back in favor of SQA.

Several modifications were suggested that can improve the performance of the Adiabatic Algorithm Albash and Lidar, 2018, Crosson et al., 2014. Some of them suggest increasing the spectral gap by the addition of a local field in the middle of the evolution Crosson et al., 2014 or by modifying the annealing schedule $\tilde{A}(s)$ Morita and Nishimori, 2008. Others suggest abandoning adiabaticity either by the initialization in an excited state Crosson et al., 2014 or by allowing diabatic transitions.

### 3.4 Quantum Annealing

Quantum Annealing was applied to a lot of optimization problem such as Traveling Salesman and graph coloring Rieffel et al., 2014, image recognition Neven et al., 2008 and the ground state search for the Sherington-Kirkpatrik model Venturelli et al., 2015. In essence, the idea behind the Quantum Annealing is to use quantum fluctuations in order to explore the solution space.

In Quantum Annealing the system's state evolves according to the Schroedinger equation (3.22) under the Hamiltonian $H(t)=a(t) H_{C}+b(t) H_{M}$ where, just as in the Adiabatic Algorithm, $H_{C}$ encodes the cost function and $H_{M}$ is a transverse field $H_{M}=\sum_{i=1}^{n} X_{i}$. There are two traditional choices for the schedule coefficients $a(t)$ and $b(t)$. The first option is to chose the coefficients as in the Adiabatic Algorithm $\left(a(t)=\frac{t}{t_{f}}, b(t)=1-\frac{t}{t_{f}}\right)$. The other approach consists in fixing a constant $a(s)=C$ and to gradually "turn off" the transverse field : from large $b(0) \gg 0$ to $b\left(t_{f}\right)=0$.

## Analog quantum hardware

The annealing process requires a special analogue device called quantum annealer that allows continuous control of the Hamiltonian of the system. The most famous such devices are manufactured by DWave on superconducting qubits, while there are also realizations on Rydberg's atoms Pichler et al., 2018. Such physical systems allow only pair interactions so Quantum Annealing is used to solve problems only encoded by an Ising model.

Due to the decoherence effect of the physical hardware, a hard time limit is imposed on the annealing time $t_{f}<\tau$ where $\tau$ is the coherence time. The limited runtime implies that the adiabatic condition 3.24 is generally not satisfied. Therefore, the system can jump in an excited state and return a suboptimal result in the end. Quantum Annealing is an heuristic, just as widely used implementations of Simulated Annealing with fixed cooling times.

The potential profit comes from the quantum tunneling that can, in principle, help the system to escape from local minima (see figure 3.2). In Simulated Quantum Annealing that also exploits


Figure 3.2: Quantum tunneling versus thermal hopping through a potential barrier.
In Simulated Annealing the potential energy landscape $E=H_{C}$ (blue curve) corresponds to the cost function. A system trapped in a local minima $x_{l o}$ can escape by thermal hopping (violet path): first go to a higher-energy state and then drop in a global minima $x^{*}$. The probability of the transition (1) to a higher-energy state decreases with the temperature $T$.
In Quantum Annealing the energy at some moment $s$ (green curve) is an interpolation $H(s)=a(s) H_{C}+b(s) H_{M}$ of the problem Hamiltonian $H_{C}$ and a transverse field $H_{M}$. The coefficient $b(s)$ decreases to zero at the final moment $s=1$ and the final energy operator encodes exactly the cost function $H(1)=H_{C}$ (blue curve). In quantum mechanics a system with energy $E<E_{\max }(s)$ can directly tunnel through the energy barrier (red path). The probability of the tunneling depends on the width $\Delta$ and the hight $h(s)$ of the barrier. In Quantum Annealing the hight $h(s)$ is gradually increased by tuning $a(s)$ and $b(s)$, so the tunneling probability gets lower.
Simulated Quantum Annealing imitates the quantum tunneling inside a classical algorithm.
the tunneling effect, the quantum evolution under the Schroedinger equation is stochastically approximated by a Quantum Monte-Carlo chain. It was shown that quantum effects lead to a more efficient exploration of the search space than the thermal hopping Kadowaki and Nishimori, 1998, Martoňák et al., 2002, Martoňák et al., 2004.

### 3.4.1 Performance scaling for different annealing times

An interesting question is how the algorithm behaves when the runtime limit $t_{f}$ changes. For practical applications the success probability is a particularly relevant metric. The success probability is defined by the overlap $\left|\left\langle x^{*} \mid \psi_{t_{f}}\right\rangle\right|^{2}$ of the final state $\left|\psi\left(t_{f}\right)\right\rangle$ with the ground state $\left|x^{*}\right\rangle$. According to the adiabatic theorem at the infinite time limit $t_{f} \rightarrow \infty$ the overlap tends to 1 , but somewhat surprisingly it doesn't grow monotonically. A numerical simulation reported in Crosson et al., 2014 reveals that, given a range $t_{f} \in\left[T_{\text {min }}, T_{\text {max }}\right]$, there is an optimal $T^{*}$ that depends on the considered instance such that the success probability on $\left[T_{\min }, T_{\text {max }}\right]$ is maximal in $T^{*}$. For most instances of the Max-2-Sat problem considered in Crosson et al., 2014 the optimal $T^{*}$ is strictly inside the interval. Hence, the solution doesn't necessarily get better with
a finite increase of $t_{f}$.

## Lieb-Robinson bound

At the moment, there are only a few works with analytical results about Quantum Annealing in the constant time regime $t_{f}=O(1)$. The reason is a lack of tools highlighted in Braida et al., 2022. Generally, theoretical analyses are based on the Lieb-Robinson bound. This bound allows approximating the impact of the whole system on a local quantum observable by the action of its local neighborhood Tran et al., 2019]. More precisely, the bound applies to a local Hamiltonian $H=\sum_{i=1}^{k} H_{k}$ where each term acts on a small subset of qubits $S_{k}$. The Hamiltonian's connectivity is given by a graph $G=(V, E)$ where each node corresponds to a qubit and there is an edge between $u$ and $v$ if there is a term $H_{k}$ that acts on both $u \in S_{k}, v \in S_{k}$. We call the $r$-neighborhood of a term $H_{k}$ all qubits that are at distance at most $r$ from $S_{k}$ in the graph $G$. The Lieb-Robinson bound implies that for a short time $t_{f}$ we can approximate:

$$
\begin{equation*}
\left\langle\psi\left(t_{f}\right)\right| H_{k}\left|\psi\left(t_{f}\right)\right\rangle \approx\left\langle\psi_{r}\left(t_{f}\right)\right| H_{k}\left|\psi_{r}\left(t_{f}\right)\right\rangle \tag{3.27}
\end{equation*}
$$

where $\left|\psi_{r}(t)\right\rangle$ is a restriction of the system state $|\psi(t)\rangle$ on the $r$-neighborhood of $H_{k}$. The size $r$ of the relevant neighborhood depends on the time $t_{f}$. This dependence is sometimes called the light cone and is interpreted as a speed of information transmission in the quantum evolution. The bound can be applied for sparse connectivity $G$ and short runtime $t_{f}$ when the local dynamic can be simulated classically in a reasonable time.

In Braida et al., 2022 the Lieb-Robinson bound was used to establish an approximation ratio guarantee for the Quantum Annealing with very short $t_{f}$ on MaxCut on triangle-free 3regular graphs. However, the obtained guarantee is not expected to be tight. Besides, the work Moosavian et al., 2022 uses the bound to expose a set of instances for which the Quantum Annealing is not believed to show a quantum advantage in any constant time.

### 3.4.2 Experimental performance

In opposition to a limited number of theoretical results, a lot of efforts are invested in the experimental evaluation of Quantum Annealing. It is mostly due to the availability of large DWave machines that since recently can be accessed on a cloud. the question if they really implement quantum effects is still a subject of active debate with a number of negative Shin et al., 2014, Martin-Mayor and Hen, 2015, Rønnow et al., 2014 and positive pieces of evidence Boixo et al., 2014. Interestingly, the study Boixo et al., 2014 claims the presence of quantum effects by showing that DWave's output strongly correlates with those of SQA while having different results with the thermal annealing simulation.

Quantum annealers suffer from multiple constraints, some of them due to the construction choices and others - to unavoidable physical phenomena.

## Embedding

One of the potential obstructions for the applicability of DWave annealer is its sparse connectivity: a physical qubit in the chip interacts with its neighbors in a Chimera or Pegasus sparse layout Systems, Therefore, the Ising model for the cost function has to be embedded in the hardware architecture. The embedding associates to a logical qubit a chain of physical qubits that are encouraged to take the same value by ferromagnetic coupling $J_{i, j}<0$.

For any initial connectivity pattern the embedding can be done with at most quadratic overhead in the number of qubits Choi, 2008. An example for a fully-connected random Ising model is given in Venturelli et al., 2015. On the downside, the embedding reduces the performance of the heuristic by adding an extra energy scale required to impose the chain coupling $J_{i, j}<0$. It also significantly limits the size of the instances that can be treated by the machine as $O\left(n^{2}\right)$ qubits are needed to embed a fully connected model on $n$ variables. Hence it is somewhat more important to increase the hardware connectivity than to increase the number of qubits Rieffel et al., 2014.

In instances issued from combinatorial optimization, it is often the case that dense connectivity is a result of penalization of hard constraints. Therefore, other techniques to deal with constraints are of great interest. The approach from Hen and Spedalieri, 2016, Hen and Sarandy, 2016 suggests modifying the mixing term in order to keep the evolution in the feasible subspace. Such a modification consists in the initialization of the system in some feasible state and the evolution under new driver $H_{M}$. A way to impose a linear equality constraint on $n$ variables using $O(n)$ physical qubits from DWave graph is presented in Vyskocil and Djidjev, 2019. For some applications it may be profitable to interchange the embedding order: the work Bian et al., 2016 on Constraint Satisfaction Problem suggests to locally embed the problem's constraints and then couple qubits that correspond to the same variable.

Another obstacle for a good performance is the limited range and precision of control parameters $J_{i, j}$ and $h_{i}$ in the physical Ising model $H=\sum_{(u, v) \in E_{H}} J_{u, v} Z_{u} Z_{v}+\sum_{u \in V_{H}} h_{u} Z_{u}$ where $G_{H}=\left(V_{H}, E_{H}\right)$ is the hardware layout graph Venturelli et al., 2015. A way to intelligently deal with these restrictions is described in Bian et al., 2014.

The sparse connectivity and the limited control are in a sense due to the construction choice (that is usually justified by experimental constraints). These characteristics define a set of hardware-native instances that can directly be solved by the annealer. Complications arise when the instance of interest doesn't naturally map to these parameters and needs to be embedded.

## Hardware noise

In addition to construction limits, such physical phenomena as decoherence and noise impact the annealing performance on all instances including hardware-native ones. Indeed, it was claimed that a non-zero environment temperature may transform a quantum annealer in a simple thermal annealer Albash et al., 2017 or make it perform even worse Martin-Mayor and Hen, 2015.

Besides, the annealer's performance also suffers from random fields and random-bond fluctuations that are present in the chip. It was shown that in order to assure close-to-correct dynamics while growing the number of qubits a significant reduction of such effects is required [Zhu et al., 2016]. Another option to tackle this issue is to use error mitigation Kandala et al., 2019, Berg et al., 2022.

### 3.5 Quantum Approximate Optimization Algorithm (QAOA)

QAOA, originally introduced in Farhi et al., 2014a is a quantum-classical heuristic for combinatorial optimization. In a same manner that Quantum Annealing, QAOA receives in input a pseudo-boolean cost function $f:\{0,1\}^{n} \rightarrow \mathbb{R}$ encoded in a diagonal energy operator $C: C|x\rangle=f(x)|x\rangle$. It returns a low-energy state $x_{o}$ that corresponds to an approximate optimum of $f(x)$.

### 3.5.1 QAOA components

QAOA uses quantum computers to prepare a state made out of three principal components: an initial state $\left|\psi_{0}\right\rangle$ and two families of operators that don't commute called phase-separation and mixing.

## Initial state

The initial state $\left|\psi_{0}\right\rangle$ should be in the first place trivial to prepare. It is usually taken to be a uniform superposition:

$$
\begin{equation*}
\left|\psi_{0}\right\rangle=\frac{1}{2^{n}} \sum_{x \in\{0,1\}^{n}}|x\rangle \tag{3.28}
\end{equation*}
$$

Such choice of the initial state maintains the link between QAOA and the Adiabatic Algorithm. However, other options may be beneficial in distinct optimization contexts. For instance, if the problem of interest has hard constraints, a version of QAOA Hadfield et al., 2019 suggests to explicitly keep the evolution in a feasible subspace $S_{b} \subsetneq\{0,1\}^{n}$. In this approach, the system is initialized in some superposition over feasible solutions.

Another alternative is to warm start the optimization with a proper choice of the initial state. The idea behind this is to integrate the information used by classical approximation algorithms in order to concentrate the search efforts on promising solutions. For the MaxCut example, the classical information can be the approximate solution returned by the Goemans Williamson optimization algorithm Egger et al., 2021 or the information gained from the corresponding semidefinite program [Egger et al., 2021, Tate et al., 2020].

## Phase separation operator

Phase-separation operators carry the information about the objective function. The family contains the following parametrized diagonal unitaries:

$$
\begin{equation*}
U_{P}(\gamma)=e^{i \gamma C} \tag{3.29}
\end{equation*}
$$

where $C$ is the problem Hamiltonian and $\gamma$ is the parameter that serves as the rotation angle.

## Mixing operators

As phase-operators are diagonal, they alone can't change the probability distribution. Therefore, a mixing family is introduced. This family is responsible for the exploration of the solution space. In a traditional QAOA the mixing is defined as:

$$
\begin{equation*}
U_{M}(\beta)=e^{i \beta B} \tag{3.30}
\end{equation*}
$$

where $B=\sum_{i=1}^{n} X_{i}$ is the transverse field Hamiltonian.
If the solution space is restricted to $S_{b} \subsetneq\{0,1\}^{n}$ by hard constraints other mixing families can be appealing Hadfield et al., 2019. Crucially, a mixing family has to connect all feasible solutions: for each pair $(x, y) \in S_{b}$ there should exist $m \in \mathbb{N}$ such that the transition probability is positive:

$$
\begin{equation*}
\left.\left|\langle y| U_{M}\left(\beta_{m-1}\right) \ldots U_{M}\left(\beta_{0}\right)\right| x\right\rangle\left.\right|^{2}>0, \quad \forall x, y \in S_{b} \tag{3.31}
\end{equation*}
$$

In the Adiabatic Algorithm the evolution starts in the ground state $\left|\psi_{0}\right\rangle$ of the driver Hamiltonian $B$. Such initialization is a part of the conditions of the adiabatic theorem. In an attempt to preserve the correspondence between QAOA and the Adiabatic Algorithm the mixing family is modified in some warm-start approach Egger et al., 2021. A matrix $B^{w s}$ is chosen such that the warm-start state $\left|\psi_{0}^{w s}\right\rangle$ is the ground state of $B^{w s}$ and the mixing family is $U_{M}(\beta)=e^{i \beta B^{w s}}$.

QAOA applies phase separation and mixing in alternation:

$$
\begin{equation*}
\left|\psi_{p}(\boldsymbol{\beta}, \gamma)\right\rangle=U_{M}\left(\beta_{m-1}\right) U_{P}\left(\gamma_{m-1}\right) \ldots U_{M}\left(\beta_{0}\right) U_{P}\left(\gamma_{0}\right)\left|\psi_{0}\right\rangle \tag{3.32}
\end{equation*}
$$

where the number of layers $p$ called depth is chosen in advance. For $p \rightarrow \infty$ the QAOA circuit is simply the Trotterization of the adiabatic evolution Farhi et al., 2014a. Interestingly, the form (3.32) is universal, i.e. it is expressive enough to encode any desired quantum computation Lloyd, 2018.

At the end the system is measured in the computational basis. The solution $|x\rangle$ is returned with probability $\left|\left\langle x \mid \psi_{p}(\boldsymbol{\beta}, \gamma)\right\rangle\right|^{2}$.

## Classical loop

The classical part of QAOA serves to determine parameter values $\boldsymbol{\beta}^{*}, \boldsymbol{\gamma}^{*}$. Each choice of parameters correspond to a different states $\left|\psi_{p}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle$ and, in consequence, to a different probability distributions

$$
\begin{equation*}
\mathbb{P}_{\beta, \gamma}(x)=\left|\left\langle x \mid \psi_{p}(\boldsymbol{\beta}, \gamma)\right\rangle\right|^{2}, \quad x \in\{0,1\}^{n} \tag{3.33}
\end{equation*}
$$

QAOA output is a sample from the parameterized distribution 3.33). We score a set of parameter values $\boldsymbol{\beta}, \boldsymbol{\gamma}$ by the ability of the corresponding distribution $\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}(x)$ to return low-energy states.

At first sight, one may suggest to discriminate distributions by the probability $\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}\left(x^{*}\right)$ to output the optimal solution $x^{*}$. This, however, doesn't make sense as such metric requires to know the solution that we are looking for in advance. Instead, QAOA compares the distributions by the corresponding expectations of the objective function:

$$
\begin{equation*}
L(\boldsymbol{\beta}, \boldsymbol{\gamma})=\langle f\rangle_{\boldsymbol{\beta}, \boldsymbol{\gamma}}=\left\langle\psi_{p}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right| C\left|\psi_{p}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle \tag{3.34}
\end{equation*}
$$

## How many measurements we need?

We expect the algorithm to return a value that is at least as good as the expectation $L(\boldsymbol{\beta}, \boldsymbol{\gamma})$. The number of repetitions $m$ required to obtain a sample $\hat{X}$ such that $f(\hat{X})<L(\boldsymbol{\beta}, \boldsymbol{\gamma})+1$ can be estimated from the Chebyshev's inequality.

We consider the average energy $\bar{f}_{m}=\frac{1}{m} \sum_{i=1}^{m} f\left(X_{i}\right)$ of $m$ independent samples $X_{1}, \ldots, X_{m}$ from $\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}$. By linearity,

$$
\begin{equation*}
E\left[\bar{f}_{m}\right]=\frac{1}{m} \sum_{i=1}^{m} E\left[f\left(X_{i}\right)\right]=L(\boldsymbol{\beta}, \gamma) \tag{3.35}
\end{equation*}
$$

The Chebyshev's inequality states:

$$
\begin{equation*}
\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}\left[\left|\bar{f}_{m}-L(\boldsymbol{\beta}, \boldsymbol{\gamma})\right| \geq 1\right] \leq \bar{\sigma}^{2} \tag{3.36}
\end{equation*}
$$

where $\bar{\sigma}^{2}$ is the variance of $\bar{f}_{m}$. As all $X_{i}$ are independent $\bar{\sigma}^{2}=\frac{\sigma^{2}}{m}$ where $\sigma^{2}$ is the variance of the mean energy in the QAOA state: $\sigma^{2}=\left\langle f^{2}\right\rangle_{\boldsymbol{\beta}, \boldsymbol{\gamma}}-\langle f\rangle_{\boldsymbol{\beta}, \boldsymbol{\gamma}}^{2}$.

After $m$ repetitions with probability at least $1-\frac{\sigma^{2}}{m}$ the average $\bar{f}_{m}<L(\boldsymbol{\beta}, \boldsymbol{\gamma})+1$. If this inequality is true, there is at least one sample $X_{k}$ such that $f\left(X_{k}\right) \leq L(\boldsymbol{\beta}, \boldsymbol{\gamma})+1$. This sample is the solution returned by QAOA.

### 3.5.2 How to compute the mean energy?

Accordingly to the previous observations, the mean energy $L(\boldsymbol{\beta}, \boldsymbol{\gamma})$ is indeed a reasonable choice for the loss function in the classical parameter optimization routine. Yet it is not clear how to compute the value $L(\boldsymbol{\beta}, \boldsymbol{\gamma})$.

The work Bravyi et al., 2021 speculates about the intrinsic difficulty to evaluate the expectation of a general observable $O=O_{1} \otimes \cdots \otimes O_{n}$ in the output distribution of a shallow quantum circuit. For instance, for a constant-depth circuit even the approximate value is \#P-hard to compute. In a special case where each $O_{i}$ is close to identity the problem admits a polynomial approximation scheme. This setting doesn't apply to QAOA where the energy operator $C$ is a sum of operators in the form $C_{k}=\prod_{i \in S_{k}} Z_{i}$ Bravyi et al., 2021.

On the other side, the QAOA circuit has a specific form that can simplify the computation of mean values. Indeed, for the MaxCut problem a formula for the mean value in QAOA with $p=1$ was presented in Wang et al., 2017. We generalized the result to the case of weighted MaxCut in Dalyac et al., 2021 and, using the ZX-calculus, to the general Ising model (see chapter 8).

Claim 3.5.1 For the weighted MaxCut problem encoded in the Ising model $C=\sum_{1 \leq u<v \leq n} e_{u, v} Z_{u} Z_{v}$ the mean value of $Z_{u} Z_{v}$ in the parameterized state $|\psi(\beta, \gamma)\rangle=U_{M}(\beta) U_{C}(\gamma)\left|\psi_{0}\right\rangle$ of $Q \bar{A} O A$ with depth 1 is:

$$
\begin{align*}
\langle\psi(\beta, \gamma)| Z_{u} Z_{v}|\psi(\beta, \gamma)\rangle & =\frac{-\sin (4 \beta) \sin \left(\gamma e_{u, v}\right)}{2}\left[\prod_{w \neq u, v} \cos \left(\gamma e_{u, w}\right)+\prod_{w \neq u, v} \cos \left(\gamma e_{v, w}\right)\right]+ \\
& +\frac{\sin ^{2}(2 \beta)}{2}\left[\prod_{w \neq u, v} \cos \left(\gamma\left(e_{u, w}-e_{v, w}\right)\right)-\prod_{w \neq u, v} \cos \left(\gamma\left(e_{u, w}+e_{v, w}\right)\right)\right] \tag{3.37}
\end{align*}
$$

We remark that the same result was obtained in an independent work Bravyi et al., 2020.
The works Basso et al., 2021, Farhi et al., 2022 suggest iterative procedures that compute mean energies. The former considers MaxCut problem on large-girth (no short cycles) regular graphs and the latter investigate the Sherrington-Kirkpatrik model in infinite size limit. Both procedures have the complexity that grows exponentially with $p$, although the complexity scalings are independent from the number of variables $n$. In Hadfield et al., 2021 the mean values for different $p$ are written as power series.

## Stochastic approximation

In all other cases one can access the expectation value $L(\boldsymbol{\beta}, \gamma)$ using stochastic approximation:

$$
\begin{equation*}
L(\boldsymbol{\beta}, \gamma) \approx \hat{L}_{N}^{\boldsymbol{\beta}, \boldsymbol{\gamma}}=\frac{1}{N} \sum_{k=1}^{N} f\left(x_{k}\right) \tag{3.38}
\end{equation*}
$$

If the stochastic approximation is used, the classical routine that minimizes $L(\boldsymbol{\beta}, \boldsymbol{\gamma})$ over possible values of $\boldsymbol{\beta}$ and $\boldsymbol{\gamma}$ makes calls to the quantum computer to get samples $X_{1}, \ldots, X_{N}$. Other
stochastic aggregation function such as Conditional Value-at-Risk can be employed to score parameter values Braine et al., 2021 with a potential improvement of the convergence speed Barkoutsos et al., 2020.

The interaction between quantum and classical parts of the algorithm is demonstrated on the figure 3.3. To resume, QAOA use quantum computer in two separate contexts. The classical procedure that search optimal parameters makes calls to the quantum processor when the loss function $L(\boldsymbol{\beta}, \boldsymbol{\gamma})$ is evaluated stochastically. Once optimal parameters $\boldsymbol{\beta}^{*}, \boldsymbol{\gamma}^{*}$ are determined, the QAOA state is measured several times and the best output is returned as solution.


Figure 3.3: The interaction between classical and quantum parts of QAOA. Green (red) rectangles correspond to classical (quantum) parts of the algorithm. Mean value evaluation depends on the implementation and may either make calls to a quantum computer or, when $p=1$ and in some special cases be done classically.

The optimal choice of parameter lead to the following guarantee on the approximation ratio of QAOA:

$$
\begin{equation*}
\alpha=\frac{\min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L(\boldsymbol{\beta}, \boldsymbol{\gamma})}{f\left(x^{*}\right)} \tag{3.39}
\end{equation*}
$$

where $x^{*}$ is the optimal solution of the problem. In practice, the parameter optimization search is difficult because of the highly irregular landscape with multiple local minimas. If the loss function $O(\boldsymbol{\beta}, \boldsymbol{\gamma})$ is evaluated stochastically, the optimization is hampered by the related stochastic noise.

For the reasons mentioned above the optimization procedure returns values $\boldsymbol{\beta}^{*} \gamma^{*}$ such that $L\left(\boldsymbol{\beta}^{*}, \gamma^{*}\right) \geq \min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L(\boldsymbol{\beta}, \boldsymbol{\gamma})$. The contribution of the optimization error can be explicitly separated in the performance metric:

$$
\begin{align*}
\epsilon & =\frac{L\left(\boldsymbol{\beta}^{*}, \boldsymbol{\gamma}^{*}\right)-f\left(x^{*}\right)}{f_{\max }-f\left(x^{*}\right)}  \tag{3.40}\\
& =\underbrace{\frac{L\left(\boldsymbol{\beta}^{*}, \boldsymbol{\gamma}^{*}\right)-\min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L(\boldsymbol{\beta}, \boldsymbol{\gamma})}{f_{\max }-f\left(x^{*}\right)}}_{E_{M}}+\underbrace{\frac{\min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L(\boldsymbol{\beta}, \boldsymbol{\gamma})-f\left(x^{*}\right)}{f_{\max }-f\left(x^{*}\right)}}_{E_{O}} \tag{3.41}
\end{align*}
$$

In the expression 3.41 the errors $E_{M}$ and $E_{O}$ are due to the model mismatch and to the optimization error respectively.

### 3.5.3 Performance guarantees for QAOA

For the big number of layers $p$ QAOA circuits are able to reproduce the Trotterized adiabatic evolution. Hence, in the limit $p \rightarrow \infty$ the algorithm achieves the exact optimum. Moreover, contrary to Quantum Annealing, the increase in quality of the solution is monotone. This fact is a direct consequence of $U_{C}(0)=U_{M}(0)=I$ :

$$
\begin{align*}
\min \left\langle\psi\left(\beta_{1}, \ldots, \beta_{\boldsymbol{p}-\mathbf{1}}\right)\right| C\left|\psi\left(\beta_{0}, \ldots, \beta_{\boldsymbol{p}-\mathbf{1}}\right)\right\rangle & =\min \left\langle\psi\left(\beta_{1}, \ldots, \beta_{\boldsymbol{p}-\mathbf{1}}, \mathbf{0}\right)\right| C\left|\psi\left(\beta_{1}, \ldots, \beta_{\boldsymbol{p}-\mathbf{1}}, \mathbf{0}\right)\right\rangle \\
& \geq \min \left\langle\psi\left(\beta_{1}, \ldots, \beta_{\boldsymbol{p}}\right)\right| C\left|\psi\left(\beta_{1}, \ldots, \beta_{\boldsymbol{p}}\right)\right\rangle \tag{3.42}
\end{align*}
$$

At low depths $p$ the analytical analysis allowed to establish some results about the approximation ratio of QAOA. The first result that caused a lot of excitement was the approximation ratio of $\left(\frac{1}{2}+\frac{1}{22 D^{3 / 4}}\right)$ achieved by QAOA with depth $p=1$ on E3LIN2 Farhi et al., 2014b. At the moment the best known classical guarantee was $\left(\frac{1}{2}+\frac{C}{D}\right)$ so $\mathrm{QAOA}_{1}$ demonstrated quantum speedup. Since then a better classical algorithm was found with $\alpha=\frac{1}{2}+\frac{C}{D^{1 / 2}}$ Barak et al., 2015]. This ratio matches the asymptotic scaling of the inapproximability bound Trevisan, 2001.

## MaxCut

Most of the analytical results were derived for QAOA applied to MaxCut. In MaxCut one aims to split the vertices of a graph $G=(V, E)$ on two sets $V=V_{1} \cup V_{2}$ such that the number of crossing edges $\left|\left\{(u, v) \in E: u \in V_{1}, v \in V_{2}\right\}\right|$ is maximized.

The computation of the approximation ratio with formulas (3.39) and (3.41) requires knowing the exact optimum that is generally difficult to compute. Hence the performance is usually reported in terms of the cut fraction $c$, i.e. the ratio of the size of QAOA cut to the total number of edges $c=\frac{L\left(\boldsymbol{\beta}^{*}, \boldsymbol{\gamma}^{*}\right)}{|E|}$. The original paper Farhi et al., 2014a demonstrated that $\mathrm{QAOA}_{1}$ achieves the cut fraction of at least 0.6924 on 3 -regular graph. On triangle-free $D$-regular graph $\mathrm{QAOA}_{1}$ obtains $c \geq \frac{1}{2}+\frac{0.3032}{\sqrt{D}}$ Hadfield, 2018. The work Wurtz and Love, 2021 provides the cut fractions for $p=2$ and $p=3$ for 3 -regular graphs with no cycles of length 5 and 7 respectively.

Most remarkably, the iterative procedure from Basso et al., 2022 permitted to show that QAOA for $p \geq 11$ outperforms the Goemans-Williamson algorithm on D-regular graphs with girth $l>2 p+1$. The Goemans-Williamson algorithm is particularly important as its performance matches the inapproximability bound for general graphs (that are not necessarily large-girth or regular). We remark that for the specific instances of large-girth regular graphs a polynomial approximation scheme exists Alaoui et al., 2021, albeit its guarantee relies on a still unproven conjecture.

The above results hold for the ideal quantum computer. When a real hardware is used, additional obstructions may significantly degrade the performance.

### 3.5.4 QAOA circuit

In physical implementation the evolution (3.32) has to be compiled into a circuit composed out of elementary gates. The set of elementary gates contains one-qubit and two-qubit gates that are supported by the hardware. Due to the limited connectivity, only a restricted subset of qubits can directly interact - which requires the addition of SWAP gates for two-qubit interactions that are not connected in the native layout.

We illustrate the compilation on a simple optimization problem encoded in the Hamiltonian:

$$
\begin{equation*}
C=\sum_{(u, v) \in E} Z_{u} Z_{v} \tag{3.43}
\end{equation*}
$$

where the interaction set $E$ corresponds to the graph $G=(V, E)$ shown on Figure 3.4. Up to a irrelevant additive constant the Hamiltonian (3.43) encodes the MaxCut problem on the graph $G$.


Figure 3.4: Example of the optimization problem.
As the Pauli-Z operators are diagonal the summands of (3.43) mutually commute, just as $X_{i}$ operators acting on different qubits. Hence we can write:

$$
\begin{align*}
& U_{C}(\gamma)=e^{i \gamma \sum_{(u, v) \in E} Z_{u} Z_{v}}=\prod_{(u, v) \in E} e^{i \gamma Z_{u} Z_{v}}  \tag{3.44}\\
& U_{M}(\beta)=e^{i \beta \sum_{i=1}^{n} X_{i}}=\prod_{i=1}^{n} e^{i \beta X_{i}} \tag{3.45}
\end{align*}
$$

If we assume that elementary gate set includes the Hadamard gate $H$, CNOT and the $Z$ rotation $R_{z}(\theta)=\left(\begin{array}{cc}e^{-i \frac{\theta}{2}} & 0 \\ 0 & e^{i \frac{\theta}{2}}\end{array}\right)$ for arbitrary angle $\gamma \in[0,2 \pi]$ we can compile $e^{i \beta X_{u}}$ to the circuit:

$$
\begin{equation*}
q_{u}:-H=R_{z}(2 \gamma)-H \tag{3.46}
\end{equation*}
$$

and $e^{i \gamma Z_{u} Z_{v}}$ to


Finally, the circuit for one QAOA layer $U_{M}(\beta) U_{C}(\gamma)$ for the Hamiltonian (3.43) is:


An optimized compilation strategy for the QAOA circuit on hardware with limited connectivity is presented in Lotshaw et al., 2022, Crooks, 2018.

In general, vanilla QAOA admits a relatively intuitive circuit representation. However, modifications of the mixing operator that are used to impose the hard constraints substantially increase the difficulty of the task. Indeed, even the simplest $X Y$ version requires the compilation of $U_{M}(\beta)=e^{i \beta B}$ where

$$
\begin{equation*}
B=\sum_{i=1}^{k} X_{i} X_{i+1}+Y_{i} Y_{i+1} \tag{3.49}
\end{equation*}
$$

Terms in the sum don't commute, so there is no simple way to decompose $U_{M}(\beta)$ on elementary gates. Although some compilation strategies were proposed Hadfield, 2018, this obstacle still seriously constrains the capacity of both numerical and physical experimentations.

### 3.5.5 Hardware implementations

QAOA was implemented on Google Sycamore superconducting hardware Harrigan et al., 2021 for instances with up to 23 variables. Experiments using different platforms were performed on 2qubit photonic computer Qiang et al., 2018 and on 40 trapped-ion qubits Pagano et al., 2020.

It was observed that limited connectivity together with noise dramatically impacts the performance of QAOA Harrigan et al., 2021. The question if the quantum advantage is possible at all in the presence of noise is still a matter of debate.

On the fundamental side, the work França and García-Patrón, 2021] shows that for the current noise level the output distribution is close to a thermal state with $T>T_{0}$. As a thermal state for $T>T_{0}$ ( $T_{0}$ depends on the instance) can be efficiently simulated on a classical computer the noise level should be significantly reduced in order to leave room for any quantum advantage.

The work Guerreschi and Matsuura, 2019 reports an attempt of a QAOA versus classical exact solvers comparative study that takes into account as much realistic conditions as possible. It was observed that, even if asymptotically QAOA (with approximate solution) is faster than exponential exact algorithms, it may require hundreds or thousands of qubits in order to get a runtime improvement. The stochastic noise in the evaluation of the loss function also negatively affects the overall runtime scaling Lotshaw et al., 2022.

The situation is slightly more optimistic for hardware-native instances i.e. instances whose connectivity pattern reproduces the qubit layout of the platform. Both works Harrigan et al., 2021, Pagano et al., 2020 observed that for these instances performance remains constant with the number of variables $n$ while it decreases for instances that require embedding. This is consistent with the intuition: embedding strategies increase the depth of the circuit and longer circuits are more impacted by noise.

## Classical simulation of QAOA

Even if some results on the real quantum hardware start to appear, most of the numerical results are obtained using simulations. As the full-scale simulation of the circuit on $n$ qubits has the complexity $2^{2 n}$ the experiments were done for only limited problem sizes.

From the fundamental side, even at the lowest depth $p=1$ QAOA exhibits quantum supremacy Farhi and Harrow, 2016. More precisely, it was demonstrated that if one could efficiently sample from the distribution $\mathbb{P}_{\beta_{1}, \gamma_{1}}(x)$ the polynomial hierarchy would collapse. As this is unlikely, it is believed that the output of QAOA can't be efficiently simulated. However, is some specific conditions of small angles en efficient simulation is possible Hadfield et al., 2021.

### 3.5.6 Locality in QAOA

One of the principal features of bounded depth QAOA is its locality. Roughly speaking, in a local optimization algorithm whenever it is quantum or classical a variable can be explicitly impacted only by its neighbors Hastings, 2019. We recall that the similar property for the Quantum Annealing follows from the Lieb-Robinson bound.

More precisely, a Hamiltonian $H=\sum_{i=1}^{N} H_{i}$ is called $k$-local if each term $H_{i}$ acts non-trivially on at most $k$ qubits $S_{k}\left(\left|S_{k}\right| \leq k\right)$ Albash and Lidar, 2018. Two variables $u$ and $v$ are called neighbors if they appear together in at least one $S_{k}$. The Ising cost operator

$$
\begin{equation*}
C=\sum_{u \in V} h_{u} Z_{u}+\sum_{(u, v) \in E} J_{u, v} Z_{u} Z_{v} \tag{3.50}
\end{equation*}
$$

is 2 -local and the neighborhood relationship is defined by the edges of the connectivity graph $G=(V, E)$. The corresponding phase separation is a product of $e^{i \gamma h_{u} Z_{u}}$ and $e^{i \gamma J_{u, v} Z_{u} Z_{v}}$ and the only 2 -qubit gates that appear in the QAOA circuit act on neighboring pairs of qubits.

This fact has some curious consequences.
Let the notation $\langle O\rangle_{p}$ stand for the expectation of an observable $O$ in some QAOA state of depth $p$ for arbitrary fixed $\boldsymbol{\beta}, \boldsymbol{\gamma} \in \mathbb{R}^{p}$. Locality implies that the mean values $\left\langle Z_{u}\right\rangle_{p}$ and $\left\langle Z_{u} Z_{v}\right\rangle_{p}$ depend only on the qubits that are at distance at most $p$ from the nodes $u$ or $v$ in the graph $G=(V, E)$. Indeed, in the Heisenberg representation the expectation of an observable $O$ in a state $|\psi\rangle=U|\psi\rangle$ is semantically equivalent to the mean value of the conjugated observable $\hat{O}=U^{\dagger} O U$ in the state $\left|\psi_{0}\right\rangle$ :

$$
\begin{equation*}
\langle\psi| O|\psi\rangle=\left\langle\psi_{0}\right| U^{\dagger} O U\left|\psi_{0}\right\rangle=\left\langle\psi_{0}\right| \hat{O}\left|\psi_{0}\right\rangle \tag{3.51}
\end{equation*}
$$

## Support of an observable

We call support of an operator $O$ a set of qubits on which $O$ acts non-trivially. The expectation of the observable $O$ with support $Q$ in a tensor product state $|\psi\rangle=\left|\psi_{1}\right\rangle \otimes \cdots \otimes\left|\psi_{n}\right\rangle$ is independent on states $\left|\psi_{i}\right\rangle, i \notin Q$ that are not in the support:

$$
\begin{align*}
\left\langle\psi_{n}\right| \otimes \cdots \otimes\left\langle\psi_{1}\right| O\left|\psi_{1}\right\rangle \otimes \cdots \otimes\left|\psi_{n}\right\rangle & =\left[\prod_{i \notin Q}\left\langle\psi_{i} \mid \psi_{i}\right\rangle\right] \times\left(\otimes_{i \in Q}\left\langle\psi_{i}\right|\right) O\left(\otimes_{i \in Q}\left|\psi_{i}\right\rangle\right) \\
& =\left(\otimes_{i \in Q}\left\langle\psi_{i}\right|\right) O\left(\otimes_{i \in Q}\left|\psi_{i}\right\rangle\right. \tag{3.52}
\end{align*}
$$

An observable $O$ with support $Q \subset V$ commutes with $e^{i \beta X_{u}}$ for all $u \notin Q$ and with $e^{i \gamma Z_{u} Z_{v}}$ for $u, v \notin Q$. Therefore, the mixing operator doesn't change the support of an observable. When one iteratively conjugates the observable $O$ with QAOA layers:

$$
\begin{equation*}
O^{i}=U_{C}^{\dagger}\left(\gamma_{i}\right) U_{M}^{\dagger}\left(\beta_{i}\right) O^{i-1} U_{M}\left(\beta_{i}\right) U_{C}\left(\gamma_{i}\right) \tag{3.53}
\end{equation*}
$$

the support is extended from $Q^{i-1}$ to $Q^{i}=Q^{i-1} \cup\left\{v \in V \mid \exists u \in Q^{i-1}:(u, v) \in E\right\}$. After $p$ conjugations the support $Q^{p}$ for $Z_{u} Z_{v}$ contains only qubits at distance at most $p$ from $u$ or $v$. With a standard initial state $\left|\psi_{0}\right\rangle=|+\rangle \otimes \cdots \otimes|+\rangle$ the expectation is:

$$
\begin{equation*}
\left\langle Z_{u} Z_{v}\right\rangle_{p}=\left(\otimes_{i \in Q^{p}}\langle+|\right) U_{Q^{p}}^{\dagger}(\boldsymbol{\beta}, \boldsymbol{\gamma}) Z_{u} Z_{v} U_{Q^{p}}(\boldsymbol{\beta}, \boldsymbol{\gamma})\left(\otimes_{i \in Q^{p}}|+\rangle\right) \tag{3.54}
\end{equation*}
$$

where $U_{Q^{p}}(\boldsymbol{\beta}, \boldsymbol{\gamma})$ is the restriction of QAOA on the qubits in $Q^{p}$ (i.e. gates that act only on qubits from $Q^{p}$ ).


Figure 3.5: The support of the observable $Z_{u} Z_{v}$ (blue link) for $p=1$ (red nodes) and $p=2$ (red and green nodes).

## The role of locality in QAOA

Locality explains the loss function concentration phenomenon reported in Brandao et al., 2018. The concentration means that for different instances of MaxCut generated according to the same distribution (for example $D$-regular graphs of different sizes $n_{1}, \ldots, n_{M}$ ) the loss function values $L_{1}(\boldsymbol{\beta}, \boldsymbol{\gamma}) \ldots L_{M}(\boldsymbol{\beta}, \boldsymbol{\gamma})$ for fixed $\boldsymbol{\beta}, \boldsymbol{\gamma}$ are approximately the same across instances.

Most of the results limiting the performance of QAOA also follows from the locality. The work Hastings, 2019 shows the cases when local classical algorithms outperform single-layer QAOA. It also conjectures that a bounded-depth QAOA can't improve over global classical methods on E3LIN2. Works Farhi et al., 2020 and Bravyi et al., 2020 use locality to bound the approximation ratio of constant-depth QAOA on Maximum Independent Set and MaxCut respectively. As stated in Farhi et al., 2020 QAOA with depth $p<\Omega(\log n)$ can't find an independent set $S$ with more than $0.854\left|S_{\text {max }}\right|$ nodes on a typical $d$-regular graphs. Locality together with the property of MaxCut cost function called $\mathbb{Z}_{2}$-symmetry was used in Bravyi et al., 2020 to construct a worst-case example for which QAOA with $p<\Omega(\log n)$ can't beat the approximation ratio of the Goemans-Williamsion algorithm.

## Approximation ratio of QAOA on Ising chain

We use locality to prove the simple fact that the approximation ratio of QAOA of depth $p$ on the ring of disagrees is bounded by $\alpha<\frac{2 p+2}{2 p+3}$. This fact was proven in Mbeng et al., 2019b using Jordan-Winger transformation and in [Mbeng et al., 2019a using boundary conditions. An independent work Wurtz and Love, 2021 provides a demonstration founded on ideas that are very close to ours.

The ring of disagree problem's cost function is:

$$
\begin{equation*}
C_{r}=\sum_{i=1}^{n} Z_{i} Z_{i+1} \tag{3.55}
\end{equation*}
$$

where we take $Z_{n+1} \equiv Z_{0}$. The Hamiltonian (3.55) encodes (up to some constant factor) the MaxCut problem on an Ising chain.

Claim 3.5.2 For the ring of disagree problem $Q A O A$ with depth $p$ achieves the approximation ratio at most $\alpha=\frac{2 p+2}{2 p+3}$.

## Proof:

For simplicity, we take $n=(2 p+3) \times m$ where $m \geq 1$ is even.
Ring of disagrees is an anti-ferromagnetic Ising model on a layout $G_{r}=\left(V_{r}, E_{r}\right)$ that is shown on the figure 3.7a. Crucially, after $p$ layers of QAOA the support of an observable $Z_{i} Z_{i+1}$
is $Q=\{u \mid u \in[i-p, \ldots i+p+1]\}$ (see Figure3.6). After the index change we can set $Q=\{u \mid u \in$ $[1, \ldots, 2 p+1]\}$. The restriction of the operator $C_{r}$ on the support $Q$ is $C_{r}^{Q}=\sum_{u=1}^{2 p} Z_{u} Z_{u+1}$.


Figure 3.6: Support of the operator $Z_{i} Z_{i+1}$ after $p$ layers of QAOA on the ring of disagree problem

Let's arbitrary fix some $\boldsymbol{\beta}, \boldsymbol{\gamma} \in \mathbb{R}^{p}$. Due to the locality property (3.54) the expectation $\left\langle\psi_{p}^{r}(\boldsymbol{\beta}, \gamma)\right| Z_{i} Z_{i+1}\left|\psi_{p}^{r}(\boldsymbol{\beta}, \gamma)\right\rangle$ in QAOA- $p$ state $\left|\psi_{p}^{r}(\boldsymbol{\beta}, \gamma)\right\rangle$ is the same as the expectation of $Z_{p} Z_{p+1}$ in the state

$$
\begin{equation*}
\left|\phi_{Q}(\boldsymbol{\beta}, \gamma)\right\rangle=e^{i \beta_{p} \sum_{q=1}^{2 p+1} X_{q}} \ldots e^{i \beta_{1} \sum_{q=1}^{2 p+1} X_{q}} e^{-i \gamma_{1} \sum_{q=1}^{2 p} Z_{q} Z_{q+1}}|+\rangle_{1} \ldots|+\rangle_{p} \tag{3.56}
\end{equation*}
$$

Thus the energy expectation is:

$$
\begin{equation*}
L_{r}(\boldsymbol{\beta}, \boldsymbol{\gamma})=\left\langle\psi_{p}^{r}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right| C_{r}\left|\psi_{p}^{r}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle=n \times\left\langle\phi_{Q}(\boldsymbol{\beta}, \gamma)\right| Z_{p} Z_{p+1}\left|\phi_{Q}(\boldsymbol{\beta}, \gamma)\right\rangle \tag{3.57}
\end{equation*}
$$

We consider another instance $C_{d}=\sum_{(u, v) \in E_{d}} Z_{u} Z_{v}$ where the pair interactions $E_{d}$ correspond to edges of the graph $G_{d}=\left(V_{d}, E_{d}\right)$ made out of $m$ disconnected cycles each of length $2 p+3$ The graph $G_{d}$ is shown on the figure 3.7 b . We remark that $\left|V_{r}\right|=\left|V_{d}\right|=n$ and $\left|E_{r}\right|=\left|E_{d}\right|=n$.

(a) The ring of disagree graph $G_{r}=$ $\left(V_{r}, E_{r}\right)$. It has $n=(2 p+3) \times m$ nodes. The Hamiltonian (3.55) encodes the MaxCut problem on this graph.

(b) A graph $G_{d}=\left(V_{d}, E_{d}\right)$ with the same $p$-local structure as $G_{r}$. Each of $m$ components has $2 p+3$ nodes.

As the each small cycle in $G_{d}$ has length $l>2 p+1$ the support of the observable $Z_{u} Z_{v}$, $(u, v) \in E_{d}$ conjugated by $p$ layers of QAOA for $C_{d}$ is exactly as on 3.6.

Using the same reasoning as before we obtain:

$$
\begin{equation*}
L_{d}(\boldsymbol{\beta}, \boldsymbol{\gamma})=\left\langle\psi_{p}^{d}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right| C_{d}\left|\psi_{p}^{d}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle=n \times\left\langle\phi_{Q}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right| Z_{p} Z_{p+1}\left|\phi_{Q}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle \tag{3.58}
\end{equation*}
$$

As $\boldsymbol{\beta}, \boldsymbol{\gamma} \in \mathbb{R}^{p}$ were chosen arbitrary the loss functions for two instances are equal: $L_{r} \equiv L_{d}$. As the loss function $L_{d}$ is the energy expectation $\left\langle C_{d}\right\rangle$ in some distribution we have:

$$
\begin{equation*}
\min _{x} C_{d}(x) \leq \min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L_{d}(\boldsymbol{\beta}, \boldsymbol{\gamma}) \equiv \min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L_{r}(\boldsymbol{\beta}, \boldsymbol{\gamma}) \tag{3.59}
\end{equation*}
$$

The ground state problem for $C=\sum_{(u, v) \in E} Z_{u} Z_{v}$ is equivalent to the MaxCut problem $\min _{x \in\{0,1\}} \operatorname{cut}(x)$ on the layout graph $G=(V, E)$. Indeed, there is a one-two-one correspondence between solutions of two problems. The maximum cut on $G_{r}$ is obtained in the partition $x^{*}=$ $[0,1,0,1 \ldots 0,1]$ so the ground state energy for $C_{r}$ is $C_{r}\left(x^{*}\right)=\left\langle x^{*}\right| C_{r}\left|x^{*}\right\rangle=-n$. As graph $G_{d}$ has $m$ odd cycles its maximum cut is at most $c_{o p t}=n-m$ achieved by variable assignment $x^{*}=[0,1,0,1 \ldots 0,1]$. Therefore, the ground energy of $C_{d}$ is $C_{d}\left(x^{*}\right)=\left\langle x^{*}\right| C_{d}\left|x^{*}\right\rangle=-n+m$.

Finally, on the ring of disagrees the QAOA with depth $p$ has the approximation ratio:

$$
\begin{align*}
\alpha_{p} & =\frac{\min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L_{r}(\boldsymbol{\beta}, \boldsymbol{\gamma})}{\min _{x} C_{r}(x)}=\frac{\min _{\boldsymbol{\beta}, \boldsymbol{\gamma}} L_{d}(\boldsymbol{\beta}, \boldsymbol{\gamma})}{\min _{x} C_{r}(x)}  \tag{3.60}\\
& \geq \frac{\min _{x} C_{d}(x)}{\min _{x} C_{r}(x)}=\frac{-n+m}{-n}=\frac{2 p+2}{2 p+3} \tag{3.61}
\end{align*}
$$

One might expect that the integration of some global information can enhance the performance of the algorithm. Indeed, such an effect was observed on warm-start modifications where the global information is passed through the initial state Tate et al., 2020 .

### 3.5.7 Parameter optimization

QAOA is a variational algorithm. In a manner, it doesn't define a program in the traditional sense but rather a model (commonly referred as ansatz) that is trained for a special purpose. Indeed, the algorithm operates a family of distributions $\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}(x)$ that can be prepared on a quantum computer. The problem to be solved determines the loss function $L(\boldsymbol{\beta}, \boldsymbol{\gamma})$. The classical part has to train the ansatz: it has to find the values $\boldsymbol{\beta}^{*}, \boldsymbol{\gamma}^{*}$ that minimize the loss function. In order for QAOA to be efficient, the training should be reasonably fast and it should provide an accurate result.

## Why it is difficult to optimize parameters?

There are several obstructions that make optimization challenging. Two of them are related to the cost function structure:

- multiple local minimas. The loss function $L(\boldsymbol{\beta}, \boldsymbol{\gamma})=\langle\psi(\boldsymbol{\beta}, \boldsymbol{\gamma})| C|\psi(\boldsymbol{\beta}, \boldsymbol{\gamma})\rangle$ usually has many arbitrary bad local minimas (see figure 3.8). Thus, if a gradient-based optimization method starts in a bad initial point it can end up in a very poor solution Larkin et al., 2020, Bittel and Kliesch, 2021.
- barren plateaus. The loss function lanscape can have vast flat zones where the gradient vanishes, so the iterative optimizer can't improve the solution even if it is not optimal. This phenomenon was explicitly observed for random parameterized quantum circuits McClean et al., 2018. Fortunately, low-depth QAOA applied to optimize local functions is not expected to exhibit such destructive behavior Cerezo et al., 2021b.

Besides, we usually have access only to a stochastic estimation of the loss function's value:

$$
\begin{equation*}
L(\boldsymbol{\beta}, \gamma) \approx \frac{1}{N} \sum_{i=1}^{N} C\left(X_{i}\right) \tag{3.62}
\end{equation*}
$$



Figure 3.8: Loss function $L\left(\beta_{1}, \gamma_{1}\right)$ landscape (on the left) in QAOA $_{1}$ for MaxCut on the weighted graph (on the right)
where $X_{i}$ are independently sampled from $\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}(x)$. This approach to evaluate the loss introduces a stochastic imprecision that linearly decreases with the number of samples $N$.

Moreover, the samples $X_{i}$ themselves are not exactly the expected ones. In fact, in noisy quantum hardware measurements are not distributed precisely as $\mathbb{P}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}(x)$ but rather according to some noisy versions $\widetilde{\mathbb{P}}_{\boldsymbol{\beta}, \boldsymbol{\gamma}}^{i}(x)$ that are different for each sample $i$. This effect was demonstrated to have a significant impact on the viability of the stochastic estimation Harrigan et al., 2021.

As was already mentioned (see equation 3.42) when $p$ grows the model gains more expressibility and better solves the problem. On the flip side, the training complexity drastically increases with $p$ as any additional layer adds two extra dimensions to the optimization landscape Shaydulin and Alexeev, 2019.

The work Bittel and Kliesch, 2021 shows an example of NP-hard training problem for the number of parameters $p \sim n$. In this work a simple ansatz $W\left(\alpha_{1}, \ldots, \alpha_{n}\right)=\prod_{u=1}^{n} e^{i \alpha_{u} Y_{u}}$ over $n$ parameters is taken to solve the MaxCut problem. For this ansatz it was proven that an efficient approximation algorithm for the optimum of the loss function, if it existed, could be used to solve the MaxCut problem.

## Global methods

Since the parameter landscape has multiple minima, more robust global optimization methods are a good choice for small $p$. The original QAOA paper Farhi et al., 2014a used the grid search. In Larkin et al., 2020 the Asynchronously Parallel Optimization Solver for finding Multiple Minima (APOSMM) was shown to find good-quality parameter values Larkin et al., 2020. In our work Dalyac et al., 2021 the differential evolution was shown to be faster than local methods.

## Initialization strategies

When local optimization methods are used the initialization strategy is extremely important Bittel and Kliesch, 2021.

The method can start from multiple random points, but in a difficult landscape such an approach drastically increases the total runtime. The work Sack and Serbyn, 2021 claims that one can obtain a solution of the same quality with one point corresponding to the approximate annealing schedule as with exponentially many random point.

In Zhou et al., 2020 two initialization strategies are presented. Both reuse the optimum
values at depth $p-1$ to initialize the search at depth $p$. The INTERP strategy uses linear interpolation, while FOURRIER implements more complex ideas of spectrum decomposition.

As the loss function concentrates for the instances from the same distribution, one can reuse optimum values over the same families of instances. The work Brandao et al., 2018 suggests using optimum parameter values for smaller instances to initialize the search for bigger ones.

The concentration phenomenon also inspired numerous machine learning approaches that aim to amortize the optimization cost over several instances. The QAOAKit package presented in Shaydulin et al., 2021] provides pretrained optimal values for different instances of the MaxCut problem. A transfer technique suggested in Shaydulin et al., 2022 allows to reuse QAOAKit for weighted MaxCut. The work Crooks, 2018 trains parameters directly on a family of instances. In Khairy et al., 2020 a kernel density estimation technique is adopted to learn a generative model of QAOA parameters that are close to the optimum.

## Learning-based methods

In addition, several meta-learning techniques were suggested Wilson et al., 2021, Khairy et al., 2020. In contrast to the previously mentioned approaches that reuse the information to output good parameter values or efficient initialization points, in meta-learning techniques the optimizer itself is trained to efficiently guide the search to the optimum. In Wilson et al., 2021 a Lost ShortTerm Memory recurrent neural network is trained to suggest new trial points $\boldsymbol{\beta}_{t+1}, \gamma_{t+1}$ taking in input the gradient of the loss in the point $\boldsymbol{\beta}_{t}, \boldsymbol{\gamma}_{t}$. In Khairy et al., 2020 the reinforcement learning framework is applied to learn a policy network that suggests good iterative updates. Unsupervised learning strategies based on clustering were considered in Moussa et al., 2022.

These learning techniques can significantly reduce the QAOA runtime making it competitive with classical heuristics Crooks, 2018.

## Local methods

In brief, global methods are robust to local minima and meta-learning techniques decrease the overall runtime while improving the solution quality. Nevertheless, of-the-shelf local optimization algorithms are still the most commonly used ones. Due to the presence of multiple local minimas the performance of local heuristics is strongly dependent on the initialization strategy. Moreover, the loss function in general is non-convex, especially when it is estimated on noisy hardware. For instance, because of non-convexity gradient-based methods fail to converge for the Variational Quantum Eigensolver ansatz Peruzzo et al., 2014.

Derivative-free local heuristics such as Nelder-Mead simplex method Gao and Han, 2012 and $C O B Y L A$ Powell, 1994 are a very popular choice for QAOA. It is partially due to a relatively small cost in terms of the number of function evaluations compared to gradient-based methods Guerreschi and Smelyanskiy, 2017. However, more sophisticated gradient-based method Koczor and Benjamin, 2022 is expected to be more efficient both in cost and in the solution quality.

Some works use traditional gradient-based methods such as quasi-Newton gradient-based BFGS Fletcher, 1987 and ADAM Kingma and Ba, 2014 in the classical loop of QAOA. ADAM was reported to find better solutions than other of-the-shell methods Tate et al., 2020.

Yet special routines that are explicitly designed for quantum variational algorithms can be significantly faster. One such technique is Quantum Analytic Descent ( $Q A D$ ) Koczor and Benjamin, 2022. This sophisticated algorithm significantly reduces the number of calls to the quantum computer by using local classical approximations of the loss function to make "big jumps". Quantum mea-
surements are used to evaluate coefficients of the approximation that is further optimized with several steps of the classical gradient descent. The approximation is refined in the new point and, unless the convergence criteria is met, the process restarts. This method was reported to be particularly suitable in the limited function evaluation budget context.

## Gradient of the loss function

During a long time the main obstacle for gradient-based routines was the absence of a convenient way to compute the gradient values, even though some efforts were made in this direction Guerreschi and Smelyanskiy, 2017. Indeed, for the loss function value

$$
\begin{equation*}
L(\boldsymbol{\alpha})=\left\langle\psi_{0}\right| U^{\dagger}(\boldsymbol{\alpha}) O U(\boldsymbol{\alpha})\left|\psi_{0}\right\rangle \tag{3.63}
\end{equation*}
$$

in most cases one have only a stochastic approximation. The situation is even more difficult for the gradient as even the stochastic option is not accessible. Indeed, the gradient of the loss has the form:

$$
\begin{equation*}
\frac{\partial\left\langle\psi_{0}\right| U^{\dagger}(\boldsymbol{\alpha}) O U(\boldsymbol{\alpha})\left|\psi_{0}\right\rangle}{\partial \alpha_{i}}=\left\langle\psi_{0}\right| U^{\dagger}(\boldsymbol{\alpha}) O \frac{\partial U(\boldsymbol{\alpha})}{\partial \alpha_{i}}\left|\psi_{0}\right\rangle+\left\langle\psi_{0}\right| \frac{\partial U^{\dagger}(\boldsymbol{\alpha})}{\partial \alpha_{i}} O U(\boldsymbol{\alpha})\left|\psi_{0}\right\rangle \tag{3.64}
\end{equation*}
$$

where in general the operator $U^{\dagger}(\boldsymbol{\alpha}) O \frac{\partial U(\boldsymbol{\alpha})}{\partial \alpha_{i}}$ as well as its complex conjugate is neither unitary not Hermitian. Therefore, it can't be stochastically approximated given samples from the QAOA distribution.

The alternative is to use finite approximation:

$$
\begin{equation*}
\frac{\partial L}{\partial \alpha_{i}}(\boldsymbol{\alpha}) \approx \frac{L\left(\boldsymbol{\alpha}+\epsilon e_{i}\right)-L\left(\boldsymbol{\alpha}-\epsilon e_{i}\right)}{2 \epsilon} \tag{3.65}
\end{equation*}
$$

where $e_{i}$ is a unit vector with only one non-zero position $i$.
In realistic settings the estimated values of $L\left(\boldsymbol{\alpha}+\epsilon e_{i}\right)$ and $L\left(\boldsymbol{\alpha}-\epsilon e_{i}\right)$ are imprecise due to the hardware noise and the stochastic error. These imprecisions together with the finite approximation error strongly increase the number of measurements that are necessary to obtain a sufficiently good estimation of the gradient value.

For the special case where $U(\boldsymbol{\alpha})=e^{i \alpha \sigma}$ where $\sigma$ is one-qubit Pauli gate $X, Y$ or $Z$ the gradient can be expressed via so-called phase-shift rule Schuld et al., 2019:

$$
\begin{equation*}
\frac{\partial L}{\partial \alpha_{i}}\left(\ldots \hat{\alpha}_{i} \ldots\right)=L\left(\ldots \hat{\alpha}_{i}+\frac{\pi}{4} \ldots\right)-L\left(\ldots \hat{\alpha}_{i}+\frac{\pi}{4} \ldots\right) \tag{3.66}
\end{equation*}
$$

Recently, parameter-shift rules appeared for more complex parameterized unitaries $U(\boldsymbol{\alpha})$ Schuld et al., 2019, Kyriienko and Elfving, 2021, Wierichs et al., 2022, Izmaylov et al., 2021. Using these rules one can evaluate the gradient by querying the function values in multiple points:

$$
\begin{equation*}
\frac{\partial L}{\partial \alpha_{i^{*}}}(\hat{\boldsymbol{\alpha}})=\sum_{k=1}^{m} \epsilon_{k} L\left(\boldsymbol{\alpha}^{\boldsymbol{k}}\right) \tag{3.67}
\end{equation*}
$$

where $\alpha_{i}^{k}=\left\{\begin{array}{ll}\hat{\alpha}_{i}, & i \neq i^{*} \\ \hat{\alpha}_{i}+\phi_{i}, & i=i^{*}\end{array}\right.$. The shifts $\phi_{i}$ and coefficients $\epsilon_{k}$ depend on the spectrum of the parameterized gate.

We highlight that the equalities in the expressions (3.66) and (3.67) are exact, i.e. only stochastic error (in estimations of $L\left(\boldsymbol{\alpha}^{\boldsymbol{k}}\right)$ ) remains when the rules are used to compute the gradient. This may significantly improve the performance of gradient-based optimization strategies.

A Python package PennyLane Bergholm et al., 2018 that was specially designed for variational quantum algorithms fully relies on these rules. We come back to parameter shift rules in chapter (9.

## Which method is better?

The speed and the solution quality are the most important metrics for parameter optimization methods. Another one is the robustness to noise that appears in the evaluations of the loss function. Meta-learning techniques are expected to be more robust as they can potentially learn and integrate the hardware specific effects Wilson et al., 2021.

The Python package scikit-quant Lavrijsen et al., 2020 contains several optimization methods (NOMAD Le Digabel, 2011, ImFil Kelley, 2011, SnobFit Huyer and Neumaier, 2008, $B O B Y Q A$ Powell, 2009). The package was specially designed to have good performance in the presence of noise. This package can directly be used in hardware implementations of variational algorithms.

### 3.6 Recursive QAOA

One of the main reasons why QAOA may fail to find a good approximate solution is its locality Farhi et al., 2020, Hastings, 2019. In other words, at low depth QAOA can't see the global structure of the instance. In order to address this drawback the work Bravyi et al., 2020 suggests to integrate QAOA in another recursive procedure that enforces a non-local treatment in the classical part. The resulting algorithm was named Recursive $Q A O A$ ( $R Q A O A$ ) and was successfully used for MaxCut Bravyi et al., 2020, Egger et al., 2021 and Max-k-Cut Bravyi et al., 2022 problems.

The works Egger et al., 2021 and Patel et al., 2022 suggest techniques for further enhancement of the performance of RQAOA. In Egger et al., 2021 the optimization is warm-started with an additional classical preprocessing. The work Patel et al., 2022 presents an even more advanced approach that integrates RQAOA in a larger reinforcement-learning routine.

Just as Quantum Annealing and the original QAOA the recursive algorithm takes a diagonal Hamiltonian $C_{\text {init }}=\sum_{k=1}^{m} h_{k} \prod_{i \in S_{k}} Z_{i}$ on $n$ variables in input and returns a low-energy state in output. The recursive part executes a variable elimination that stops when the instance is small enough $(n \leq \eta)$ to be addressed by brute-force classical routines.

## Algorithm

At each recursion step the variable elimination routine considers a reduced instance $C_{t}$ over $n-t$ variables $V_{t}$. Initially, $C_{0}$ is taken as $C_{0}=C_{\text {init }}, V_{0}=[1, \ldots, n]$. The choice of the leaving variable is assisted by the QAOA routine for the Hamiltonian $C_{t}$. The procedure consists of the following steps:

1. Consider the family of QAOA states $\left|\psi_{t}(\boldsymbol{\beta}, \gamma)\right\rangle$ for the instance $C^{t}$ over variables $V_{t}$
2. Find the optimal parameter values $\boldsymbol{\beta}_{t}^{*}$ and $\boldsymbol{\gamma}_{t}^{*}$, i.e. solve the loss minimization problem

$$
\begin{equation*}
\boldsymbol{\beta}_{t}^{*}, \boldsymbol{\gamma}_{t}^{*}=\underset{\boldsymbol{\beta}, \boldsymbol{\gamma}}{\operatorname{argmin}}\left\langle\psi_{t}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right| C_{t}\left|\psi_{t}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle \tag{3.68}
\end{equation*}
$$

3. For each variable $u \in V_{t}$ compute the expectation of the observables $Z_{u}$ :

$$
\begin{equation*}
E_{u}=\left\langle\psi_{t}(\boldsymbol{\beta}, \gamma)\right| Z_{u}\left|\psi_{t}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle \tag{3.69}
\end{equation*}
$$

Find the variable $u^{*}$ that has the biggest absolute expectation value:

$$
\begin{equation*}
u^{*}=\operatorname{argmax}\left|E_{u}\right| \tag{3.70}
\end{equation*}
$$

4. Do the same for each pair of variables $w, v \in V_{t}$, i.e. compute:

$$
\begin{equation*}
E_{w, v}=\left\langle\psi_{t}(\boldsymbol{\beta}, \gamma)\right| Z_{w} Z_{v}\left|\psi_{t}(\boldsymbol{\beta}, \boldsymbol{\gamma})\right\rangle \tag{3.71}
\end{equation*}
$$

and find the pair

$$
\begin{equation*}
p^{*}=\left(w^{*}, v^{*}\right)=\operatorname{argmax}\left|E_{w, v}\right| \tag{3.72}
\end{equation*}
$$

5. If $\left|E_{u^{*}}\right| \geq\left|E_{w^{*}, v^{*}}\right|$ fix the value of the variable $u^{*}$ to $\operatorname{val}\left(u^{*}\right)=\sigma\left(E_{u}\right)$ and remove the variable $u^{*}: V_{t+1}=V_{t} \backslash\left\{u^{*}\right\}$. We denote by $\sigma(\cdot)$ the sign function: $\sigma(x)=\left\{\begin{array}{ll}-1, & x \leq 0 \\ 1, & x>0\end{array}\right.$.
6. Otherwise fix the correlation of the variables $w^{*}, v^{*}$ to $\operatorname{val}\left(w^{*}\right) \operatorname{val}\left(v^{*}\right)=\sigma\left(E_{w^{*}, v^{*}}\right)$ and remove the variable $w^{*}: V_{t+1}=V_{t} \backslash\left\{w^{*}\right\}$.
7. Repeat until the number of variables is small: $\left|V_{t}\right| \leq \eta$ and then find the ground state of $C_{t}$ with a brute-force routine.

## How to modify the Hamiltonian

When the routine fixes a value of a variable or a correlation it modifies the Hamiltonian $C_{t} \rightarrow$ $C_{t+1}$. The new instance $C_{t+1}$ is designed to be equivalent to its predecessor on solutions that satisfy the fixed constraint. If $C_{t}$ is an Ising model the modification is straightforward. Let

$$
\begin{equation*}
C_{t}=\sum_{u \in V_{t}} h_{u} Z_{u}+\sum_{u, v \in V_{t}} h_{u, v} Z_{u} Z_{v} \tag{3.73}
\end{equation*}
$$

Fixing $\operatorname{val}\left(u^{*}\right)=\sigma$ for $\sigma \in\{-1,1\}$ implies $Z_{u}=\sigma$. Therefore, the new Hamiltonian $C_{t+1}$ that is equivalent to $C_{t}$ on variable assignments satisfying $\operatorname{val}\left(u^{*}\right)=\sigma$ is

$$
\begin{equation*}
C_{t+1}=\sum_{u \in V_{t} \backslash\left\{u^{*}\right\}} h_{u} Z_{u}+\sum_{u, v \in V_{t} \backslash\left\{u^{*}\right\}} h_{u, v} Z_{u} Z_{v}+\sigma \sum_{v \in V_{t} \backslash\left\{u^{*}\right\}} h_{u^{*}, v} Z_{v} \tag{3.74}
\end{equation*}
$$

Similarly, on assignments that respect $\operatorname{val}\left(w^{*}\right) \operatorname{val}\left(v^{*}\right)=\sigma, \sigma \in\{-1,1\}$ we have $Z_{w^{*}} Z_{v^{*}}=\sigma$ or $Z_{w^{*}}=\sigma Z_{v^{*}}\left(\right.$ as $\left.Z^{2} \equiv I\right)$. The Hamiltonian $C_{t+1}$ is:

$$
\begin{align*}
C_{t+1}= & \sum_{u \in V_{t} \backslash\left\{w^{*}, v^{*}\right\}} h_{u} Z_{u}+\left(h_{v^{*}}+\sigma h_{w^{*}}\right) Z_{v^{*}} \\
& +\sum_{u, v \in V_{t} \backslash\left\{w^{*}, v^{*}\right\}} h_{u, v} Z_{u} Z_{v}+\sum_{v \in V_{t} \backslash\left\{w^{*}\right\}}\left(h_{v^{*}, v}+\sigma h_{v^{*}, v}\right) Z_{v^{*}} Z_{v} \tag{3.75}
\end{align*}
$$

## Motivation of the variable elimination routine

To make the intuition behind the variable elimination clear let's consider an observable $O$ such that $O|x\rangle= \pm r|x\rangle$ on vectors $|x\rangle$ from the computational basis. If in some state $|\psi\rangle$ the expectation $\langle\psi| O|\psi\rangle \approx+r$ one can conclude that the probability distribution $\mathbb{P}(x)=|\langle x \mid \psi\rangle|^{2}$ concentrates on states $|x\rangle$ that corresponds to the eigenvalue $+r$.

QAOA distribution is chosen to favor solutions corresponding to low values of the objective function. Therefore, if in such distribution the mean value of an observable $Z_{u}: Z_{u}|x\rangle= \pm 1|x\rangle$ is close to $\sigma=1$ one can expect that the variable takes the value 1 in most low-energy states. So it makes sense to set $\operatorname{val}(u)=1$.

## Classical simulation

A particularly pertinent fact about RQAOA is that for low depths it can be fully simulated on a classical computer Bravyi et al., 2022, Egger et al., 2021. Indeed, for $p=1$ all expectations $\left\langle\psi_{1}(\beta, \gamma)\right| Z_{u}\left|\psi_{1}(\beta, \gamma)\right\rangle,\left\langle\psi_{1}(\beta, \gamma)\right| Z_{u} Z_{v}\left|\psi_{1}(\beta, \gamma)\right\rangle$ and $\left\langle\psi_{1}(\beta, \gamma)\right| C\left|\psi_{1}(\beta, \gamma)\right\rangle$ can be evaluated using analytical formulas, for example our formula (3.37).

We recall that for the original QAOA even at $p=1$ an efficient simulation algorithm would make the polynomial hierarchy collapse Farhi and Harrow, 2016. Thereby, computing the mean value of an operator in the QAOA state (at least for low depths) is fundamentally easier in terms of the complexity scaling on the instance size than to obtain a sample from the corresponding QAOA distribution. In particular, for $p=1$ we have a polynomial classical algorithm that returns the values $\left\langle\psi_{1}(\beta, \gamma)\right| O\left|\psi_{1}(\beta, \gamma)\right\rangle$ for $O \in\left\{Z_{u}, Z_{u} Z_{v}, C\right\}$ but, unless the polynomial hierarchy collapses, there is no classical algorithm that returns $X \sim \mathbb{P}_{1}(\beta, \gamma)$ in polynomial time.

## Chapter 4

## Smart Scheduling

Due to the rapid development of electric mobility, problems related to vehicle charging gain in importance. In this chapter we consider a scheduling problem, the goal is to assign a set of charges to different charging stations.

Scheduling problems are extremely important in industrial applications. Essentially, a schedule is an ordered assignment of $n$ jobs to $m$ machines. At every moment a job can be processed on only one machine and each machine can process no more than one job.

There are many variants of scheduling problems with diverse machine environments, job characteristics and optimality criteria. Possible machine environments include single processor setup as well as multiple processor setups. In the latter case machines can be parallel, uniform or unrelated. On the other side, jobs themselves have different parameters. For instance, they can be preemptive or not. There may also be a partial order defined on the set of jobs. Such order usually models a precedence requirement where a job can be executed only after some preparation tasks are completed. Diverse problem settings and possible solution techniques are presented in the survey Graham et al., 1979. The survey also introduces the widely adopted classification for scheduling problems in the form $\alpha|\beta| \gamma$. In this expression $\alpha$ relates to the machine settings, $\beta$ - to diverse job characteristics and $\gamma$ - to the optimality criteria.

### 4.1 Problem statement

In the smart scheduling problem we are given a set of charge demands $J,|J|=n$. Each job $j \in J$ has a fixed processing time $t_{j} \in \mathbb{R}^{+}$and priority $w_{j} \in \mathbb{R}^{+}$. Jobs have to be assigned to a set $M$ of $m$ parallel (identical) charging stations. Identical stations imply that the processing time for a job is always $t_{j}$, i.e it doesn't depend on a specific terminal it is assigned to. We assume that jobs are non-preemptive. In other words, once a processor $i \in M$ starts to treat a job the execution can't be interrupted until the job is completed.

A given schedule $S$ determines the completion time $C_{j}$ for each job $j \in J$. Intuitively, the completion time equals the processing time $t_{j}$ plus the time that the job has to wait before being executed:

$$
\begin{equation*}
C_{j}=W_{j}+t_{j}=\left(\sum_{\substack{i \in J \\ i \measuredangle s j}} t_{i}\right)+t_{j} \tag{4.1}
\end{equation*}
$$

where the condition $i \prec_{S} j$ is true if the schedule $S$ assigns the job $i$ to the same machine as $j$ and $i$ is executed before $j$.


Figure 4.1: In realistic setting priority can relate to the importance of a particular vehicle. For example, for a manager of a municipal car park it can be less desirable to delay a police car than a bus.

We are interested in a schedule that minimizes the total weighted completion time: $\sum_{j} w_{j} C_{j}$. In the classification of Graham et al., 1979] the problem can be written as $P \| \sum_{j} w_{j} C_{j}$ if the number of stations is a part of the input and $P m \| \sum_{j} w_{j} C_{j}$ if it is fixed.

### 4.1.1 Complexity and approximations

In some particular cases this scheduling problem can be efficiently solved. For instance, in a single station case $1 \| \sum_{j} w_{j} C_{j}$ a polynomial algorithm directly follows from the Smith's Rule Smith, 1956. According to this rule, it suffices to process jobs in a non-increasing order given by ratios $w_{i} / t_{i}$. For multiple station case the problem can also be solved in $O(n \log n)$ if all priorities $w_{j}$ are equal $\left(P \| \sum_{j} C_{j}\right)$.

However, for general positive priorities and for fixed $m \geq 2$ the problem is NP-hard in the ordinary sense Skutella and Woeginger, 2000. Moreover, if the number of stations $m$ is a part of the input the scheduling problem $P \| \sum_{j} w_{j} C_{j}$ is $N P$-hard in the strong sense Garey and Johnson, 1990. Yet, even if the optimal solution is hard to compute a relatively efficient approximation is achievable. The work Skutella and Woeginger, 2000 presents a polynomialtime approximation scheme (PTAS) that, given an arbitrarily fixed approximation ratio $\alpha=1+\epsilon$ for $\epsilon>0$ computes the approximate optimum in time polynomial in $n$ and $m$. The suggested procedure has an exponential complexity on the precision $1 / \epsilon$. We recall that a family of $1+\epsilon$ approximation algorithms with complexity growing polynomially in $1 / \epsilon$ is called fully polynomialtime approximation scheme (FPTAS). As problems that are NP-hard in a strong sense are not expected to have a FPTAS, the PTAS for $P \| \sum_{j} w_{j} C_{j}$ can't be significantly outperformed.

On the other hand, for a fixed number of stations $m$ the problem $P m \| \sum_{j} w_{j} C_{j}$ is pseudopolynomia $\sqrt{6}$ and admits a FPTAS [Sahni, 1976, Woeginger, 2001. These results rely on a the dynamic programming formulating of the scheduling problem.

The dynamic program maintains a set of partial solutions. Partial solutions are iteratively extended to full solutions. The dynamic algorithm finds an exact optimum but may have an exponential runtime. In order to guarantee a polynomial-time execution the work [Sahni, 1976] suggests an approximate scheme based on rounding of the input data. For $m=2$ and arbitrary $\epsilon>0$ the complexity of the scheme is $O\left(n^{2} / \epsilon\right)$.

[^3]Alternatively, the work Woeginger, 2001 obtains an approximation scheme by performing a trimming of the partial solution space at each step of the dynamic program. This approach can be applied to a wide range of problems that admit a dynamic programming formulation satisfying certain properties (are DP-benevolant). The trimming technique has also inspired a class of polynomial-time evolutionary algorithms with guaranteed approximation ratios Doerr et al., 2011.

Although it cannot be expected that quantum algorithms outperform the specific FPTAS on the given scheduling problem, it is nevertheless interesting to analyze what performance quantum heuristics can achieve.

### 4.1.2 Max-m-Cut reformulation

In order to apply quantum algorithms to a scheduling problem one has to find a corresponding QUBO formulation. One QUBO formulation for $P 2 \| \sum_{j} w_{j} C_{j}$ was suggested in Alidaee et al., 1994. In our work Dalyac et al., 2021 we consider the mapping Skutella, 1998 that associates to each scheduling instance $P m \| \sum_{j} w_{j} C_{j}$ an instance of the weighted Max-m-Cut problem.

More precisely, we define a complete weighted graph $G=(V, E)$ where the nodes $V$ correspond to charging jobs $J$. To each pair of nodes $(u, v), u \in V, v \in V$ is associated a weight $e_{u, v}=\min \left\{w_{u} t_{v}, w_{v} t_{u}\right\}$. We aim to find a partition $V=V_{1} \sqcup \cdots \sqcup V_{m}$ of the set of nodes such that the total weights of edges that are cut by the partition is maximal. An edge $(u, v)$ is cut if its endpoints belong to different members of the partition: $u \in V_{i}, v \in V_{j}, i \neq j$. The Max-m-Cut problem can be written as:

$$
\begin{equation*}
\max _{\mathcal{V}=V_{1} \sqcup \cdots \sqcup V_{m}} \operatorname{cut}(\mathcal{V})=\sum_{1 \leq s<s^{\prime} \leq m} \sum_{u \in V_{s}} \sum_{v \in V_{s^{\prime}}} e_{u, v} \tag{4.2}
\end{equation*}
$$

A solution for the Max-m-Cut problem 4.2 defines an assignment of the charge jobs to different processors. Jobs that belong to the same partition $V_{s}$ are assigned to the same station $s \in M$. Then for each station a locally optimal schedule is obtained using the Smith's Rule. An example of a scheduling problem on 3 stations and the corresponding instance of Max-3-Cut is shown on the figure 4.2 .

In the rest of the section we explain why this reformulation is indeed correct.
For the total completion time in the scheduling problem we have:

$$
\begin{equation*}
\sum_{j} w_{j} C_{j}=\underbrace{\sum_{j} w_{j} W_{j}}_{\mathcal{W}}+\underbrace{\sum_{j} w_{j} t_{j}}_{\mathcal{E}} \tag{4.3}
\end{equation*}
$$

where $\mathcal{E}=$ const that doesn't depend on a particular schedule. The waiting time is $W_{j}=$ $\sum_{i \in P_{j}} t_{i}$ where $P_{j} \subset J$ is the set of jobs scheduled before $j$ on the same station.

If two jobs $i$ and $j$ are assigned to the same station $s \in M$ one of them has to wait for another to be completed. According to the Smith Rule, if $w_{i} / t_{i}>w_{j} / t_{j}$ (or, equivalently, $w_{i} t_{j}>w_{j} t_{i}$ ) the job $j$ will be scheduled after $i$. Therefore, the job $i \in P_{j}$ and the value $w_{j} t_{i}$ will contribute to component $\mathcal{W}$ in the sum 4.3. In other words, if $i$ and $j$ are assigned to the same part $V_{s}$ the weight of the edge $e_{u, v}$ appears in the total weighted completion time:

$$
\begin{aligned}
\sum_{j} w_{j} W_{j} & =\sum_{u, v \in V} e_{u, v} I\left(\exists s: i \in V_{s}, j \in V_{s}\right) \\
& =\sum_{u, v \in V} e_{u, v}\left(1-I\left(u \in V_{s}, v \in V_{s^{\prime}}, s \neq s^{\prime}\right)\right)
\end{aligned}
$$

## Input:

$$
\begin{aligned}
V & =\{a, b, c, d, e\} \\
T & =[1,2,2,1,1] \\
W & =[1,1,1,1,5]
\end{aligned}
$$

## Output:

st. $1: e$
st. $2: b, a$
st. $3: c, d$


Figure 4.2: Example of smart-charging instance for $m=3$ and the corresponding instance of the Max-3-Cut problem. An optimal partition is given by colors.

$$
\begin{equation*}
\equiv \sum_{u, v \in V} e_{u, v}-\operatorname{cut}(\mathcal{V}) \tag{4.4}
\end{equation*}
$$

The last expression (4.4) shows that a partition that maximizes $\operatorname{cut}(\mathcal{V})$ correspond to a schedule that minimizes the total weighted completion time $\mathcal{W}+\mathcal{E}$. Therefore, an algorithm with an approximation ratio $\alpha_{M C}$ for Max-m-Cut can be adapted to solve the scheduling problem with the ratio $\alpha_{S}=\alpha_{M C}+m\left(1-\alpha_{M C}\right)$ Skutella, 1998. For instance, when the Max-2Cut subproblem is addressed with Goemans-Williamson algorithm the completion time of the approximate optimum schedule is of factor $\alpha_{S}=1.122$ from the optimum value.

The finite ratio $\alpha_{S}=1.122$ can't compete with the solution quality of FPTAS for the scheduling problem. However, if Max-m-Cut is approximately solved with an efficient heuristic (for example a quantum heuristic) the runtime advantage can make the Max-m-Cut model attractive for practical use. Besides, a qualification of quantum heuristics on easy-to-approximate instances allows to better understand their computational performance.

### 4.2 Max-m-Cut

In the Max-m-Cut problem one aims to color a set of nodes of a weighted graph $G=(V, E)$ in a way that maximizes the total weight of properly colored edges (see the equation 4.2). A coloring $\mathcal{C}: V \rightarrow[1, \ldots, m]$ defines a partition $\mathcal{V}=\sqcup_{c=1}^{m} V_{c}$ where each part corresponds to one color: $V_{c}=\{v \in V \mid \mathcal{C}(v)=c\}$.

### 4.2.1 The case $m=2$

The case $m=2$ is by far the most studied one both in classical and in quantum computer science. The Max-2-Cut problem is also referred as MaxCut or Max-2-XORSAT. The problem is NPhard. Moreover, if $P \neq N P$ there is no algorithm for unweighted MaxCut with approximation
ratio better than $0.941 \ldots$ Arora et al., 1998. It was shown in Crescenzi et al., 2001 that a weighted version of the problem has essentially the same hardness as the unweighted one.

The best proven approximation ratio for MaxCut is $\alpha_{G W} \geq 0.87856$. This ratio is achieved by the algorithm suggested by Goemans and Williamson Goemans and Williamson, 1995. In the Goemans-Williamson algorithm a SDP-relaxation of the cut problem is solved. The discrete partition is extracted from the continuous solution of the SDP by a special randomized rounding procedure. The bound $\alpha_{G W} \geq 0.87856$ was proven to be tight Karloff, 1999, i.e. there exist instances for which the algorithm achieves exactly the ratio 0.87856 . Moreover, under the Unique Game Conjecture Khot, 2002] the ratio 0.87856 is optimal Khot et al., 2007.

## Unique Game Conjecture

The Unique Game is a one round game with one verifier and two provers that can't communicate. A verifier sends two inputs $s$ and $t$ to the first and the second prover respectively. Each prover sends back one answer from the alphabet $\mathcal{A}$ (also called domain) of size $k=|\mathcal{A}|$. The verifier receives $a, b \in \mathcal{A}$ and accepts the answers if $\sigma_{s t}(a)=b$ where $\sigma_{s t}$ is some permutation on the set $\mathcal{A}$. A value $\nu$ of the game is a maximum success probability that provers can achieve.

Roughly speaking, the Unique Game Conjecture (UGC) states that it is NP-hard to distinguish between almost satisfiable and almost unsatisfiable unique games:

Conjecture: For any $\epsilon>0$ and $\delta>0$ there exists $k=k(\epsilon, \delta)$ such that for a Unique Game with answers from the domain of cardinality $k(\epsilon, \delta)$ it is NP-hard to say if its value $\nu$ is $\nu>1-\epsilon$ or $\nu<\delta$.

Despite significant efforts, the UGC remains unproven. An peculiar results about the validity of the conjecture for quantum Unique Games was reported in Kempe et al., 2007]. In the quantum version of the Unique Game problem provers, while still not allowed to communicate, share an entangled state.

Entanglement is quantum resource that sometimes leads to an increase in the success probability of games. For instance, for a famous CHSH game classical non-entangled provers achieve a value of $75 \%$ while an entangled strategy succeed in $85 \%$ of cases Bell, 1964.

The value of a quantum Unique Game problem is usually referred as entangled value $\nu_{e}$. It turns out that, contrary to the value of classical Unique Games, the entangled value can be efficiently approximated by a semidefinite program Kempe et al., 2007. We recall that the approximation hardness for MaxCut for ratios $\alpha>\alpha_{G W}$ follows from the reduction from the classical Unique Game problem and the conjecture about the classical value $\nu$. Therefore, an efficient approximation for the entangled value $\nu_{e}$ doesn't invalidate the inapproximability bound.

Other than the Goemans-Williamson algorithm, there exist multiple polynomial heuristics solving the MaxCut problem. For instance, MaxCut can be straightforwardly mapped to a QUBO and solved with general QUBO heuristics. It was shown that the Tabu Search have a good performance on QUBO formulation of MaxCut Kochenberger et al., 2013].

Moreover, some restricted classes of instances can be solved to optimality or to any desired approximation ratio $1-\epsilon$ in polynomial time. This is the case for planar graphs Deza and Laurent, 1994. For $k$-regular large girth graphs a near-optimal solution may be computed in $O(n k l)$ under some widely-believed conjecture Alaoui et al., 2021].

Regular graphs often appear in experimental benchmarks for quantum heuristics Farhi et al., 2014a, Basso et al., 2021. If this is the case, the classical competitor has to be carefully chosen. For
instance, for graphs with maximum degree 3 that often appear in QAOA literature an SDP-based algorithm achieves the approximation ratio $0.9326>\alpha_{G W}$ Halperin et al., 2004.

Quantum heuristics for optimization are commonly tested on the MaxCut problem. For lowdepth QAOA several theoretical guarantees were established Hadfield, 2018, Basso et al., 2021, Wurtz and Love, 2021. Yet, these guarantees don't demonstrate a proven quantum advantage. Moreover, there exists a family of instances such that QAOA with depth $p=\Omega(\log n)$ has worse approximation ratio than the Goemans-Williamson algorithm.

Generally, the performance of QAOA heavily depends on the considered instances Larkin et al., 2020. It was observed that QAOA performs better on graphs with many small cycles Wurtz and Love, 2021. A machine-learning based approach that predicts the performance of QAOA on different instances of MaxCut was presented in Moussa et al., 2020.

An instance of the Smart Scheduling problem maps to an instance of weighted MaxCut on a complete weighted graph. Previously, QAOA was applied to MaxCut on complete graphs issued from the Sherrington-Kirkpatrick model Farhi et al., 2022.

### 4.2.2 Complexity and approximations for $m>2$

For general $m>2$ the Max-m-Cut problem is known to be NP-hard to approximate with a factor $\alpha>1-1 /(132 m)$ Kann et al., 1997. A naive heuristic that randomly choses a color for each node has an expected performance of $\alpha_{R}=1-1 / m$.

The work Frieze and Jerrum, 1997 presented an algorithm that outperforms the randomized heuristic. The approximation ratio $\alpha_{m}$ of this algorithm asymptotically scales as:

$$
\begin{equation*}
\alpha_{m} \sim\left(1-\frac{1}{m}\right)+\frac{2 \ln m}{m^{2}} \tag{4.5}
\end{equation*}
$$

If the Unique Game Conjecture holds, the asymptotic scaling 4.5 is essentially optimal. In other words, if there were an algorithm with approximation ratio $\alpha_{m}>1-\frac{1}{m}+\frac{2 \ln m}{m^{2}}$ the Unique Game problem would be easy Khot et al., 2007.

### 4.3 QAOA for Smart Scheduling

In previous sections we have described the mapping of the smart scheduling problem to the Max-m-Cut problem. In the current section we provide a detailed protocol that allows to run QAOA on the Max-m-Cut instances issued from this mapping. Such protocol should specify the building components of QAOA, i.e. the families $U_{P}(\gamma)$ and $U_{M}(\beta)$ of phase and mixing operators respectively. In addition, it should contain a strategy for parameter optimization. In this section we present our protocol (previously reported in our work Dalyac et al., 2021) together with its performance evaluation.

### 4.3.1 Encodings

A natural formulation of the Max-m-Cut problem (sometimes referred as Max-m-colorable subgraph problem) employs integer variables $x_{i} \in[1, \ldots, m]$. The variables correspond to colors assigned to nodes. For a graph with $n$ nodes we take $n$ integer variables $x_{i}$. Then the problem can be formulated as:

$$
\begin{equation*}
\max _{x \in[1, \ldots, m]^{n}} \sum_{(u, v) \in E} e_{u, v} I\left(x_{u} \neq x_{v}\right) \tag{4.6}
\end{equation*}
$$

In theory, the Max-m-Cut problem can be directly encoded using $m$-level physical systems called qudits. QAOA and RQAOA implementions on qudits were applied to the Max-3-Cut problem in Bravyi et al., 2022. According to presented numerical results, both algorithms outperform the classical Newman method while RQAOA achieves better results than the traditional QAOA.

## Unary encoding

The work Wang et al., 2020 suggests to map an integer variable $x_{i} \in[1, \ldots, m]$ to a sequence of $m$ binary variables $\boldsymbol{u}_{\boldsymbol{i}}=\left[u_{i}^{1}, \ldots, u_{i}^{m}\right]$. If an initial variable has the value $x_{i}=k$ all bits in $\boldsymbol{u}_{\boldsymbol{i}}$ are set to zero except $u_{k}$ that is set to one. Such mapping is called unary encoding, it constitutes a state-of-the art approach to deal with bounded integer variables Hadfield et al., 2019.

A sequence on $m$ bits is a valid codeword if the following equality holds:

$$
\begin{equation*}
\sum_{1 \leq j \leq m} u_{i}^{j}=1 \tag{4.7}
\end{equation*}
$$

The energy of the system is given by the Hamiltonian:

$$
\begin{equation*}
C=\sum_{(u, v) \in E} e_{u, v} C_{u, v} \tag{4.8}
\end{equation*}
$$

where for every edge $(u, v) \in E$ the operator $C_{u, v}$ can be as the following Ising model:

$$
\begin{equation*}
C_{u, v}=\frac{1}{4} \sum_{j \in m}\left(1+Z_{u}^{j} Z_{v}^{j}\right) \tag{4.9}
\end{equation*}
$$

Acting on a bitstring $|\boldsymbol{x}\rangle$ that is a valid codeword (i.e. $x$ satisfies 4.7) the operator $C_{u, v}$ takes the value 1 if the edge $(u, v)$ is cut and 0 otherwise.

For the Ising energy operator $C$ the circuit implementing the energy evolution $U_{C}(\gamma)=e^{i \gamma C}$ is straightforward. The genuine implementation challenge emerges from the validity constraints (4.7).

One approach is to penalize unfeasible solutions in the objective function with the term $\lambda\left(1-\sum_{1 \leq j \leq m} u_{i}^{j}\right)^{2}$. For $\lambda$ that is big enough unfeasible solutions will be significantly suboptimal. However, the penalization may impact the performance of the quantum heuristic by "flattening" the energy landscape around feasible solutions.

Another approach is to modify the initial state $\left|\psi_{0}\right\rangle$ and the mixing family $U_{M}(\beta)$ in order to maintain the evolution in a subspace of valid codewords. This approach was developed in Wang et al., 2020 and Hadfield et al., 2019. It proposes to replace the default operator $B$ with an XY-Hamiltonian:

$$
\begin{align*}
B & =\sum_{u \in V} B_{u}  \tag{4.10}\\
B_{u} & =\sum_{1 \leq i<j \leq m} X_{i} X_{j}+Y_{i} Y_{j} \tag{4.11}
\end{align*}
$$

The unary encoding is demanding in number of qubits: $n \times m$ binary variables are necessary for a graph with $n$ nodes. Furthermore, there is no obvious way to compile QAOA components into a sequence of elementary gates. Indeed, the unitary $e^{i \beta B_{u}}$ with an all-to-all connections
can't be decomposed in a simple product of multiqubit Pauli rotations. As the terms $\sigma_{i} \sigma_{i+1}$ and $\sigma_{i-1} \sigma_{i}$ don't commute ( $\sigma$ is Pauli-X or Pauli-Y gate), the compilation issue remains open even for a simpler ring mixer:

$$
\begin{equation*}
B_{u}=\sum_{1 \leq i \leq m} X_{i} X_{i+1}+Y_{i} Y_{i+1} \tag{4.12}
\end{equation*}
$$

## Binary encoding

In our work Dalyac et al., 2021 we suggested to use a more resource-efficient binary encoding. For $m=2^{l}$ each integer variable $x_{i} \in[0, \ldots, m-1]$ is mapped to a sequence of $l=\log _{2} m$ binary variables. A bitstring $\boldsymbol{b}_{\boldsymbol{i}}=\left[b_{0}, \ldots, b_{l-1}\right]$ corresponds to the integer $x_{i}=\sum_{0 \leq j \leq l-1} 2^{j} b_{j}$. In this encoding every binary string of length $l$ is a valid codeword. The initial state and mixing can be taken as in constraint-free QAOA:

$$
\begin{align*}
& \left|\psi_{0}\right\rangle=\frac{1}{\sqrt{2^{n}}} \sum_{x \in\{0,1\}^{n}}|x\rangle  \tag{4.13}\\
& U_{M}(\beta)=e^{i \beta B}, \quad B=\sum_{v \in V} \sum_{0 \leq i \leq l-1} X_{v}^{i} \tag{4.14}
\end{align*}
$$

In binary encoding the energy can be written as a sum over edge set:

$$
\begin{equation*}
C=\sum_{(u, v) \in E} e_{u, v} C_{u, v} \tag{4.15}
\end{equation*}
$$

where each edge in represented by the operator:

$$
\begin{equation*}
C_{u, v}=I(\mathcal{C}(u) \equiv \mathcal{C}(v))=\prod_{0 \leq i<l} \frac{\left(1-Z_{u}^{i} Z_{v}^{i}\right)}{2} \tag{4.17}
\end{equation*}
$$

The expression 4.17 contains higher order monomials $M_{Q}=\prod_{i \in Q} Z_{i},|Q|>2$. Thus, it is beyond the realm of a pure Ising model. Nevertheless, the phase operator

$$
\begin{equation*}
U_{C}(\gamma)=e^{i \gamma C}=\prod_{Q} e^{i \gamma c_{q} M_{Q}} \tag{4.18}
\end{equation*}
$$

can be easily compiled to a quantum circuit. Indeed, for the monomial $M_{Q}==\prod_{i \in Q} Z_{i}$ of degree $|Q|=d$ the circuit for $e^{i \gamma c_{q} M_{Q}}$ is given by:

where $Q=\left\{q_{1}, \ldots q_{d}\right\}$. The decomposition 4.19) contains $O(d)$ CNOTS.
Due to the presence of CNOTS the binary encoding, while using less qubits $\left(n \times \log _{2} m\right)$, leads to circuits of higher depth. The tradeoff between the number of qubits and the depth of the circuit
for different encodings of the Traveling Salesman problem was discussed in Glos et al., 2022. Deeper circuits are more affected by noise. Nevertheless, in certain settings (for instance for quantum computers with small number of qubits) a higher-degree formulation may be beneficial Glos et al., 2022.

Our approach allows to tackle problems with $m$ that is power of 2. In Fuchs et al., 2020 the binary encoding was extended to deal with arbitrary natural $m \in \mathbb{N}$. A sequence of $l=\left\lceil\log _{2} m\right\rceil$ qubits is used to encode the color of a node. Bitstrings that are decoded in "overlong" integers $x \in\left[m, \ldots, 2^{l}-1\right]$ are "stitched" to represent the same color. This extension removes the requirement $m \equiv 2^{l}$. The standard initial state 4.13 and mixing operators 4.14 are still accurate is this scheme. On the down side, stitching leads to sophisticated phase operators that compile to circuits with an exponential amount $O\left(e^{\left(2^{l}-m\right)}\right)$ of CNOT gates.

### 4.4 Numerical results

### 4.4.1 Experimental setup

We analyzed the performance of QAOA on graphs corresponding to the Smart Scheduling problem. Instances were generated out of a real-world dataset containing about 2250 loads performed during May 2017 on identical charging points of the Belib's network. Considered Belib load stations are located in Paris, France Bel, 2017]. The dataset provided realistic charge durations $t_{j} \in \mathbb{R}^{+}$but not the priorities. We considered integer priorities $p_{j} \in \mathbb{N}$. The values of $p_{j}$ were independently sampled from the Poisson distribution:

$$
\begin{equation*}
\mathbb{P}\left(p_{i}=k\right)=\frac{\lambda^{k} e^{-\lambda}}{k!} \tag{4.20}
\end{equation*}
$$

Reported numerical results correspond to an ideal (noise-free) QAOA. For $p=1$ QAOA was simulated on a regular laptop. For higher depths $p>1$ the computations were performed on the Quantum Learning Machine provided by Atos QLM, 2022.

For each number of nodes ranging from 6 to 250 we generated $N=100$ different instances of the $P 2 \| \sum_{j} w_{j} C_{j}$ problem.

We decided to restrict the numerical simulations to the case $m=2$ for multiple practical reasons.

Firstly, even on the QLM supercomputer the runtime of the brute-force circuit simulation is prohibitive for $n \geq 30$ binary variables. Therefore, even for $m=3$ we can't simulate QAOA on instances containing more than 10 jobs. Unfortunately, numerical results on such small instances can't lead to any sound conclusions about the scaling of the algorithm McClean et al., 2021.

In addition, we were able to derive an analytical formula for the loss function for the weighted MaxCut 3.37). The analytical formula allows to evaluate the approximation ratio of QAOA at $p=1$ for relatively big instances. For instance, we were able to treat the instances with $n_{\max }=250$ nodes.

### 4.4.2 Establishing the protocol for QAOA

## Initialization strategy

The work Brandao et al., 2018 reports the concentration phenomena for QAOA on unweighted MaxCut. Roughly speaking, according to this phenomenon QAOA parameters have similar
values on instances that are sampled from the same distribution. We have observed the same behavior for the optimal parameters on instances of the scheduling problem (see fig. 4.3).


Figure 4.3: Optimal parameters $\beta_{1}$ and $\gamma_{1}$ for QAOA with depth $p=1$ on instances of Smart Scheduling of sizes $n \in[6,8,10,15,30,50,70,100,150,170,200,220,250]$. Optimal values are computed with grid search. We observe that for the instances of the same size $n$ the optimal parameters take values that are close. Concentration is more important for bigger instances. Moreover, the absolute values of parameters get closer to zero with $n$, i.e. QAOA approaches the randomized algorithm.

Having confirmed the expectation about parameter concentration, we decided to use this phenomenon to predict good initialization points for different sizes $n$. Notably, in our experiments we initialized local search routines in a barycenter of previously seen parameter values $\left(\beta_{i n i t}^{n}, \gamma_{i n i t}^{n}\right)=\left(\sum \beta_{i}^{n} / N, \sum \gamma_{i}^{n} / N\right)$. We have observed that starting from such point, local methods were able to converge to the global optimum on all instances from the dataset.

For higher depths $p>1$ we used the INTERP approach presented in Zhou et al., 2020. The INTERP strategy relies on the assumption that optimal parameters at depth $p+1$ are close to optimal parameters at depth $p$. More precisely, the strategy suggests to initialize a local search at depth $p+1$ in the point $\boldsymbol{\beta}^{i}=\left[\beta_{1}^{i}, \ldots, \beta_{p+1}^{i}\right], \gamma^{i}=\left[\gamma_{1}^{i}, \ldots, \gamma_{p+1}^{i}\right]$ such that

$$
\begin{align*}
\beta_{j}^{i} & =\frac{j-1}{p} \bar{\beta}_{j-1}+\frac{p-i+1}{p} \bar{\beta}_{j}  \tag{4.21}\\
\gamma_{j}^{i} & =\frac{j-1}{p} \bar{\gamma}_{j-1}+\frac{p-i+1}{p} \bar{\gamma}_{j} \tag{4.22}
\end{align*}
$$

where $\left[\bar{\beta}_{1}, \ldots \bar{\beta}_{p}\right]$ and $\left[\bar{\gamma}_{1}, \ldots \bar{\gamma}_{p}\right]$ are optimal parameters for QAOA at depth $p$. The auxiliary values $\bar{\gamma}_{0}, \bar{\gamma}_{p+1}, \bar{\beta}_{0}$ and $\bar{\beta}_{p+1}$ are set to zero. Our observations match the assumption behind the INTERP strategy (see fig. 4.4).


Figure 4.4: Optimal QAOA parameters at different depth $p$ for an instance with $n=8$ charges. We can observe that the shape of parameter curve at some depth doesn't significantly change when the depth is incremented by 1 .

## Optimization method

For the parameter optimization we compared different local optimization methods implemented in the scikit package for Python (see fig. 4.5). Methods were evaluated on two metrics. The first one is optimization cost. We measured the cost in number of function evaluations made by the method until convergence. The second metric is the quality of the returned optimum that is measured by the approximation ratio:

$$
\begin{equation*}
\alpha=\frac{L\left(\boldsymbol{\beta}^{o}, \gamma^{o}\right)}{c^{*}}, \quad L((\beta), \gamma)=\langle(\beta), \gamma| C|(\beta), \gamma\rangle \tag{4.23}
\end{equation*}
$$

where $c^{*}$ is the optimum cut and $\boldsymbol{\beta}^{o}, \gamma^{o}$ are parameter values returned by the evaluated local routine. In our experiment the exact solutions $c^{*}$ were obtained with the dynamic program for the scheduling problem Sahni, 1976.

According to our simulations, Nelder-Mead is the best choice both in terms of the solution quality and the expected runtime. However, we didn't evaluate its resilience to the imprecision in the evaluation of the loss $L(\boldsymbol{\beta}, \gamma)$.

In the simulations of the QAOA on a perfect hardware the stochastic noise appears when the loss function is approximated by the average energy of samples returned by a quantum computer:

$$
\begin{equation*}
L(\boldsymbol{\beta}, \boldsymbol{\gamma})=\frac{1}{M} \sum_{i=1}^{M} C\left(X_{i}\right), \quad X_{i} \sim \mathbb{P}(\boldsymbol{\beta}, \boldsymbol{\gamma}) \tag{4.24}
\end{equation*}
$$

In realistic settings the noise is induced both by the hardware imperfections and by the stochastic imprecision. In the comparative study reported in Wilson et al., 2021 Nelder-Mead was shown to be particularly affected by the noise.

Taking in account previously mentioned observations, we suggest a protocol that consists of a manually tuned initialization strategy and a privileged optimization method. For $p=1$ our initialization strategy recycles already known optimal parameters. For $p>1$ we adopt the INTERP approach as its motivation seems to be valid in our application. We also suggest to use the Nelder-Mead local routine as it has good performance and low cost compared to other evaluated methods.

Average QAOA performance on graphs with $|V|=10$ nodes for different numerical optimization methods


Optimization cost on graphs with $|V|=10$ nodes for different numerical optimization methods


Figure 4.5: A comparison of local optimization routines for the training of QAOA parameters. Vertical lines correspond to the error bars. The derivative-free Nelder-Mead method Gao and Han, 2012 achieves almost the same quality as the quasi-Newton BFGS method Fletcher, 1987 (on the top figure Nelder-Mead curve hides the BFGS one). Both routines perform better than the frequently used COBYLA optimizer Powell, 1994. Still Nelder-Mead needs less time to converge.

### 4.4.3 Performance evaluation

As expected, we have observed that QAOA performance grows with the depth $p$ (see fig. 4.6). The increase of the approximation ratio with $p$ also indicates that our approach for parameter optimization finds fairly good parameter values. Indeed, as the parameter optimization gets more challenging with $p$ a poor strategy can lead to the decrease in the quality of returned solutions Shaydulin and Alexeev, 2019.

Surprisingly, we also noticed that the approximation ratio gets better with the size $n$ of the scheduling instance. As illustrated on Figure 4.7 the similar behavior is observed for the random assignment algorithm. From the curves on Figure 4.7 we can conclude that QAOA outperforms


Figure 4.6: QAOA performance on instances of MaxCut generated from the Smart Scheduling on realistic data. The running time of the simulations scales exponentially with size on the instance.


Figure 4.7: Approximation ratio achieved by $\mathrm{QAOA}_{1}$ and a random algorithm on Smart Scheduling instances of different sizes $n$.
the randomized heuristics. Yet, the advantage vanishes with $n$. The good performance of the randomized algorithm suggests that the difficulty of the instances under consideration decreases with the size.

In what follows we sketch a probabilistic argument supporting this conclusion. This argument is not a formal proof but rather a reflection about the reasons that can explain the scaling behavior.

## Why the problem gets easier with size?

We have an instance of MaxCut on complete weighted graph $G=(V, E)$. A complete graph on $n$ nodes has $m=\frac{n(n-1)}{2}$ edges. A randomized algorithm on such instance returns the cut with
expected value

$$
\begin{equation*}
\nu_{r}=\frac{\sum_{(u, v) \in E} w_{u, v}}{2} \tag{4.25}
\end{equation*}
$$

Let's assume that the weights of edges are independently sampled from some probability distribution $w_{u, v} \sim \mathcal{D}\left(\mu, \sigma^{2}\right)$ with mean value $\mu$ and variance $\sigma^{2}$. We aim to evaluate the scaling of the average-case behavior of the randomized algorithm with $n$.

In a complete graph a partition can't cut more that $L=\frac{n}{2} \times \frac{n}{2}=\frac{m}{2}+\frac{n}{4}$ edges. The worst performance of the randomized algorithm corresponds to the case when the optimal cut has exactly $L$ edges and, moreover, these are edges of the highest weights. In other words, we can bound the value of the optimal cut by:

$$
\begin{equation*}
\nu^{*} \leq \max _{\substack{E_{l} \subset E \\\left|E_{l}\right|=L}} \sum_{e \in E_{l}} w_{e} \tag{4.26}
\end{equation*}
$$

The number of subsets $E_{l}$ that have the cardinality $L$ can be evaluated with the binomial coefficient:

$$
\begin{equation*}
N=\binom{|E|}{L}=\binom{m}{m / 2+n / 4} \leq\binom{ m}{m / 2} \leq\binom{ 2 L}{L} \approx \frac{2^{2 L}}{\sqrt{\pi m}} \tag{4.27}
\end{equation*}
$$

The last approximate equality in (4.27) is derived from to the Stirling's formula Keller and Vanden-Broeck, 2007.
For each subset $E_{l}$ we introduce new random variable

$$
\begin{equation*}
X_{l}=\frac{\sum_{e \in E_{l}} w_{e}}{L} \tag{4.28}
\end{equation*}
$$

The central limit theorem states that for big $L$ variables $X_{l}$ are distributed according to the normal distribution CLT, 2008:

$$
\begin{equation*}
X_{l} \sim \mathcal{N}\left(\mu, \frac{\sigma^{2}}{L}\right) \tag{4.29}
\end{equation*}
$$

The cumulative distribution function $(C D F) \Phi(x)=\mathbb{P}\left[X_{l} \leq x\right]$ for the distribution $\mathcal{N}\left(\mu, \frac{\sigma^{2}}{L}\right)$ can be written as:

$$
\begin{equation*}
\Phi(x)=\frac{1}{2}\left[1+\operatorname{erf}\left(\frac{\sqrt{L}(x-\mu)}{\sigma \sqrt{2}}\right)\right] \tag{4.30}
\end{equation*}
$$

where the error function $\operatorname{erf}(x)$ satisfies the bound Chiani et al., 2003:

$$
\begin{equation*}
\operatorname{erf}(x) \geq 1-e^{-x^{2}} \tag{4.31}
\end{equation*}
$$

Using this bound we can get for any $a>0$ :

$$
\begin{align*}
\Phi(\mu+a) & =\frac{1}{2}\left[1+\operatorname{erf}\left(\frac{\sqrt{L} a}{\sigma \sqrt{2}}\right)\right]  \tag{4.32}\\
& \geq 1-\frac{1}{2} e^{-L a^{2} / 2 \sigma^{2}} \tag{4.33}
\end{align*}
$$

We now consider $N$ independent variables $\left[X_{1}, \ldots, X_{N}\right]$ identically distributed as 4.29). Actually, the initial variables $X_{l}=\frac{\sum_{e \in E_{l}} w_{e}}{L}$ are not independent as two sets $E_{l}$ may have nonempty intersection. We will come back to the impact of the independence assumption later in this section.

We evaluate the probability that the maximum value of these variables doesn't exceed the mean value by a fixed shift $a$ :

$$
\begin{align*}
\mathbb{P}\left[\max _{l \in[1, \ldots, N]} X_{l} \leq \mu+a\right] & =\mathbb{P}\left[\cap_{l \in[1, \ldots, N]}\left(X_{l} \leq \mu+a\right)\right]  \tag{4.34}\\
& =\prod_{l \in[1, \ldots, N]} \mathbb{P}\left[X_{l} \leq \mu+a\right]=\Phi(\mu+a)^{N} \tag{4.35}
\end{align*}
$$

In this expression both $N$ and $\Phi(\mu+a)$ have an exponential dependence on the number $L$. In order to evaluate how this probability scales with $L$ we use 4.33, 4.27):

$$
\begin{equation*}
\log \Phi(\mu+a)^{N}=N \log \Phi(\mu+a) \geq \frac{2^{2 L}}{\sqrt{\pi L}} \log \left(1-\frac{1}{2} e^{-L a^{2} / 2 \sigma^{2}}\right)=A(L) \tag{4.36}
\end{equation*}
$$

where $L=\frac{n^{2}}{4}$.
Unfortunately, the bound 4.36 is too loose. For big $n$ it only leads to a trivial bound $\mathbb{P}\left[\max _{l \in[1, \ldots, N]} X_{l} \leq \mu+a\right] \geq 0$.

The poor result is caused by the large value $N=\binom{|E|}{m}$. In reality, there is only $K=\binom{n}{n / 2}$ partitions that cut exactly $L$ edges. Intuitively, for identically distributed edge weights the optimal partition on large instances should be one that splits nodes on two equal parts. Assuming that only one of these partitions can lead to an optimal value, the probability that the optimal cut is not too far from the mean value becomes:

$$
\begin{align*}
& \mathbb{P}\left[\max _{l \in[1, \ldots, \boldsymbol{K}]} X_{l} \leq \mu+a\right]=\Phi(\mu+a)^{\boldsymbol{K}} \geq e^{B(n)}  \tag{4.37}\\
& B(n)=\frac{2^{n}}{\sqrt{\pi n / 2}} \log \left(1-\frac{1}{2} e^{-n^{2} a^{2} / 2 \sigma^{2}}\right)  \tag{4.38}\\
& K=\binom{n}{n / 2} \tag{4.39}
\end{align*}
$$

The graph for $B(n)$ is shown on fig. 4.8. According to this plot, the probability that the optimum is close to the mean value converges to 1 for large $n$.

Finally, we recall that we have assumed that the variables $X_{l}$ are independent while this is not exactly true in our setup. Indeed, two variables $X_{1}=\sum_{e \in E_{1}} w_{e} / L$ and $X_{2}=\sum_{e \in E_{2}} w_{e} / L$ can depend on the same terms. In such case they are positively correlated. Indeed, we can decompose:

$$
\begin{align*}
& X_{1}=A+B  \tag{4.40}\\
& X_{2}=A+C \tag{4.41}
\end{align*}
$$

where $A=\sum_{e \in E_{1} \cup E_{2}} w_{e} / L, B=\sum_{e \in E_{1} \backslash E_{2}} w_{e} / L$ and $C=\sum_{e \in E_{2} \backslash E_{1}} w_{e} / L$ are independent from each other. The covariance of $X_{1}$ and $X_{2}$ is:

$$
\begin{equation*}
E\left[X_{1} X_{2}\right]-E\left[X_{1}\right] E\left[X_{2}\right]=E[(A+B)(A+C)]-E[A+B] E[A+C]=E\left[A^{2}\right]-E[A]^{2}>0 \tag{4.43}
\end{equation*}
$$



Figure 4.8: On both graphs we take $\sigma=1$ and $a=\sqrt{2}$. The plot on the left corresponds to the bound $A(L)=\frac{2^{2 L}}{\sqrt{\pi L}} \log \left(1-\frac{1}{2} e^{-L a^{2} / 2 \sigma^{2}}\right)$.
On the right we show $B(n)=\frac{2^{n}}{\sqrt{\pi n / 2}} \log \left(1-\frac{1}{2} e^{-n^{2} a^{2} / 2 \sigma^{2}}\right)$

According to Nadarajah and Kotz, 2008 the maximum of two Gaussian random variables with correlation $r$ is distributed as:

$$
\begin{equation*}
p(x)=2 \phi(x) \Phi\left(\frac{1-r}{\sqrt{1-r^{2}}} x\right) \tag{4.44}
\end{equation*}
$$

where $\phi(x)$ and $\Phi(x)$ are density and CDF for the standard normal distribution $\mathcal{N}(0,1)$. Probability distributions for different positive correlation values $r$ are shown on the figure 4.9. We observe that for positively correlated random variables the distribution weight "goes left", i.e. we the expected maximum gets smaller. Motivated by this observation, we bound the maximum of positively correlated normal variables by the maximum of independent variables sampled from the same individual distributions.


Figure 4.9: Distribution of $\max \left\{X_{1}, X_{2}\right\}$ for different correlation values $r=[0, \ldots, 1]$. The value $r=0$ corresponds to independent variables.

### 4.4.4 Possible improvements

Our QAOA protocol for the Smart Scheduling problem may be improved in several ways.

On one hand, we can profit from the concentration to design advanced training techniques to accelerate the the search of optimal parameters. For instance, a set of pretrained parameters for different instances of MaxCut Shaydulin et al., 2021 can be transfered to weighted instances Shaydulin et al., 2022. We can use these results to gain in the runtime and quality of the optimization routine.

On the other hand, it is possible to enhance the performance of QAOA by adding some classical preprocessing. For instance, classical information about the optimum can be extracted from the solution of the SDP relaxation or, alternatively, the low-rank Burer-Monteiro relaxations. Warm-start approaches Egger et al., 2021, Tate et al., 2020 suggest to integrate this classical information inside the initial state $\left|\psi_{0}\right\rangle$.

## Chapter 5

## Charging Task Selection

In the Smart Scheduling problem the starting time of the jobs $J$ could be chosen freely. In some other popular scheduling problems jobs are provided with release dates $r_{j}, j \in J$ and deadlines $d_{j}, j \in J$. If the difference $d_{j}-r_{j}$ is strictly equal to the duration $t_{j}$ of the job then it has to be executed precisely in the interval $\left[r_{j}, d_{j}\right]$. In such setup the decision-maker is given a set of jobs with intervals and it has to chose which jobs will be accepted for the execution.

### 5.1 Problem statement

In the Charge Task Selection problem each load request is given a predefined starting time $T_{j}^{0} \in \mathbb{R}^{+}$and processing time $t_{j} \in \mathbb{R}^{+}$. Charges are non-preemptive: once started they can't be interrupted until completion. If accepted, a job $j$ brings a profit $p_{j} \in \mathbb{R}^{+}$. Decision-maker aims to select a valid schedule of highest profit. A schedule is valid if selected jobs are simultaneously satisfiable. For instance, if the charges are accomplished on a single station one can't accept two jobs $i$ and $j$ with overlapping time intervals $\left[T_{i}^{0}, T_{i}^{0}+t_{i}\right] \cap\left[T_{j}^{0}, T_{j}^{0}+t_{j}\right] \neq \emptyset$.

Charge Task Selection belongs to the class of Interval Scheduling problems. In essence, this class implements a demand-oriented philosophy: a set of mandatory constrains is given by clients and a logistic manager should satisfy them with minimum amount of resources. In the absence of additional constraints the Interval Scheduling problem on one station can be solved in polynomial time Gavril, 1972. However, most modifications are known to be NP-hard. The survey Kolen et al., 2007 describes different versions of Interval Scheduling problems and provides references to algorithmic solutions.

### 5.1.1 Group constraint

We consider here the Interval Scheduling Problem with additional group constraints. As before, each job comes with an interval $\left[T_{j}^{0}, T_{j}^{0}+t_{j}\right]$. The goal is to select a maximal subset of tasks that can completed on a single station. Obviously, two demands with overlapping time intervals can't be simultaneously satisfied. In addition, each job is assigned to one group. A schedule is valid is it accepts no more than one charge from each group.

Although the one-per-group constraint seems rather artificial, it adapts to different realistic restrictions. For instance, such constraints model an environment where electric vehicles belong to different companies and no company should be privileged by the logistic manager.

Furthermore, a popular Job Interval Selection problem can be naturally formulated as Interval Scheduling problem with group constraint Chuzhoy et al., 2006. In the smart charging
terminology the Job Interval Selection problem suggests for each vehicle a list of intervals when in can be charged. The decision-maker has to select at most one interval per vehicle while avoiding time conflicts.

### 5.1.2 Conflict graph representation



Figure 5.1: Example of the task selection problem. Colors of vehicles represent different groups. We depict the edges of the conflict graph for the yellow vehicle inside the green circle.

An instance of the Charge Task Selection problem with groups can be represented by a conflict graph $G=(V, E)$. Each node $v_{j}$ corresponds to a job $j \in J$. There is an edge ( $v_{i}, v_{j}$ ) between two nodes if the jobs $i$ and $j$ can't be simultaneously satisfied, i.e. they either belong to the same group or their intervals overlap (see fig. 5.1). The resulting graph is a union $G=\left(V, E_{I} \cup E_{G}\right)$ of an interval graph $G=\left(V, E_{I}\right)$ and disconnected cliques modeling the group constraint $G=\left(V, E_{G}\right)$. In literature the term strip graphs is used for the union of an interval graph and a graph made out of cliques Halldórsson and Karlsson, 2006.

Each valid solution for the Charge Task Selection problem corresponds to an independent set on the conflict graph $G$. If we associate to each node a weight $p_{j}$ task selection of maximum profit is equivalent to the maximum weighted independent set on $G$.
We denote by $k \in \mathbb{N}$ the cardinality of the largest group. For each instance one can create an equivalent one where each group has exactly $k$ jobs. The transformation can be done by adding copies of some jobs. An optimum solution for the original instance can be extracted from the solution of the new instance. Therefore, from now we assume that each group contains exactly the same amount $k \in \mathbb{N}$ of jobs. We also take all profits equal $p_{j}=1$. In other words, we want to maximize the amount of satisfied charges.

If each group has only one element the problem is polynomial. However, if $k \geq 2$ the problem is NP-Hard and it doesn't admit a polynomial time approximation scheme [Spieksma, 1999.

A greedy algorithm Spieksma, 1999] achieves the approximation ratio of $1 / 2$. This ratio was improved to $\frac{e-1}{e}$ where $e=2.718 \ldots$ is Euler number in Chuzhoy et al., 2006.

### 5.2 Maximum Independent Set

An independent set (also called stable) in a graph $G=(V, E)$ is a subset of nodes $S \subset V$ such that none of them are connected by an edge: $\forall u, v \in S:(u, v) \notin E$. We denote by $I(G)$ the ensemble of all stable sets of the graph $G$.

A Maximum Independent Set (MIS) in $G=(V, E)$ is a stable $S_{\text {max }}$ of largest cardinality:

$$
\begin{equation*}
S_{\max }=\underset{s \in I(G)}{\operatorname{argmax}}|s| \tag{5.1}
\end{equation*}
$$

The size $\left|S_{\text {max }}\right|$ is called the stability number of the graph and is conventionally denoted by $\alpha(G)$. Maximum independent sets have to be distinguished from maximal independent sets. A stable $S$ is maximal if it can't be extended to a bigger independent set:

$$
\begin{equation*}
\forall v \in V \backslash S \quad \exists u \in S:(u, v) \in E \tag{5.2}
\end{equation*}
$$

Obviously, a maximum independent set is maximal but the inverse is not true.
In the weighted version of MIS (referred as MWIS) a real number $w_{v} \in \mathbb{R}, v \in V$ is associated with each node. The goal is to chose an independent set of biggest weight:

$$
\begin{equation*}
S_{w \max }=\underset{s \in I(G)}{\operatorname{argmax}}|W(s)| \tag{5.3}
\end{equation*}
$$

where for every set $S$ the expression $W(S)=\sum_{v \in S} w_{v}$ is the total weight of all vertices in $S$.

## Related problems

A clique $C$ in a graph $G=(V, E)$ is a set of pairwise connected vertices: $\forall u, v \in C:(u, v) \in E$. The cardinality of a maximum clique in a graph $G$ is called the clique number $\omega(G)$. Clearly, a maximum independent set on a graph $G=(V, E)$ is equivalent to a maximum clique (MC) in a complement graph $\bar{G}=(V, \bar{E})$ where $(u, v) \in \bar{E}$ if and only if $(u, v) \notin E$. Therefore, all results for one problem can be transposed for the second problem.

Graph coloring is another problem closely related to MIS and MC. This problem is an optimization over proper vertex colorings. A proper vertex coloring is a partition of the set of nodes $V$ on $k \in \mathbb{N}$ independent sets $V=S_{1} \sqcup \cdots \sqcup S_{k}, S_{i} \in I(G)$. A minimal natural number $k$ that allows a proper coloring is called chromatic number $\chi(G)$.

The sandwich theorem LLovász, 1979] relates different graph characteristics mentioned above:

$$
\begin{equation*}
\omega(G)<\theta(\bar{G}) \leq \chi(G) \tag{5.4}
\end{equation*}
$$

In this expression $\theta(\cdot)$ is the Lovasz number of the graph. A Lovasz number is equal to the value of some specific SDP program. Contrary to $\omega(G)$ and $\chi(G)$ that are NP-hard to compute, the Lovatz number can be evaluated in polynomial time.

### 5.2.1 Binary formulations of the $\mathrm{M}(\mathrm{W})$ IS problem

A subset of vertices $S$ set can be represented by a binary characteristic vector $x^{S} \in\{0,1\}^{|V|}$ such that:

$$
x_{v}^{s}= \begin{cases}1, & v \in S  \tag{5.5}\\ 0, & v \in V \backslash S\end{cases}
$$

$\mathrm{M}(\mathrm{W}) \mathrm{IS}$ can be formulated as integer linear program Bomze et al., 1999:

$$
\begin{align*}
\max & \sum_{v \in V} w_{v} x_{v}  \tag{5.6}\\
& x_{u}+x_{v} \leq 1, \quad \forall(u, v) \in E  \tag{5.7}\\
& x_{v} \in\{0,1\}, \quad \forall v \in V \tag{5.8}
\end{align*}
$$

More importantly for quantum algorithms, the MIS problem is equivalent to the following QUBO:

$$
\begin{align*}
\min & x^{t} A x  \tag{5.9}\\
& x \in\{0,1\}^{|V|} \tag{5.10}
\end{align*}
$$

where $A=A_{G}-I$ and $A_{G}$ is the adjacency matrix of the graph Bomze et al., 1999. For the weighted version the matrix $A$ is:

$$
a_{i, j}= \begin{cases}-w_{i}, & \text { if } i=j  \tag{5.11}\\ -\frac{1}{2}\left(w_{i}+w_{j}\right), & \text { if }(i, j) \in E \\ 0, & \text { if }(i, j) \in \bar{E}\end{cases}
$$

Motzkin and Strauss presented a similar continuous quadratic optimization problem Motzkin and Straus, 1965

$$
\begin{align*}
& \max x^{T} A_{G} x  \tag{5.12}\\
& \quad \sum_{v \in V} x_{v}=1  \tag{5.13}\\
& x_{v} \geq 0, \quad \forall v \in V \tag{5.14}
\end{align*}
$$

Interestingly, the optimum value of the Motzkin Strauss problem is precisely $1-1 / \omega(G)$.

### 5.2.2 Different graph types

An instance of the Charging task Selection problem without group constraint corresponds to an interval graph. A graph is called an interval graph if there exists a map $I: V \rightarrow \mathbb{R}^{+} \times \mathbb{R}^{+}$that assigns to each node $v$ an interval $\left[T_{v}^{0}, T_{v}^{f}\right]$ on the time line such that two nodes are connected by an edge if and only if the corresponding intervals overlap. Determining if $G=(V, E)$ is an interval graph can be done in linear time $O(|V|+|E|)$ Habib et al., 2000].

Interval graphs are a special case of more general chordal graph. In chordal graphs (also called triangulated graphs) any cycle $\left\{u_{i}\right\}_{i \in C}$ of length greater than four has a chord. Chord in a cycle is an internal edge $\left(u_{i}, u_{j}\right)$ that connects two non-subsequent nodes: $j \neq i+1$.

Chordal graphs in turn form a subset of perfect graphs. A graph is perfect if its clique number equals its characteristic number: $\omega(G)=\chi(G)$. Interestingly, the complement of a perfect graph is also perfect graph Grötschel et al., 1984.

When the group constraint is imposed, the Charging Task Selection problem corresponds to a strip graph - a union of interval graph $G=\left(V, E_{I}\right)$ and a graph made out of several cliques $G=\left(V, E_{G}\right)$.

Our real-world application correspond to a strip graph. On the other hand, some quantum hardware (such as neutral atom chip developed by Pasqal Henriet et al., 2020 is particularly well suited to solve the MIS problem on a distinct family of graphs called unit disks. A graph is called unit disk if there exists a transformation mapping each node $v$ to a point $p_{v}=\left(x_{v}, y_{v}\right)$ in two-dimensional Euclidian space such that two nodes $u$ and $v$ are connected by an edge: $(u, v) \in E$ if and only if corresponding points $p_{u}$ and $p_{v}$ are at distance at most one in $\mathbb{R}^{2}$.

### 5.2.3 Complexity and approximations

The Maximum Independent Set problem on general graph is NP-complete Karp, 1972. Hence, known exact algorithms solving it have an exponential runtime in the worst case.

## Exact algorithms

Most of exact algorithms are based on a Branch \& Bound technique that uses bounds to explore the solution space more efficiently. For instance, the algorithm presented in Balas and Yu, 1986 uses an MIS on induced triangulated subgraph to find a good lower bound. This technique was extended to weighted MIS in Balas and Xue, 1991 and tested on instances with up to 2000 nodes.

The runtime performance of exact algorithms for MIS significantly varies depending on considered instances. Thus, CLIQUER solver developed in [Östergård, 1999] is fast on graphs with high density, while on sparse graphs it is outperformed by the Branch \& Bound approach presented in Held et al., 2012.

## Approximation

The approximation hardness of MIS was proven in several steps. Firstly, the work Garey and Johnson, 1976 demonstrated that if one can approximate MIS to an arbitrary fixed ratio $\alpha$ then the problem admits a PTAS. As this was considered to be unlikely, MIS was not expected to have an approximation algorithm with constant $\alpha$. A rigorous inapproximability result appeared in Håstad, 1999, Zuckerman, 2007. According to this result for all $\epsilon>0$ there is no polynomial algorithm returning for every instance a stable $S^{*}$ such that $\left|S_{\max }\right| /\left|S^{*}\right| \leq n^{1-\epsilon}$ unless $P=$ NP. The best known approximation algorithm Boppana, 1990 achieves the ratio $\left|S_{\text {max }}\right| /\left|S^{*}\right|=O\left(n / \log ^{2} n\right)$ that asymptotically matches the inapproximability bound.

## Special cases

The hardness and the inapproximability results mentioned above describe the worst-case complexity. As expected, they don't hold for various special families of the instances. For instance, there exist polynomial algorithms for MIS on chordal graphs Gavril, 1972 and, more generally, for perfect graphs Grötschel et al., 1984.

On strip graphs MIS problem is fixed-parameter tractable (FPT) with respect to the number of groups $k$. In other words, for fixed $k$ there exists an algorithm polynomial in $|V|$ and $|E|$ that solves the problem. One such algorithm based on dynamic programming was presented in Halldórsson and Karlsson, 2006. It has a runtime complexity $O\left(2^{k} n\right)$. Better time complexity is obtainable for the construction presented in Bevern et al., 2014.

For MIS on unit disk graphs a polynomial time approximation scheme based on local search was presented in Chan and Har-Peled, 2012. A iterative procedure Nieberg et al., 2004 implies a PTAS for the weighted case.

Finally, for random graphs sampled from the same family the numbers $\alpha(G)$ and $\omega(G)$ tend to concentrate. Indeed, it was shown in Matula, 1976 that for Erdos-Renyi random graphs with density $\delta$ the clique number belongs to an interval

$$
\begin{equation*}
\lfloor M(n, \delta)\rfloor \leq \omega(G) \leq\lceil M(n, \delta)\rceil \tag{5.15}
\end{equation*}
$$

with probability going to 1 where

$$
\begin{equation*}
M(n, \delta)=2 \log _{\frac{1}{\delta}} n-2 \log _{\frac{1}{\delta}} \log _{\frac{1}{\delta}} n+2 \log _{\frac{1}{\delta}} \frac{e}{2}+1 \tag{5.17}
\end{equation*}
$$

### 5.2.4 Classical Heuristics

Maximum Independent Set is a very challenging computational problem, so even the most rapid exact algorithms can't tackle large instances. The approximation algorithm Boppana, 1990 has poor guarantees and is relatively slow. As a consequence, in practical applications the MIS problem is usually solved with heuristics.

A recently-reported sophisticated heuristic based on binary search finds optimal solutions on real-world instances with about $10^{6}$ nodes in less than a minute Lu et al., 2017. It can tackle instances with up to $10^{8}$ vertices in several days.

## DIMACS benchmark

Heuristics are commonly characterized by experimental performance. A diversified benchmark for Maximum Clique (MC) solvers was suggested in the DIMACS challenge Johnson and Trick, 1996 Since then it became a reference test dataset in most works on algorithms for MIS and MC.

DIMACS benchmark contains sixty-five instances from eleven different families of graphs issued from wide variety of applications. CFat instances are issued from fault diagnosis, Johnson and Hamming from coding theory. San and Sanr are generated random graphs with predefined clique size. Another family of random graphs is PHat. PHat graphs have wider node degree spread and larger $\omega(G)$ that standard Erdos-Renyi graphs. Brock instances have maximum clique that is much larger than the expected size. MANN instances comes from covering problems. Keller graphs are based on Keller's conjecture on tiling hypercubes.

## Greedy heuristics

Some heuristics such as Lu et al., 2017 are inspired by Branch \& Bound with limited time budget. Another popular research direction is greedy algorithms. In a nutshell, greedy algorithms maintain some partial solution and iteratively add or remove nodes from it depending on the values of indicators.

A naive greedy heuristic Johnson, 1974 suggests to add a new node to an independent set in order defined by degrees. For weighted case the popular greedy strategies consider nodes $v \in A$ according to one of the following orders Held et al., 2012:

- maximum weight: $v=\operatorname{argmax}_{v \in \boldsymbol{A}} w_{v}$
- maximum static surplus: $v=\operatorname{argmax}_{v \in \boldsymbol{A}}\left[w_{v}-\sum_{u \in V,(u, v) \in E} w_{u}\right]$
- maximum dynamic surplus: $v=\operatorname{argmax}_{v \in \boldsymbol{A}}\left[w_{v}-\sum_{\substack{u \in \boldsymbol{A} \\(u, v) \in E}} w_{u}\right]$

In these expressions $\boldsymbol{A}$ denotes the set of available vertices. If in a current step the partial solution is an independent set $S^{t}$ then the set $\boldsymbol{A}$ is $V \backslash\left(S^{t} \cup\left\{u \mid \exists v \in S^{t}:(u, v) \in E\right\}\right)$.

Greedy approaches are particularly fast compared to other heuristics. Moreover, an integration of sophisticated post-processing allows to achieve very good performances. For instance, a Greedy Randomized Adaptive Search (GRAS) heuristic combines local search with a greedy initialization Feo et al., 1994. In experimental evaluation on random graphs with variable densities GRAS was able to find optimal or expected optimal solutions on most instances.

## Local search

Greedy algorithms iteratively construct the solution. A distinct strategy is adopted by local search methods. A typical local search procedure is initialized by some feasible solution $x^{t}$. Thereupon it explores the local neighborhood $N\left(x^{t}\right)$ of the initialization point $x^{t}$. If a better solution $x^{*} \in N\left(x^{t}\right)$ is found, the methods makes a move towards it: $x^{t+1}=x^{*}$. Otherwise, the exploration either stops or moves to some suboptimal solution $x^{t+1}=\hat{x}$ with probability $p\left(x^{t}, \hat{x}\right)$.

There exists a wide variety of local search methods with different notion of neighborhood $N(\cdot)$ and diverse modification strategies $p\left(x^{t}, \hat{x}\right)$. For instance, 2-interchange heuristic replaces a single vertex in a maximal stable $x^{t}$ with two others if this improves the value of the solution. The work Andrade et al., 2008 presents a method that checks if there is a 2 -improvement in a linear time. It also shows that such local search strategy is systematically able to improve a naive greedy solution.

Simulated Annealing is another way to apprehend the local search framework. In this metaheuristic the probability of transition to a suboptimal solution $p\left(x^{t}, \hat{x}\right)$ is proportional to the global temperature parameter $T$. Emulating the controlled cooling in metallurgy, the temperature $T$ is slowly decreased. Simulated Annealing as well as greedy algorithms showed good performance for the Maximum Clique problem on dense graphs Homer and Peinado, 1994.

Tabu search is a local strategy that use the short-term memory to restrict the set of legal neighbors. The basic version is sometimes tuned with intensification and diversification strategies that allows to explore the solution space more efficiently. A refined tabu search Battiti and Protasi, 2001 demonstrated excellent results on the DIMACS benchmark: on most instances it found already known best solutions. Moreover, for some graphs the algorithm improved previous bounds.

Greedy and local search strategies and their various combinations usually achieve satisfactory performance on MIS problem. Up to know, global optimization methods such as genetic algorithms Marchiori, 1998, continuous-based heuristics Gibbons et al., 1993 and neural networks Grossi et al., 2006 were unable to significantly outperform them. Nevertheless, a genetic algorithm presented in Marchiori, 1998 showed results comparable to simulated annealing and tabu search on DIMACS benchmark. The continuous based heuristic Gibbons et al., 1993, in turn, doesn't require a parameter calibration that has a significant impact on performance of metaheuristics.

### 5.3 Quantum heuristics for Charging Task Selection

### 5.3.1 Encoding

The Maximum Independent Set problem is equivalent to the QUBO formulation (5.9). This formulation can be equivalently written as :

$$
\begin{align*}
& \max \sum_{v \in V} x_{v}-\lambda \sum_{(u, v) \in E} x_{u} x_{v}  \tag{5.18}\\
& \quad x \in\{0,1\}^{|V|} \tag{5.19}
\end{align*}
$$

where $\lambda=2$.

## Penalization

The value $\lambda=2$ is sufficient to make solutions that are not independent sets $x^{S}, S \notin I(G)$ suboptimal. However, quantum heuristics don't return an exact optimum but rather some low-energy state. In this settings it is extremely important to properly discriminate unfeasible solutions in the objective function. Therefore, the penalty value $\lambda$ is typically set to a larger value $\lambda \gg 2$. Yet too large value of $\lambda$ will "flatten" the relative energy landscape around feasible solutions reducing the performance of heuristics. Bisection search can be used to accurately fix the value of the penalty term Stollenwerk et al., 2018]. The bisection search simply checks the feasibility of solutions returned by a quantum heuristic on formulations with different values $\lambda$.

Another approach to ensure feasibility was adopted in Farhi et al., 2020]. In this work the value $\lambda$ is taken to be $\lambda=1$ and independent sets are extracted from the output of QAOA by a post-processing procedure.

## Modified mixer Hamiltonian

Furthermore, the feasibility of solution returned by a quantum heuristic can be achieved by an appropriate modification of the driver Hamiltonian. The original QAOA paper Farhi et al., 2014a suggests to use the mixer $\bar{U}_{M}(\beta)=e^{i \beta B^{\prime}}$ where

$$
\langle x| B^{\prime}\left|x^{\prime}\right\rangle= \begin{cases}1, & x, x^{\prime} \in I(G), \quad x \text { and } x^{\prime} \text { differ in one bit }  \tag{5.20}\\ 0, & \text { otherwise }\end{cases}
$$

The operator $\bar{U}_{M}(\beta)$ preserves the evolution in the feasible subspace. However, it is not clear how to design a circuit implementing such evolution. Taking in account the implementability, the work Hadfield, 2018 introduces a sequential mixing $U_{M}^{s}(\beta)=\prod_{v \in V} U_{M, v}^{s}(\beta)$. For each node with $l$ neighbors $N(v)$ the unitary $U_{M, v}^{s}(\beta)$ is $e^{i \beta B_{v}}$ where the operator $B_{v}$ is:

$$
\begin{equation*}
B_{v}=\frac{1}{2^{l}} X_{v} \prod_{u \in N(v)}\left(I+Z_{u}\right) \tag{5.21}
\end{equation*}
$$

### 5.3.2 Maximum Independent set on neutral atoms

For unit disk instances the MIS problem (UD-MIS) is native for neutral atom platforms Henriet et al., 2020, Pichler et al., 2018. Using laser control system such platform creates a energy landscape corresponding to the Hamiltonian:

$$
\begin{equation*}
H=\sum_{i=1}^{N} \frac{\hbar \Omega}{2} X_{i}-\sum_{i=1}^{N} \frac{\hbar \delta}{2} Z_{i}+\sum_{j<i} \frac{C_{6}}{\left|\mathbf{r}_{i}-\mathbf{r}_{j}\right|^{6}} \frac{Z_{i}+1}{2} \frac{Z_{j}+1}{2} \tag{5.22}
\end{equation*}
$$

where $r_{i}$ is a coordinate of a node in the unit disk embedding, $\Omega$ is Rabi frequency and $\delta$ is detuning of the laser system. The energy term for pair interactions prevents two atoms to be in an excited state if they are at distance at most $r_{b}=\left(C_{6} / \hbar \Omega\right)^{\frac{1}{6}}$. At the same time, the interaction is negligible for other pair. This phenomenon is called Rydberg blockage [Urban et al., 2009. The Rydberg blockade imposes the system to be in a state that corresponds to an independent set of the chosen unit disk layout.

A neural atom processor works in an analog mode. Using laser control, one can drive the system to an approximate solution of the UD-MIS problem.

## Embedding in unit disk

In order to apply this technique to generic Charge Task Selection instances one has to find a proper embedding of the instance into a unit disk. In other words, we have to find a mapping (usually referred as realization) from the set of nodes to a set of points in 2D. If we denote by $d\left(p_{1}, p_{2}\right)$ a Euclidean distance between two-dimensional point $p_{1} \in \mathbb{R}^{2}$ and $p_{2} \in \mathbb{R}^{2}$ the embedding problem ( $E P$ ) can be formulated as:

$$
\begin{array}{lr}
\text { Find }\left\{\left(\boldsymbol{v}_{\boldsymbol{i}}\right)\right\}_{i \in V} & \\
& \\
& \\
& \rho \leq d\left(\boldsymbol{v}_{\boldsymbol{i}}, \boldsymbol{v}_{\boldsymbol{j}}\right) \leq r_{b}, \\
d\left(\boldsymbol{v}_{\boldsymbol{i}}, \boldsymbol{v}_{\boldsymbol{j}}\right)>r_{b}, & \forall(i, j) \in \mathbb{E}  \tag{5.26}\\
\boldsymbol{v}_{\boldsymbol{i}} \in[0, \bar{L}]^{2}, & \forall(i, j) \notin \mathbb{E} \\
& \forall i \in V
\end{array}
$$

where $\rho \sim r_{b} / 3$ is the minimal spacing between separated atoms and $\bar{L}$ is a maximum square area to place atoms in a quantum chip. In a physical system typical values of the parameters are $r_{b} \approx 15 \mu m$ for the Rydberg blockade radius, $\rho \approx 5 \mu m$ and $\bar{L} \approx 100 \mu m$ for the platform characteristics.

We remark that the embedding problem doesn't always admit a feasible solution. For instance, a limited number of vertices can be placed in a disk $D\left(\rho, r_{b}\right)$. Therefore, a node with a sufficiently high degree can't be embedded in such layout.

Besides, in order to leave a room for quantum advantage embedding should be efficiently computable.

Because of the constraints $5.24,5.25$ the formulation (EP) is non-convex, thus it is expected to be difficult Park and Boyd, 2017. Indeed, it is NP-hard to determine if the problem admits a solution at all Breu and Kirkpatrick, 1998. Still, in practice (EP) can be transformed to a mixed integer linear program and optimized by a conventional solver such as CPLEX Cplex, 2009. Unfortunately, our experiments revealed that the solution time of this (EP) formulation is too large so it can compromise the potential quantum advantage.

In Dalyac et al., 2021 we suggested a more restrictive formulation (EPR):

$$
\begin{array}{lr}
\text { Find }\left\{\left(x_{i}, y_{i}\right)\right\}_{i \in V} & \\
& \begin{array}{lr}
\text { s.t. } & \\
& \bigvee_{\phi \in \Phi} f_{\phi}\left(x_{i}, x_{j}, y_{i}, y_{j}\right), \\
& \forall(i, j) \in E \\
& \left.\left|x_{i}-x_{j}\right|>r_{b}\right) \vee\left(\left|y_{i}-y_{j}\right|>r_{b}\right), \\
x_{i}, y_{i} \in[0, \bar{L}], & \forall(i, j) \in \bar{E} \\
& \forall i \in V
\end{array}
\end{array}
$$

where $\Phi \subset[0,2 \pi]$ is a manually fixed discrete set of angles and the function $f_{\phi}: \mathbb{R}^{4} \rightarrow \mathbb{B}$ indicates that the point $\left(x_{j}, y_{j}\right)$ is on the radius $\phi$ of the disk $D\left(\rho, r_{b}\right)$ centered in $\left(x_{i}, y_{i}\right)$ :

$$
\begin{equation*}
f_{\phi}\left(x_{i}, x_{j}, y_{i}, y_{j}\right)=\left(\rho \cos (\phi) \leq\left|x_{i}-x_{j}\right| \leq r_{b} \cos (\phi)\right) \wedge\left(\rho \sin (\phi) \leq\left|x_{i}-x_{j}\right| \leq r_{b} \sin (\phi)\right) \tag{5.31}
\end{equation*}
$$

Constraints in the new formulation (EPR) are stronger that the ones in the original (ER). Therefore, the program (EPR) can be infeasible even if the graph actually admits a valid unit disk realization.

On the other hand, (EPR) contains only "or" and "and" logic combinators and absolute values of linear terms. It makes the formulation more tractable by the CPLEX solver. We tested
the modified embedding strategy on 100 instances of the Charge Task Selection problem with number of nodes from 12 to 15 . CPLEX was able to solve the (EPR) problem in $60 s$ for 84 instances. We remark that our test set contained only small graphs. For instances of realistic size the embedding procedure needs further refinements to allow a competitive runtime behavior.

Once an instance is properly embedded in the unit disk layout, the MIS problem can be directly solved on a neutral atom machine. A recent work Dalyac and Henriet, 2022 reported a procedure that finds in polynomial time for graphs with maximum degree 6 an embedding in unit disk of three dimensions where Rydberg atoms are placed in $3 D$ atomic arrays.

If an instance has no UD realization it can be addressed by some classical-quantum heuristic that outsources a part of computation to a quantum chip. One can imagine an heuristic implementing a subgraph approach similar to one in Balas and Yu, 1986 where Branch \& Bound procedure uses independent sets on induced triangulated subgraphs. A classical-quantum heuristic for general MIS can extract unit disk subgraphs from an instance and call a neutral atom chip to approximate MIS on the subgraph.

### 5.3.3 Numerical results

## Experimental setup

In order to evaluate the performance of quantum heuristics on the MIS problem we created a benchmark that contains random regular graphs with degrees in $[3,5,7,9]$ and sizes ranging from 100 to 1000 . We also included DIMACS graphs and real-world instances issued from a simulator of the charging demand scenarios developed internally by EDF.

A full scale simulation of QAOA and Quantum Annealing is limited to small instances $(n \leq$ 30). Therefore, we decided to evaluate the performance of RQAOA. Contrary to QAOA, RQAOA can be efficiently simulated at depth $p=1$ (see section 3.6). Due to the efficient simulation, RQAOA at depth $p=1$ is a purely classical algorithm. Nevertheless, our analysis sheds a light on the typical behavior of quantum heuristics.

Our implementation of $\mathrm{RQAOA}_{1}$ is in a sense the simplest possible one. It uses the analytic formula (3.37) to evaluate the mean values $\left\langle Z_{u}\right\rangle$ and $\left\langle Z_{u} Z_{v}\right\rangle$. In the first iteration of the recursive variable elimination QAOA parameters are optimized with a global differential evolution method. For the following steps the local Nelder-Mead routine is used. Conducting numerical experiments we observed that the parameter optimization step is a genuine bottleneck for our implementation. Indeed, without any specific knowledge about the neighborhood of the optimal solution a global method takes a long time to find a sufficiently good initial point. Local Nelder-Mead optimization requires less function evaluations to converge. However, for each $n^{2}$ pairs the values $\left\langle Z_{u} Z_{v}\right\rangle$ are computed in $O\left(D_{m} a x\right)$ with the formula 3.37). So for dense graphs $d=\Omega(n)$ the runtime becomes prohibitive for graphs with more than 100 nodes.

We believe that more refined parameter optimization strategies (such as described in the section 3.5.7 can accelerate the program and increase the performance of RQAOA.

## RQAOA performance

In the beginning we compared $\mathrm{RQAOA}_{1}$ to the approximation algorithm ( $C A$ ) from Boppana, 1990 on 121 random regular graphs with up to 1000 vertices. For the approximation algorithm we used the implementation from the networkx Python package. We observed that on all but 7 instances (mostly of small sizes 100 and 300 ) $\mathrm{RQAOA}_{1}$ returns a solution that is at least as
good as the one found by the classical routine. Moreover, on 108 instances $\mathrm{RQAOA}_{1}$ is strictly superior.

The DIMACS benchmark is designed for the Maximum Clique problem. A program for MIS can solve the Maximum Clique problem on $G=(V, E)$ by returning an independent set in a complement graph $G=(V, \bar{E})$. The complexity of our program for MIS significantly grows with the number of edges in the graph. For this reasons, we were able to obtain numerical results only for 40 out of 65 DIMACS instances that are either small or sufficiently sparse.

The experiment on DIMACS graphs revealed that RQAOA at $p=1$ shows good performance on some particular instances (such as ones from the Hamming family). Still it returns a poor approximate solution in many cases. This implies that it is necessary to increase the depth $p$ in order to observe the quantum advantage.

On the Charge Task Selection problem we compared RQAOA to an exact algorithm as well as to the classical heuristic both presented in Held et al., 2012 (see fig. 5.2). The exact algorithm is based on Branch \& Bound, it can tackle instances wit up to 500 nodes. The reference heuristic combines a greedy initialization with 2-interchange local search. In order to make a fair evaluation, we decided to improve the RQAOA solutions with the same local search routine. Our results demonstrates that RQAOA, while used as initialization for the local search routine, can improve the solution compared to the classical method.

In summary, our numerical experiments about RQAOA applied to MIS provides encouraging results. While we were unable to definitely outperform classical heuristics, the obtained solutions were still competitive in many cases. Moreover, our analysis was restricted to RQAOA with depth $p=1$. We believe that the performance of the algorithm can be significantly increased if we consider QAOA anzats with higher depth $p>1$.

On the down side, we have observed that the solution quality is not robust across instances. Therefore, the selection of right usecases that lead to instances well-suited for quantum heuristics remains an extremely important challenge for the near-future quantum computing.


Figure 5.2: Comparison of RQAOA with a classical heuristic on Charge Task Selection problem. On the top panel is shown the exact approximation ratio obtained by both algorithms. The bottom panel illustrates the relative performance of two heuristics in the regime $|V|>500$ marked by a violet line. In this regime the exact algorithm gets too slow. Vertical yellow line shows the limit on size of instances that can be analyzed by a full-scale simulation of quantum evolution. The full-scale simulation is necessary to get the output of such heuristic as QAOA.

| Instance | $\|V\|$ | $\|E\|$ | exact $\omega(G)$ | RQAOA |
| :--- | :--- | :--- | :--- | :--- |
| brock200_1 | 200 | 14834 | 21 | 17 |
| brock200_-2 | 200 | 9876 | 12 | 9 |
| brock200_3 | 200 | 12048 | 15 | 10 |
| brock200_-4 | 200 | 13089 | 17 | 13 |
| brock400_-1 | 400 | 59723 | 27 | 17 |
| brock400_-2 | 400 | 59786 | 29 | 18 |
| brock400_3 | 400 | 59681 | 31 | 17 |
| brock400_4 | 400 | 59765 | 33 | 16 |
| c-fat200-1 | 200 | 1534 | 12 | 12 |
| c-fat200-2 | 200 | 3235 | 24 | 23 |
| c-fat200-5 | 200 | 8473 | 58 | F |
| hamming6-2 | 64 | 1824 | 32 | 32 |
| hamming6-4 | 64 | 704 | 4 | 4 |
| hamming8-2 | 256 | 31616 | 128 | 128 |
| hamming8-4 | 256 | 20864 | 16 | 10 |
| hamming10-2 | 1024 | 518656 | 512 | 510 |
| johnson8-2-4 | 28 | 210 | 4 | 4 |
| johnson8-4-4 | 70 | 1855 | 14 | 8 |
| johnson16-2-4 | 120 | 5460 | 8 | 8 |
| keller4 | 171 | 9435 | 11 | 8 |
| MANN_a9 | 45 | 918 | 16 | 12 |
| MANN_a27 | 378 | 70551 | 126 | 117 |
| MANN_a45 | 1035 | 533155 | 345 | 330 |
| p_hat300-1 | 300 | 10933 | 8 | 6 |
| p_hat300-2 | 300 | 21928 | 25 | 19 |
| p_hat300-3 | 300 | 33390 | 36 | 28 |
| san200_0.7_1 | 200 | 13930 | 30 | 15 |
| san200_0.7_2 | 200 | 13930 | 18 | 12 |
| san200_0.9_1 | 200 | 17910 | 70 | 45 |
| san200_0.9_2 | 200 | 17910 | 60 | 35 |
| san200_0.9_3 | 200 | 17910 | 44 | 30 |
| san400_0.5_1 | 400 | 39900 | 13 | F |
| san400_0.7_1 | 400 | 55860 | 40 | 20 |
| san400_0.7_2 | 400 | 55860 | 30 | F |
| san400_0.7_3 | 400 | 55860 | 22 | 12 |
| san400_0.9-1 | 400 | 71820 | 100 | 50 |
| sanr200_0.7 | 200 | 13868 | 18 | 15 |
| sanr200_0.9 | 200 | 17863 |  | 35 |
| sanr400_0.5 | 400 | 39984 | 13 | 8 |
| sanr400_0.7 | 400 | 55869 |  | 17 |

Table 5.1: Performance of RQAOA on DIMACS instances of the Maximum Clique problem. "F" entry for RQAOA signifies that the algorithm returned an infeasible solution on the corresponding instance. The feasibility issue can be addressed by increasing the penalty for infeasible solutions or with a special postprocessing.

## Chapter 6

## Hybrid Quantum-Classical Decomposition Scheme

According to current estimations, the near-future quantum hardware will have a relatively small number of qubits with limited connectivity Preskill, 2018. Besides, existing algorithms adapted to NISQ devices can deal with only several special types of applications. In the context of combinatorial optimization attainable applications are the problems that can be formulated as the optimization of an unconstrained quadratic binary function. Moreover, in order to execute the algorithm on a quantum hardware we have to specify the mapping between the connectivity pattern presented by quadratic terms and the hardware layout.

This obstructions hardly restrict the spectrum of problems that can be addressed directly by quantum algorithms. In order to circumvent these limits, for most applications quantum hardware is assisted by classical routines. Classical routines can be used in preprocessing, postprocessing as well as be directly integrated in variational algorithms such as QAOA Farhi et al., 2014a and VQE Peruzzo et al., 2014]. An algorithm that combines quantum and classical parts is called hybrid algorithm.

## Embedding

The famous example of a classical preprocessing is the DWave embedding procedure that maps logical variables to chains of physical qubits in the hardware layout Choi, 2008. Previously (see section 5.3.2), we presented a classical program that finds a unit disk representation required to address the problem on neutral atom platforms for instances of the Maximum Independent Set problem Pichler et al., 2018, Henriet et al., 2020. In general, embedding extends the applicability of quantum algorithms to more general classes of non-native instances. It is crucial for real-world applications.

In addition to the management of hardware constraints, hybridization can be used to address more complicated formulations that a vanilla QUBO. For instance, sophisticated hybrid approaches were suggested to handle linear inequality constraints Braine et al., 2021, Yonaga et al., 2020.

## Integration of inequality constraints

Management of inequality constraints that are often present in optimization problem is essential but also very challenging. Indeed, contrary to a linear equality constraint $a^{T} x+b=0$ that can be straightforwardly integrated in the QUBO formalism with the penalty term $\left(a^{T} x+b\right)^{2}$,
penalization of inequalities requires additional slack variables. With slack variables an inequality constraint

$$
\begin{equation*}
a^{T} x+b \leq 0 \tag{6.1}
\end{equation*}
$$

is transformed to

$$
\begin{align*}
& a^{T} x+b+s=0  \tag{6.2}\\
& s \geq 0 \tag{6.3}
\end{align*}
$$

The slack variable $s$ is in general non-binary (and even non-integral). Usually, we can estimate only a poor upper bound $s \leq u \leq 2^{k}$. If the parameters $a$ and $b$ are integer-valued, the upper bound allows to represent $s$ with a binary expansion $s=\sum_{q=0}^{k} 2^{q} y_{q}, y \in\{0,1\}^{k}$. In a naive approach for the management of inequality constraint the penalty term $\left(a^{T} x+b+\sum_{q=0}^{k} 2^{q} y_{q}\right)^{2}$ is directly integrated in the objective function. Such approach requires additional $k=\left\lceil\log _{2} u\right\rceil$ strongly-connected qubits. Thus, it dramatically reduces the size of the instances that can be tackled on the quantum hardware.

To deal with this issue, the work Yonaga et al., 2020 suggests to use a formulation with auxiliary continuous variables for slacks. The resulting objective function over binary and continuous variables is further optimized by a hybrid heuristic. This hybrid heuristic implements the alternating direction method of multipliers assisted by a quantum QUBO optimizer. In the work Yonaga et al., 2020 this method was illustrated on the Quadratic Knapsack application.

Motivated by the need for management of slack variables, the work Braine et al., 2021 presents a hybrid heuristic that handles mixed binary problems (MBO):

$$
\begin{equation*}
\min _{\substack{x \in\{0,1\}^{n} \\ y \in \mathbb{R}^{m}}} x^{t} A(y) x+b(y)^{T} x+c(y) \tag{6.4}
\end{equation*}
$$

where quadratic coefficients $A: \mathbb{R}^{m} \rightarrow \mathbb{R}^{n} \times \mathbb{R}^{n}$ and linear coefficients $b: \mathbb{R}^{m} \rightarrow \mathbb{R}^{n}$ depend on continuous variables.

The hybrid algorithm for MBO is a modified version of some variational quantum algorithm for QUBO (such as QAOA or VQE). For fixed values of $y \in \mathbb{R}^{m}$ quantum circuit is used to find good values for binary variables. The continuous variables $y \in \mathbb{R}^{m}$ are optimized in the same classical loop that optimizes the parameters of the quantum circuit ( $\boldsymbol{\beta}$ and $\boldsymbol{\gamma}$ for QAOA ). The heuristic iteratively modifies binary and continuous variables until the convergence condition is satisfied. In Braine et al., 2021 this hybrid method was applied to the transaction settlement problem.

A different approach presented in Ohzeki, 2020 is designed to address fully-connected Ising models whith the full connectivity that follows from the constraint penalization. This approach uses the Hubbard-Stratanovich transformation to replace the original quadratic terms by linear terms with stochastic coefficients. The values of coefficients are determined in an iterative procedure that evaluates the energy expectation value conditioned on the coefficient values by making calls to quantum annealer. In Yu and Nabil, 2021 the resulting hybrid procedure was applied to the bus charging scheduling problem. It was shown that this approach allows to use the D-Wave quantum annealer to handle much bigger instances than if the naive penalization of constraints is applied.

## Warm-starting quantum algorithms

Hybridization approaches presented above extend the applicability of quantum algorithms. Alternatively, the integration of classical routines can be used to enhance the performance of quantum heuristics. For instance, several warm-start techniques were proposed for QAOA on the Maximum Cut problem Egger et al., 2021, Tate et al., 2020. These techniques use specially-designed initial state $\left|\psi_{0}\right\rangle$ instead of the trivial uniform superposition. The preparation of the initial state is guided by the information computed classically. In Egger et al., 2021 two alternative options are presented. The first one involves the solution of the continuous relaxation of the quadratic program. The second approach employs the solution returned by the Goemans-Williamson algorithm. The Goemans-Williamson algorithm is a rounding of the SDP relaxation. Alternatively, Burer-Monteiro relaxation is used in the warm starting presented in Tate et al., 2020.

Finally, we can design hybrid methods where quantum routines are used to accelerate classical schemes. A version of quantum-assisted Branch \& Price was presented in [Svensson et al., 2021]. In this work quantum heuristics are used to approximately solve the master problem. In the next section we suggest a different original way to integrate quantum routines in the Branch \& Price. We believe that our hybrid approach can improve the runtime of the routine compared to the fully classical procedure.

### 6.1 Integer Linear Programs

A huge variety of real-world optimization problems can be modeled as integer linear programs. An Integer Linear Program (ILP) in general form is an optimization problem that is written as follows:

$$
\begin{align*}
&(I L P): \quad \min c^{T} x  \tag{6.5}\\
& A x=b  \tag{6.6}\\
& D x \leq d  \tag{6.7}\\
& x \in \mathbb{N}^{n} \tag{6.8}
\end{align*}
$$

where $A$ and $D$ are matrices with real values.
For instance, we can consider a delivery problem where a set of $N$ items of weights $\omega_{i}, i \in$ $[1, \ldots, N]$ have to be assigned to one of $M$ trucks of different capacities $C_{t}, t \in[1, \ldots, M]$. We use binary variables $x_{i}^{t}$ that indicate that the item $i$ is delivered by the truck $t$. Each truck, if used, generated the cost $c_{t}$. We denote by $y_{t}$ the variable indicating that the truck $t$ was selected for delivery. We wish to assure the delivery demands at lowest cost:

$$
\begin{array}{lr}
\min c^{T} \boldsymbol{y} & \\
& \sum_{t=1}^{M} x_{i}^{t}=1, \\
& i \in[1, \ldots, N] \\
\sum_{i=1}^{N} \omega_{i} x_{i}^{t} \leq C_{t} y_{t}, & t \in[1, \ldots, M]  \tag{6.12}\\
x_{i}^{t} \in\{0,1\}, y_{t} \in\{0,1\} & i \in[1, \ldots, N], t \in[1, \ldots, M]
\end{array}
$$

One conventional way to exactly solve an integer linear program is Branch $\mathcal{B}$ Bound Burke and Kendall, 2 At each iteration the Branch \& Bound algorithm considers different restrictions of the feasible
subspace. In order to evaluate if a particular subspace $S$ is promising, i.e. if it can contain the optimal solution, the procedure uses bounds. Bounds indicates the range $\left[f_{\min }, f_{\max }\right] \subset \mathbb{R}$ of the values that the linear function $f(x)=c^{T} x$ take on the elements of the subspace $S$. For the procedure to be practical, bounds have to be efficiently computable.

### 6.1.1 Linear relaxation

In most implementations bounds are obtained via the linear relaxation $(L P)$ of the integer linear program:

$$
\begin{align*}
(L P): \quad \min & c^{T} x  \tag{6.13}\\
& A x=b  \tag{6.14}\\
& D x \leq c  \tag{6.15}\\
& x_{i} \geq 0, \quad x_{i} \in \mathbb{R} \tag{6.16}
\end{align*}
$$

In the linear relaxation we optimize the objective function over continuous variables instead of integer ones. Contrary to the problem over integer variables that is NP-hard, the continuous relaxation can be efficiently solved. In practice the simplex method (even if exponential in the worst case) is particularly efficient Nash, 2000.

Obviously, the optimum value $\nu_{L P}$ of the linear relaxation is smaller or equal to the optimum $\nu_{I L P}$ of (ILP). The difference $g=\nu_{I L P}-\nu_{L P}$ is called the relaxation gap. The solution of the (LP) provides a lower bound for $\nu_{I L P}$. The tightness of the lower bound is characterized by the relaxation gap.

Linear relaxation is not the only efficient way to evaluate the bounds of (ILP). Furthermore, in many cases the linear relaxation provides very poor bounds. For instance, if we consider an (ILP) for the Maximum Independent Set problem on a complete graph $K_{n}$ :

$$
\begin{array}{rr}
\max & \sum_{i \in V} x_{i} \\
& \\
& x_{i}+x_{j} \leq 1,  \tag{6.19}\\
& x_{i} \in\{0,1\},
\end{array} \quad \forall i \in V, j \in V, i \neq j
$$

the linear relaxation has the value $n / 2$ (achieved in the point $x=[0.5, \ldots, 0.5]$ ) while the optimum of the integer problem is 1 .

Better bounds may be obtained via the Lagrangian approach where only a part of the constraints (coupling constraints) are relaxed Geoffrion, 1974. Alternatively, a polyhedral approach (also called Branch \& Cut) improves the linear relaxation by iterative addition of new constraints that eliminate non-integral solutions Hoffman and Padberg, 1985. A specification of the polyhedral approach for mixed binary programs is given in Balas et al., 1993.

### 6.1.2 Branch \& Bound

Branch \& Bound uses a tree representation to implicitly explore the entire solution space $S=$ $\left\{x \mid A x=b, D x \leq d, x \in \mathbb{N}^{n}\right\}$. Each node $n$ of the tree corresponds to a subspace $S_{n} \subset S$. In the traditional Branch \& Bound subspaces are defined by additional linear constraints:

$$
\begin{align*}
S_{n}: & x \in S  \tag{6.20}\\
& A_{n} x=b_{n} \tag{6.21}
\end{align*}
$$

$$
\begin{equation*}
D_{n} x \leq d_{n} \tag{6.22}
\end{equation*}
$$

The root of the tree represents the entire space of feasible solutions $S \subset \mathbb{N}^{n}$. The subspace $S_{c}$ corresponding to a child node $c$ is contained in the subspace $S_{p}$ of the parent node $p$. Moreover, for each two children $c_{1}$ and $c_{2}$ of the same parent node the subspaces $S_{c_{1}}$ and $S_{c_{2}}$ have an empty intersection. In brief, for each node parent $p$ we have:

$$
\begin{equation*}
S_{p}=S_{c_{1}} \sqcup S_{c_{2}} \tag{6.24}
\end{equation*}
$$

where $c_{1}$ and $c_{2}$ are children of the node $p$.
At each node $n$ of the tree the algorithm considers the optimization problem on a restricted solution space:

$$
\begin{align*}
& \left(\operatorname{ILP}_{n}\right): \quad \min c^{T} x  \tag{6.25}\\
& \quad x \in S_{n} \tag{6.26}
\end{align*}
$$

A lower bound $l_{n}$ for the optimum of the integer linear program $I L P_{n}$ can be obtained from the linear relaxation. On the other side, when the algorithm treats a node the best known feasible solution $x^{*} \in S$ provides an upper bound $u$ on the value of the optimal solutions. Crucially, if $u \leq l_{n}$ further exploration of the node can't lead to an improvement in the value of the objective function. In such case the node is pruned, i.e. eliminated from the succeeding consideration.

Conversely, if $l_{n}<u$ the subspace $S_{n}$ may contain a solution that is better than $x^{*}$. If, in addition, the solution $x_{n}^{L P}$ of the linear relaxation is integral (or, equivalently, $x_{n}^{L P} \in S_{n}$ ) it is precisely the solution of the integer linear problem $\left(\operatorname{ILP}_{n}\right)$. In such case we update the upper bound to $u=l_{n}$ and the best known solution to $x^{*}=x_{n}^{L P}$.

The linear relaxation in a node $n$ is an optimization problem (LP-n) over a continuous set $L_{n} \subset \mathbb{R}^{n}$. If the node can't be pruned and the solution $x_{n}^{L P}$ of (LP-n) contains non-integral elements we have to further explore the subspace $S_{n}$.

Here comes the branching.
Essentially, in the branching step we expand the tree by adding two children $c_{1}$ and $c_{2}$ to the node $n$. We use additional linear constraints to partition the subspace $S_{n}=S_{c_{1}} \sqcup S_{c_{2}}$ in subsets

$$
\begin{align*}
S_{c_{i}}: & x \in S_{n}  \tag{6.27}\\
& P_{c_{i}}(x)=1 \tag{6.28}
\end{align*}
$$

where $P_{c_{i}}: \mathbb{R}^{n} \rightarrow \mathbb{B}$ is some predicate. Traditionally, the predicate $P_{c_{i}}$ is an inequality constraint.
The intention behind the branching is to eliminate the non-integral solution $x_{n}^{L P}$. More precisely, we chose the predicates $P_{c_{1}}$ and $P_{c_{2}}$ in a way that linear relaxations $L_{c_{1}}$ and $L_{c_{2}}$ of the sets $S_{c_{1}}$ and $S_{c_{2}}$ don't contain the partial value $x_{n}^{L P}$, i.e.

$$
\begin{equation*}
L_{c_{1}} \sqcup L_{c_{2}} \subsetneq L_{n} \text { while } S_{c_{1}} \sqcup S_{c_{2}} \equiv S_{n} \tag{6.29}
\end{equation*}
$$

Traditionally, for problems over binary variables the branching divides the solution space on subspaces corresponding to $x_{i}=0$ and $x_{i}=1$ for some variable $x_{i}$ that has a fractional value in the relaxed solution $x_{n}^{L P}$. If a general integer variable $x_{i} \in \mathbb{N}$ gets the fractional value $\tilde{x}_{i}$, the conventional brunching approach creates the nodes corresponding to $x_{i} \leq\left\lfloor\tilde{x}_{i}\right\rfloor$ and $x_{i} \geq\left\lceil\tilde{x}_{i}\right\rceil$.

The Branch \& Bound tree exploration is illustrated on the figure 6.1.
We highlight that the good quality bounds are essential for an efficient pruning. The pruning is extremely important as the runtime of the algorithm heavily depends on the ability to eliminate suboptimal nodes.


Figure 6.1: Branch and bound tree. At each node $n$ one of three situations happens. If $u \leq l_{n}$ the node is pruned. Otherwise, if the solution of the relaxed problem $x_{n}^{L P}$ is integral and better than the actual best solution the upper bound is updated: $x^{*}=x_{n}^{L P}, u=l_{n}$. Finally, if $l_{n}<u$ and $x_{n}^{L P}$ has non-integral entries the node is branched.

### 6.2 Classical Branch \& Price

We aim to address integer linear programs with a hybrid quantum-classical algorithm. For this purpose, we suggest a hybrid algorithm with a classical part that essentially reproduces the Branch \& Price technique developed in Vanderbeck, 1994. Branch \& Price is well-suited to address ILP with many variables. In Branch \& Price a huge number of variables in tackled with the column generation approach that allows to consider only a restricted set of variables at each moment.

### 6.2.1 Dantzig-Wolfe decomposition

Column generation was originally introduced in Ford and Fulkerson, 1958 to solve large linear programs over continuous variables arising in relaxations of the path-flow formulation of the network-flow problem. Column generation is particularly useful for (LP) instances when only few constraints in the standard formulation act on many variables. These constraints are referred as coupling. In what follows we assume that in the formulation (LP) the coupling constrains are given by the equalities $A x=b$.

The (LP) instances are well-suited for the column generation if most constraints ( $D x \leq d$ ) act only on small subsets of variables. If, in addition, the matrix $D$ has a block structure:

$$
D=\left(\begin{array}{lll}
D_{1} & &  \tag{6.30}\\
& \ddots & \\
& & D_{k}
\end{array}\right)
$$

the problem is particularly adapted for the decomposition. We denote by $m_{i}$ and $l_{i}$ the number of columns and rows respectively in each $D_{i}$.

To reason about the decomposable problem we introduce the notations:

$$
x=\left(\begin{array}{c}
x_{1}  \tag{6.31}\\
\vdots \\
x_{k}
\end{array}\right), \quad c=\left(\begin{array}{c}
c_{1} \\
\vdots \\
c_{k}
\end{array}\right), \quad d=\left(\begin{array}{c}
d_{1} \\
\vdots \\
d_{k}
\end{array}\right), \quad A=\left(\begin{array}{lll}
A_{1} & \ldots & A_{k}
\end{array}\right)
$$

where $x_{i}$ and $c_{i}$ are columns with $m_{i}$ elements, $d_{i} \in \mathbb{R}^{l_{i}}$ and $A_{i} \in \mathbb{R}^{m_{i} \times l_{i}}$.

The decomposition principle Dantzig and Wolfe, 1960 relies on the fact that every point in the bounded convex space $S=\left\{x \in \mathbb{R}^{m} \mid D x \leq d\right\}$ can be represented as a convex combination of its extreme points $\left\{p^{q}\right\}_{q \in Q}$ :

$$
\begin{equation*}
x=\sum_{q \in Q} \lambda^{q} p^{q}, \quad \sum_{q \in Q} \lambda_{q}=1 \tag{6.32}
\end{equation*}
$$

Representing the vectors $x_{i} \in\left\{y \in \mathbb{R}^{m_{i}} \mid D_{i} y \leq d_{i}\right\}$ by convex combinations we can obtain an equivalent formulation for (LP) that is called the master program:

$$
\begin{array}{rlr}
\text { (M-LP) : } \min & \sum_{i=1}^{k} c_{i}^{T}\left(\sum_{q \in Q_{i}} \lambda_{i}^{q} p_{i}^{q}\right) & \\
& \sum_{i=1}^{k} A_{i}\left(\sum_{q \in Q_{i}} \lambda_{i}^{q} p_{i}^{q}\right)=b & \\
& \sum_{q \in Q_{i}} \lambda_{i}^{q}=1, & \forall i \in[1, \ldots, k] \\
& \lambda_{i}^{q} \geq 0, & \forall i \in[1, \ldots, k], \quad q \in Q_{i} \tag{6.36}
\end{array}
$$

where $\left\{p_{i}^{q}\right\}_{q \in Q_{i}}$ are extreme points of the sets $S_{i}=\left\{y \mid D_{i} y \leq d_{i}\right\}$.
If we denote $c_{i}^{q}=c_{i}^{T} p_{i}^{q}$ and $a_{i}^{q}=A_{i} p_{i}^{q}$ the master program can be written as:

$$
\begin{array}{lr}
\min & \sum_{\substack{i \in[1, \ldots, k] \\
q \in Q_{i}}} c_{i}^{q} \lambda_{i}^{q} \\
\sum_{\substack{i \in[1, \ldots, k] \\
q \in Q_{i}}} a_{i}^{q} \lambda_{i}^{q}=b & \\
\sum_{\sum_{q \in Q_{i}}} \lambda_{i}^{q}=1, & \forall i \in[1, \ldots, k] \\
\lambda_{i}^{q} \geq 0, & \forall i \in[1, \ldots, k], q \in Q_{i} \tag{6.40}
\end{array}
$$

The number of variables in (M-LP) is $N=\left|Q_{1}\right|+\cdots+\left|Q_{k}\right|$. In practice $N$ is very large. We highlight that in the approach that we present in the following the huge number of variables is not a problem as we never explicitly consider all variables together. Instead, the set of variables is given implicitly by the constraints $D x \leq d$. To resume, using the decomposition we obtained for a linear program (LP) an equivalent formulation (M-LP) that has less constraints but much more variables.

### 6.2.2 Decomposition for integer linear programs

Now let's consider the ILP problem:

$$
\begin{array}{ll}
(\mathrm{ILP}) \min & c^{T} x \\
& A x=b \\
& D x \leq d \\
& x \in \mathbb{N}^{n} \tag{6.44}
\end{array}
$$

where (6.42) are the coupling constraints and the matrix $D$ has a block-diagonal structure (6.30).
A bounded set $S_{i}=\left\{y \in \mathbb{N}^{m_{i}} \mid D x \leq d\right\}$ of integer points can be explicitly enumerated: $S_{i}=\left\{p_{q}\right\}_{q \in Q_{i}}$. Each part $x_{i}$ of an integer vector $x=\left(\begin{array}{c}x_{1} \\ \vdots \\ x_{k}\end{array}\right)$ can be trivially written as

$$
\begin{equation*}
x_{i}=\sum_{q \in Q_{i}} \lambda_{i}^{q} p_{i}^{q} \text { where } \sum_{q \in Q_{i}} \lambda_{i}^{q}=1, \lambda_{i} \in\{0,1\}^{\left|Q_{i}\right|} \tag{6.45}
\end{equation*}
$$

Similarly to the continuous problem, we use decomposition to rewrite the (ILP) formulation in an equivalent master integer program (M-ILP):

$$
\begin{array}{cc}
\text { (M-ILP) : } \quad \min & \sum_{\substack{i \in[1, \ldots, k] \\
q \in Q_{i}}} c_{i}^{q} \lambda_{i}^{q} \\
& \sum_{\substack{i \in[1, \ldots, k] \\
q \in Q_{i}}} a_{i}^{q} \lambda_{i}^{q}=b \\
& \sum_{\substack{q \in Q_{i}}} \lambda_{i}^{q}=1, \\
& \forall i \in[1, \ldots, k]  \tag{6.49}\\
\lambda_{i}^{q} \in\{0,1\}, & \forall i \in[1, \ldots, k], q \in Q_{i}
\end{array}
$$

where the difference with (M-LP) is the integrality condition 6.49. We also remark that for continuous case the set $\left\{p_{q}\right\}_{q \in Q}$ is made out of extreme points of the polyhedron $\left\{y \in \mathbb{R}^{m} \mid D y \leq\right.$ $d\}$ while for ILP the set $\left\{p_{q}\right\}_{q \in Q}$ represents all points in $\left\{y \in \mathbb{N}^{m} \mid D y \leq d\right\}$.

We highlight that the formulations (ILP) and (M-ILP) have the same integral solution value. However, their linear relaxations in general are not equivalent. Indeed, the very benefit of considering the extensive formulation (M-ILP) is due to the fact that sometimes the linear relaxation of (M-ILP) is tighter. In other words, if $\nu^{*}$ denotes the optimal value of (ILP) and, equivalently, (M-ILP) while the respective linear relaxations have the values $\nu_{I L P}$ and $\nu_{M-I L P}$ the following inequalities hold:

$$
\begin{equation*}
\nu_{I L P} \geq \nu_{M-I L P} \geq \nu^{*} \tag{6.50}
\end{equation*}
$$

When an integer program is solved with Branch \& Bound the tightness of the bounds is extremely important for an efficient elimination of suboptimal subspaces. Therefore, decomposition may lead to a more tractable formulation.

### 6.2.3 ILP formulations with large number of variables

Most optimization problems admit multiple ILP formulations. For instance, the graph coloring problem on $G=(V, E)$ can be naturally formulated using $K+|V| K$ binary variables:

$$
\begin{array}{ll}
\text { (GC-c) } \min & \sum_{k=1}^{K} y_{k} \\
& \\
\sum_{k=1}^{K} x_{i}^{k}=1, & \forall i \in V  \tag{6.53}\\
x_{i}^{k}+x_{j}^{k} \leq 1, & \forall(i, j) \in E
\end{array}
$$

$$
\begin{array}{ll}
x_{i}^{k} \leq y_{k}, & \forall i \in V, k \in[1, \ldots, K] \\
x_{i}^{k} \in\{0,1\}, y_{k} \in\{0,1\}, & \forall i \in V, k \in[1, \ldots, K]
\end{array}
$$

where $K$ is some upper bound on the chromatic number $\chi(G)$. The variable $x_{i}^{k}$ is equal to 1 if the node $i$ is assigned a color $k$. The variable $y_{k}$ indicates if the color $k$ is used. The constraint 6.52 requires each node to be colored and 6.53) assures that the coloring is proper.

Such formulation has multiple disadvantages that make it difficult for the Branch \& Bound. Firstly, the linear relaxation provides a very poor bound $l=1$ (achieved by the partial solution $\left.x_{i}^{k}=1 / K, y_{i}=1 / K\right)$. Secondly, the branching is not able to eliminate symmetric solutions that can be obtained with a simple permutation of colors. Therefore, the branching rarely improves the relaxed solution.

Alternatively, we can represent a proper coloring as a partition of nodes in independent sets: $V=S_{1} \sqcup \cdots \sqcup S_{k}$. If we take one binary variable $y_{s}$ per independent set $s \in I(G)$ the optimization problem can be written as follows:

$$
\begin{align*}
(\mathrm{GC}-\mathrm{s}) \quad \min & \sum_{s \in I(G)} y_{s}  \tag{6.56}\\
& \sum_{\substack{s \in I(G) \\
v \in s}} y_{s} \geq 1, \quad \forall v \in V  \tag{6.57}\\
& y_{s} \in\{0,1\}, \quad \forall s \in I(G) \tag{6.58}
\end{align*}
$$

The constraint 6.57 ensures that each node is contained in at least one independent set selected for the resulting partition. We remark that in general case the number of independent sets $|I(G)|$ in a graph is exponential on $|V|$.

A variety of set covering problems can also be formulated as ILP with many variables Vanderbeck, 1994. In set covering problems we have to select an ensemble of subsets of a set of ground elements. Depending on the particular problem, special constraints allow to decide if a subset is appropriate for selection. The objective is to find a ensemble of valid subsets that cover each element of the ground set.

Interestingly, certain QUBO instances can be also encoded in ILP formulations suitable for decomposition. For such instances a procedure solving QUBO with column generation was presented in Mauri and Lorena, 2012.

### 6.2.4 Column generation

In previous sections we described different situations that lead to (ILP) with large amount of variables. Motivated by the terminology used in Dantzig-Wolfe decomposition we refer to such formulations as master programs (MP). In general, a master program is written in canonical form as follows:

$$
\begin{align*}
(\mathrm{MP}): \quad \min & \sum_{y \in S} c_{y} \lambda_{y}  \tag{6.59}\\
& \sum_{y \in S} a_{y} \lambda_{y} \geq b  \tag{6.60}\\
& \lambda_{y} \in\{0,1\}, \quad \forall \lambda_{y} \in S \tag{6.61}
\end{align*}
$$

where the set of variables $S$ is large. In this section we consider an approach that is specifically designed for such type of formulations. We remark that this approach remains tractable as it never considers a whole set of variables as once.

In practical applications the variables $S$ are implicitly given by a compact formulation.

Example: In the set covering formulation for the graph coloring (GC-s) problem the set of variables $S$ is the ensemble of stables $I(G)$.

If the master program is a result of decomposition, there is in $S$ one variable per every point of the bounded set $\left\{y \in \mathbb{N}^{n} \mid D y \leq d\right\}$. The compact formulation also indicates the values for the coefficients $c_{y} \in \mathbb{R}$ and $a_{y} \in \mathbb{R}^{n}$.

A linear relaxation of the master problem (LMP) is

$$
\begin{align*}
(\mathrm{LMP}): \quad \min & \sum_{y \in S} c_{y} \lambda_{y}  \tag{6.62}\\
& \sum_{y \in S} a_{y} \lambda_{y} \geq b  \tag{6.63}\\
& \lambda_{y} \leq 1, \quad \forall y \in S  \tag{6.64}\\
& \lambda_{y} \geq 0, \quad \forall y \in S \tag{6.65}
\end{align*}
$$

## Dual problem

Each bound $\lambda_{y} \leq 1$ 6.64 corresponds to one dual variable $\pi_{y}$ in the dual problem for (LMP). We use the notation $\boldsymbol{\pi}$ for the vector of $m$ dual variables corresponding to coupling equality constraints 6.63). The dual problem (D-LMP) for (LMP) is formulated as:

$$
\begin{align*}
& (\mathrm{D}-\mathrm{LMP}): \quad \max b^{T} \boldsymbol{\pi}-\sum_{y \in S} \pi_{y}  \tag{6.66}\\
& a_{y}^{T} \boldsymbol{\pi}-\pi_{y} \leq c_{y}, \quad \forall y \in S  \tag{6.67}\\
&  \tag{6.68}\\
& \pi_{y} \geq 0, \quad \forall y \in S  \tag{6.69}\\
& \\
& \boldsymbol{\pi} \geq \mathbf{0}, \boldsymbol{\pi} \in \mathbb{R}^{m}
\end{align*}
$$

We refer to the appendix A. 1 for for more details about the duality for linear programs.
If both primal and dual problems are bounded, the strong duality theorem states that their optimum values are equal Gass and Harris, 2001. Moreover, a feasible solution of the primal problem $\boldsymbol{\lambda}^{*}$ correspond to a feasible solution $\boldsymbol{\pi}^{*}$ of the dual problem if and only if the primal solution is optimal.

When the set of variables $S$ is too large, the master problem can't be directly solved by conventional methods for ILP. Indeed, due to an exponential amount of variables even the linear relaxation is intractable. To tackle this issue, the work Vanderbeck, 1994 adopts the column generation approach originally introduced for the cutting stock problem Gilmore and Gomory, 1963.

## Simplex method for large formulations

Column generation allows to solve the linear relaxation of the master problem (LMP) by considering a restricted set of variables $S^{\prime}$. New variables are added to the set only if they can
improve the solution. Promising variables are selected using the dual values of the restricted linear relaxation.

The procedure is inspired by the simplex method Nash, 2000 that we briefly present in the appendix A.2. We remark that primal and dual problems are simultaneously solved by the simplex method.

The simplex method operates a set $B$ of basis variables of cardinality $m$ where $m$ is the number of inequality constraints in the canonical formulation. Remaining variables $F$ are called free variables. At each iteration the method consider a basic feasible solution that is obtained by setting all free variables to zero. If a non-zero value of a free variable $\lambda_{i} \in F$ improves the value of the objective function, the basis is updated with a pivot operation. The pivot operation modifies the basis $B \rightarrow B^{\prime}$ by replacing a carefully selected leaving variable $\lambda_{e} \in B$ with the promising free variable.

Successive basis modifications drives the method to a basic feasible solution that is optimal for (LP). Crucially, in an optimum obtained that way only $m$ basis variables have non-zero values. Motivated by this observation, column generation adds the variables in a formulation only if they may appear in the optimum basis.

## Pricing subproblem

At each iteration $t$ column generation solves a restricted relaxation of master problem ( t -LMP):

$$
\begin{array}{ccc}
(\mathrm{t}-\mathrm{LMP}): \quad \min & \sum_{y \in S^{t}} c_{y} \lambda_{y} & \\
& \sum_{y \in S^{t}} a_{y} \lambda_{y} \geq b & \\
& \lambda_{y} \leq 1, & \forall y \in S^{t} \\
& \lambda_{y} \geq 0, & \forall y \in S^{t} \tag{6.73}
\end{array}
$$

where $S^{t} \subset S$. We denote by $y^{*}$ the optimal solution for (t-LMP) and by $\left(\pi_{y}^{*}, \pi^{*}\right)$ the corresponding values of the dual variables.

An optimal solution for the restricted problem (t-LMP) can be extended to a feasible solution for the initial (LMP) by setting to zero all variables $\lambda_{y}$ for $y \in S \backslash S^{t}$. Similarly, we may obtain an assignment (not necessarily feasible) for the values of dual variables for (LMP). In the assignment for dual variables we take $\pi_{y}=0, \forall y \in S \backslash S^{t}$. We can check if the extended primal solution is optimal for (LMP) by verifying if the extended dual solution is feasible for (DMP) (i.e. it satisfies the constraints 6.67).

For bounded problems the dual solution is infeasible if and only if there exists at least one primal variable $y^{\prime} \in S \backslash S^{t}$ such that the inequality 6.67 is violated, i.e. $a_{y^{\prime}}^{T} \pi^{*}>c_{y^{\prime}}$. Such variable can be found by solving the problem:

$$
\begin{gather*}
\max a^{T}(y) \pi^{*}-c(y)  \tag{6.74}\\
y \in S \backslash S^{t} \tag{6.75}
\end{gather*}
$$

called pricing subproblem $(P S)$. For each variable $y \in S$ we refer to the value $r_{y}=a_{y}^{T} \pi^{*}-c_{y}$ as reduced cost.

When the optimum solution $y^{\prime}$ of the pricing subproblem has a value that is smaller or equal to zero the extended dual solution is feasible. Therefore, the optimum solution for the
restricted problem ( t -LMP) is the same as the optimum of the original relaxation (LMP) over the entire set of variables. Otherwise, the variable found by (PS) can improve the value of the objective function. In such case column generation extends the set $S^{t}$ with the new variable: $S^{t+1}=S^{t} \cup\left\{y^{\prime}\right\}$ and the process restarts. The classical column generation is illustrated on the green part of the figure (6.2).


Figure 6.2: The quantum-classical column generation for huge linear relaxations. The relaxation is solved at each node of the Branch \& Bound tree. The green nodes corresponds to the classical column generation and the red ones - to integrated quantum routines.

### 6.2.5 Branch \& Price for Integer Linear Programs

Column generation is an approach that handles large linear optimization problems over continuous variables. The Branch \& Price technique originally presented in Vanderbeck, 1994 integrates column generation in Branch \& Bound algorithm to derive an exact algorithm for large integer linear programs.

Branch \& Price is a method for the formulations over many integer variables. Such formulations appear either naturally or as a result of decomposition of some compact (ILP) formulation. Just as Branch \& Bound, Branch \& Price implicitly explores the entire solution space in a tree search fashion. It uses the linear relaxation to bound the values of the objective function in each node of the tree. Nodes that can't lead to an improved solution are pruned from the tree.

The major distinction of the Branch \& Price from the Branch \& Bound consists in the way
they manage variables. During the entire execution Branch \& Bound explicitly consider all variables of the problem. Branch \& Price, on the contrary, operates a large set $S$ of implicitly given variables. Generally, $S$ is represented by a set on inequalities $D p \leq d$ on elements $p$ of a vector space with some finite dimension $Q$. Each integer point $p \in \mathbb{N}^{Q}$ satisfying the inequalities corresponds to one variable $y_{p} \in S$. Coefficients $c_{p}$ and $a_{p}$ for the variable $y_{p}$ are usually given by some (often linear) functions $c: \mathbb{N}^{Q} \rightarrow \mathbb{R}$ and $a_{p}: \mathbb{N}^{Q} \rightarrow \mathbb{R}^{m}$ evaluated in the corresponding point $p$.

Example: An ensemble of independent sets $I(G)$ of a graph $G=(V, E)$ corresponds to points $x \in\{0,1\}^{|V|}$ satisfying the inequalities $x_{u}+x_{v} \leq 1$ for every edge $(u, v) \in E$. In the set covering formulation for the graph coloring problem cost function coefficients for every variable are equal to one: $c \equiv 1$. The coefficients in covering constraints $a_{x}$ are given by the characteristic vector $x$ of the independent set.

In each node $n$ of the branching tree the method solves the linear relaxation of some (ILP) problem over variables $S$. It begins by considering a restricted set of variables $S^{0}$ such that the relaxed problem admits a feasible solution on $S^{0}$. The set of variables is iteratively extended $S^{t} \rightarrow S^{t+1}=S^{t} \cup\left\{y_{p}\right\}$ by the column generation that solves the pricing subproblem:

$$
\begin{align*}
& \text { Find } p  \tag{6.76}\\
& \qquad \begin{array}{l}
a_{p}^{T} \boldsymbol{\pi}_{t}^{*}>c_{p} \\
D p \leq d \\
\\
p \in \mathbb{N}^{Q}
\end{array} \tag{6.77}
\end{align*}
$$

where $a_{p}=a(p), c_{p}=c(p)$ and $\boldsymbol{\pi}_{t}^{*}$ is the dual solution of the restricted linear program.
The variable generation continues until the linear relaxation is solved to optimality. Equivalently, the generation stops when the pricing subproblem becomes infeasible. Thereafter if the relaxed solution is non-integral and the bound doesn't permit pruning the branching occurs (see figure 6.3).

## Ryan-Foster branching rule

Due to the additional variables that join the formulation in the column generation process, traditional branching techniques are not adapted for the Branch \& Price. Indeed, if we create two branches corresponding to $\lambda_{y}=0$ and $\lambda_{y}=1$ it is unclear how to ensure that the pricing problem doesn't generate the same variable again. Moreover, in many cases the constraint $\lambda_{y}=0$ is very weak. For instance, in set covering problems it eliminates only one subset from an exponential number. Therefore, more sophisticated techniques are required that take in account the structure of the pricing problem.

When the master problem is a variant of the set covering problem we can use the Ryan-Foster rule Ryan and Foster, 1981. The rule creates two branches. In the first branch we require two elements $x_{1}$ and $x_{2}$ from the ground set to be covered by the same subset. In the second branch the opposite condition is imposed: valid subset can't simultaneously contain both variables.

Example: For the graph coloring problem the Ryan-Foster rule admits a natural interpretation. In this problem the ground set is the set of nodes $V$ and variables correspond to independent sets. The rule firstly selects two nodes $v$ and $u$ that are not connected in the graph $(u, v) \notin E$. In the first branch the exploration continues on the variables $\lambda_{y}$ that assign $u$ and $v$ to the same


Figure 6.3: Branch \& Price. The column generation stops when the pricing subproblem becomes infeasible. Then the method proceeds as traditional Branch \& Bound.
color $(u \in y \Longleftrightarrow v \in y)$. This constraint may be communicated to the pricing subproblem by "merging" the nodes $v$ and $u$ in the graph. In the second branch the separation is enforced by addition of the new edge $(u, v)$ to the graph.

### 6.3 Quantum-assisted Branch \& Price

### 6.3.1 Embedding (R)QAOA into the Branch \& Price

A hybrid method presented in Svensson et al., 2021 use QAOA as an heuristic solver for the master problem when it corresponds to a set covering problem.

Instead, we suggest a different approach that integrates quantum heuristics in the pricing step of the column generation (see fig. 6.2). From the dual solution $\boldsymbol{\pi}$, the pricing problem (PS) is typically formulated as:

$$
\begin{gather*}
\max \boldsymbol{\pi}^{T} A p-c_{p}  \tag{6.80}\\
D p \leq d  \tag{6.81}\\
p \in \mathbb{N}^{Q} \tag{6.82}
\end{gather*}
$$

Such formulation is an integer linear program and, thus, NP-hard in general. In some cases (PS) has a particular structure that make it tractable for polynomial algorithms. For instance, in
the path-flow formulation of the network flow problem the optimum for the pricing subproblem corresponds to the shortest path in a graph Desrosiers and Lübbecke, 2006. More generally, if (PS) satisfy the integrality property $\square^{7}$ it can be efficiently solved with the simplex method.

However, in most cases (PS) is an NP-hard problem.

Example: in the Branch \& Price for graph coloring the pricing subproblem searches for an improving independent set. Each independent set $s$ has a cost function coefficient $c_{s}=1$ and it contributes to a covering constraint for every node $v \in s$. We denote by $x \in\{0,1\}^{|V|}$ the incidence vector of a subset of $V$. Then (PS) can be formulates as follows:

$$
\begin{align*}
\max & \sum_{v \in V} \pi_{v} x_{v}  \tag{6.83}\\
& x_{u}+x_{v} \leq 1, \quad(u, v) \in E  \tag{6.84}\\
& x \in\{0,1\}^{|V|} \tag{6.85}
\end{align*}
$$

which exactly corresponds to NP-hard maximum weighted independent set problem on the graph $G=(V, E)$ with node weights $w_{v}=\pi_{v}$.

When (PS) is NP-hard, finding an exact solution may take a long time. However, in Branch \& Price we don't need precisely the exact solution but rather an answer to the question if there exists an improving variable. A variable $y_{p}$ improves the solutions of the relaxation if it violates dual feasibility, i.e. if $\boldsymbol{\pi}^{T} A p-c_{p}>0$. Therefore, if an approximate solution $\hat{p}$ for (PS) has the value $\nu_{\hat{p}}=\pi^{T} A \hat{p}-c_{\hat{p}}>0$ we can accept it as a solution for (PS). For hard problems a sufficiently good heuristic algorithm can rapidly generate such variables.

On the other side, the algorithm ceases to generate columns only when there is no more improving variables. To verify if this termination condition is satisfied, one has to assure that the exact optimum value $\nu^{*}$ of (PS) is lower or equal to zero. Therefore, at each node of the branching tree we have to launch the exact algorithm for (PS) at least once to confirm that the relaxed problem is solved to optimality.

Taking in account previous remarks, we suggest a quantum-classical approach to efficiently solve the pricing subproblem. Our method proceeds in three steps. Firstly, it tries to find a solution with a classical heuristic. If the value of the obtained approximate solution is $\nu_{c}>0$ we declare the (PS) solved and add the corresponding variable to the formulation. Otherwise, we launch a quantum heuristic. If both heuristics fail, en exact algorithm is called. The exact algorithm either finds an improving variable or guarantees the optimality of the relaxed solution.

For the graph coloring example the resulting hybrid algorithm for the pricing subproblem is sketched on Figure 6.4. A scheme for the overall Quantum-assisted Branch \& Price technique is given on Figure 6.2.

### 6.4 Graph Coloring with Quantum-assisted Branch \& Price

We illustrate our hybrid approach on the graph coloring problem.
The graph coloring problem searches for a proper vertex coloring with minimal number of colors. The optimal value of the problem is called chromatic number $\chi(G)$ of the graph. In an extensive formulation of the coloring problem we aim to minimize the number of independent sets required to cover every vertex.

[^4]Finding a minimal graph coloring is NP-hard Garey and Johnson, 1990. The decision version of the problem called $k$-colorability belongs to Karp's 21 NP-complete problems Karp, 1972. For $k$-colorability the asymptoticly best exact algorithm based on dynamic programming has the complexity $O\left(2.4423^{n}\right)$ Lawler, 1976.

In addition, minimal graph coloring is NP-hard to approximate within a factor $n^{1-\epsilon}$ for all $\epsilon>0$ [Zuckerman, 2007. The best known approximation algorithm achieves the approximation ratio of $O\left(n(\log \log n)^{2}(\log n)^{3}\right)$ Halldórsson, 1993.

Graph coloring is used to model multiple real-world problems from different domains such as scheduling, resource allocations and many others [Marx, 2003, Lewis, 2021]. Therefore, efficient algorithms for this problem are of extreme importance. Naive approaches are based on greedy ideas that color vertices following some order Brélaz, 1979]. We refer to [Husfeldt, 2015] for a complete review of diverse classical algorithms for the graph coloring.

Quantum heuristics can be directly applied to the problem if some QUBO formulation is provided. One such formulation was presented in Kochenberger et al., 2005. Interestingly, a generic tabu search algorithm for QUBO was able to find competitive solutions on this last formulation.

More refined quantum approach results from the Quantum Alternating Operator Ansatz framework. An implementation presented in Hadfield et al., 2019 encodes the color of each node in a sequence of $D_{G}+2$ qubits where $D_{G}$ is a maximum node degree in the graph. The quantum system is initialized in a trivial feasible solution that uses $D_{G}+1$ colors. Mixer operator is tailored to preserve the evolution is the subspace of proper colorings.

### 6.4.1 Column generation

An algorithm for the graph coloring based on column generation was originally presented in Mehrotra and Trick, 1996. This algorithm considers a set-covering formulation (GC-s)

$$
\begin{align*}
\text { (GC-s) : } \quad \min & \sum_{s \in I(G)} \lambda_{s}  \tag{6.86}\\
& \sum_{\substack{s \in I(G) \\
v \in s}} \lambda_{s} \geq 1,  \tag{6.87}\\
& \lambda_{s} \in\{0,1\}, \tag{6.88}
\end{align*}
$$

This formulation contains one variable per independent set. In general case, the number of variables in (GC-s) is exponential.

The optimum value of the problem is equal to the chromatic number of the graph $\chi(G)$. Adopting the notation from Held et al., 2012, we use $\chi_{f}(G)$ to refer to the optimum of the linear relaxation also called fractional chromatic number. Clearly, $\chi_{f}(G)$ provides a lower bound of the chromatic number. Although there exists efficient algorithms for linear programs over continuous variables, the value $\chi_{f}(G)$ is also NP-hard to approximate with a factor $n^{1-\epsilon}$ for all $n>\epsilon$ Zuckerman, 2007]. The complexity of the relaxed problem is consistent with the exponential amount of variables in the formulation.

Column generation is used to optimize the linear relaxation of (GC-s), i.e. to evaluate $\chi_{f}(G)$. It may be is further integrated in an exact coloring algorithm implementing the Branch \& Price approach.

The state-of-the-art program exactcolor solving the graph coloring problem with Branch \& Price was presented in Held et al., 2012. In this program the initial set of variables $S^{0}$ for the
restricted master problem is obtained by the DSatur heuristic Brélaz, 1979. In the column generation a generic LP-solver Gurobi Optimization, LLC, 2022 optimizes the linear relaxation over a restricted set of variables $S^{t}$. The LP-solver returns values for dual variables $\pi_{v}, v \in V$ that define the pricing subproblem (MWIS):

$$
\begin{array}{rlr}
\text { (MWIS) : } \quad \max & \sum_{v \in V} \pi_{v} x_{v} & \\
& x_{u}+x_{v} \leq 1, & (u, v) \in E \\
& x_{v} \in\{0,1\}, & v \in V
\end{array}
$$

The generation halts when the optimum independent set $s^{*}$ found by the pricing problem has weight $\pi\left(s^{*}\right)=\sum_{v \in s^{*}} \pi_{v} \leq 1$. If the pricing subproblem returns a solution with $\pi\left(s^{*}\right)>1$ the restricted set of variables $S^{t}$ is extended by $s^{*}$.

The exactcolor program solves the pricing subproblem with a two-step approach. Firstly, it search for an heuristic solution with a method that combines greedy initialization with 2interchange local search. If the heuristic fails to find an improving variable, an exact Branch \& Bound algorithm is used that either returns a promising variable or confirms that it doesn't exist.

If the linear relaxation has non-integral solution the node is branched according to the RyanFoster rule.

Numerical experiments confirming the interest of this approach were reported in Held et al., 2012. In one of these experiments column generation was used to evaluate the fractional chromatic number $\chi_{f}(G)$. On 119 out of 136 instances from the DIMACS benchmark Johnson and Trick, 1996 the program terminated in less then three days. The fractional chromatic number computed by the program improved previously known bounds for 5 instances.

### 6.4.2 A hybrid approach to the pricing subproblem

We suggests a hybrid method that adopts a combined quantum-classical approach for the pricing MWIS subproblem. Our motivation lies in good results that quantum heuristics were able to achieve on MWIS (see section 5.3.3).

The work [Svensson et al., 2021] suggests a different combination of quantum and classical routines in the Branch \& Price framework. In their approach column generation extends the set of variables until the linear relaxation is solved to optimality. The resulting variables form a set $S^{\prime} \subset S$. In the next step a quantum heuristic is used to computes an approximate integer solution for the restricted master problem over $S^{\prime}$. We remark that when the master problem can be interpreted as a set packing problem it admits a natural QUBO formulation Alidaee et al., 2008.

Contrary to our approach, this second technique doesn't naturally integrate in the traditional Branch \& Price framework. Indeed, even an exact integer optimum on the restricted set $S^{\prime}$ isn't necessary equal to the optimum over $S$. In other words, although the variables from $S \backslash S^{\prime}$ are null in some optimum of the relaxed problem, they can have positive values in the optimal integer solution. Therefore, the only information that we can obtain with a quantum heuristic applied to the restricted master problem is a new upped bound $u$ on the optimal solution. This information is practical if, for instance, $u$ matches the rounded relaxed optimum: $u=\left\lceil l_{S^{\prime}}\right\rceil \equiv\left\lceil l_{S}\right\rceil$ or if it improves the best known upper bound.

In addition, as the set $S^{\prime}$ grows when new variables are generated the restricted master problem can rapidly become intractable for the quantum hardware of limited size. We remark that in our approach quantum heuristics are always called on a fixed number of variables $|V|$.

In our hybrid technique the instance of MWIS problem that appears at the pricing step is tackled with a three-fold procedure that combines the classical heuristic from Held et al., 2012, RQAOA and the exact algorithm (see fig 6.4).


Figure 6.4: Hybrid algorithm for the pricing subproblem for graph coloring. In our implementation, we use the same classical heuristic and Branch \& Bound implementation as in Held et al., 2012 .

Different execution scenarios for the pricing step are described in the table 6.1.
We define the $R Q A O A$ improvement rate as a ratio of pricing calls in which the classical heuristic failed but the RQAOA succeeded (green line in the table 6.1) to the total amount of cases when there were an improving solution but the classical heuristic failed (green and red lines in the table 6.1).

Quantum algorithms are expected to improve over classical heuristics only in some special cases McClean et al., 2021. If the pricing encounters one of such instances, the use of quantum heuristic allows to postpone the necessity to execute a costly exact algorithm. Therefore, in the

| Classical heuristic | Quantum heuristic | Solution exists <br> (call exact algorithm) |
| :---: | :---: | :---: |
| Succeeded | - | True |
| Failed | Succeeded | True |
| Failed | Failed | True |
| Failed | Failed | False |

Table 6.1: Possible execution scenarios of the pricing step
quantum heuristic it sufficiently fast, the overall procedure is accelerated.

### 6.4.3 Numerical results

We evaluated the $R Q A O A$ improvement rate on 10 Erdos-Renyi random graph with 50 nodes.

| Density | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 | 0.6 | 0.7 | 0.8 | 0.9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Classical heuristic success rate | 0.95 | 0.84 | 0.76 | 0.92 | 0.89 | 0.95 | 0.96 | 1 | 1 |
| RQAOA improvement rate | 0.75 | 0.10 | 0.03 | 0.07 | 0.04 | 0 | 0.2 | - | - |

Table 6.2: Classical heuristic success rate (the fraction of improving variables found by the classical heuristic algorithm) and RQAOA improvement rate.

This table shows that it is indeed profitable for some instances to use the hybrid method instead of the traditional classical approach. We recall that the RQAOA is launched only if the classical heuristic fails to find the improving variable, i.e. only on rather difficult instances.

To test our approach in realistic settings we considered instances issued from a modified version of the Charge Task Selection problem described in chapter ??(5). In the new coloring version the conflict graph remains the same, while the new objective requires to cover all charges with a minimal number of stations. Previously (see section 5.3.3) we have reported that RQAOA sometimes finds better approximation that the classical heuristic for unweighted MIS on the instances issued from the Charge Task Selection problem.

Unfortunately, for the instances of the coloring version of the Charge Task Selection problem we've observed that there is no need for RQAOA as the classical heuristic always finds the solution of the pricing problem whenever it exists at all. A spectacular performance of the classical heuristic compared to our previous result is probably due to the presence of weights that have a special structure determined by the dual solution of the relaxation.

Variable results on different instances indicates that the performance enhancement resulting from to the integration of quantum heuristics is not robust across different families of instances.

### 6.5 Discussion

In this chapter we presented a hybrid quantum-assisted Branch \& Price method for large integer linear programs. In principle, it can be adapted to any formulation over huge set of variables that is implicitly given by another pricing integer program.

However, the additional cost related to the integration of quantum heuristics should be justified by two conditions:

- the computational difficulty of the pricing subproblem for classical algorithms
- amenability of the pricing subproblem to a QUBO formulation

This conditions are fulfilled by the set covering formulation of the graph coloring problem. Another example is the cutting stock problem for which the pricing subproblem is equivalent to $0-1$ knapsack problem Gilmore and Gomory, 1963. The $0-1$ knapsack problem was addressed with QAOA in de la Grand'rive and Hullo, 2019.

In addition, if the formulation is a result of decomposition for block-diagonal $D$ the size of resulting $k$ disconnected pricing subproblems is small compared to the size of the compact formulation. It may be that the compact formulation has too many variables but the subproblems are small enough to be handled on near-future quantum hardware.

## Part III

Variational algorithms with ZX calculus

## Chapter 7

## Introduction to ZX calculus

ZX-calculus, originally introduced in Coecke and Duncan, 2011, is a graphical language that allows to reason about quantum computing. In this language complex computations on qubits are represented with diagrams. Diagrams are made out of elementary generators. Each diagram corresponds to a linear transformation between qubit spaces. In a sense, a ZX-diagram is a tensor network representation of a linear map. A compact set of rewrite rules allows to transform diagrams into equivalent ones. The notable advantage of ZX-calculus compared to other representation (including linear maps, circuits and tensor networks) is that computations may be done entirely graphically. In other words, we can manipulate matrices of exponential size using local transformations on more compact diagrammatic representation.

The work van de Wetering, 2020 provides an excellent introduction to ZX-calculus together with a fairly detailed review of important results. This introduction assumes a general background in quantum computing. A much more exhaustive book Coecke and Kissinger, 2017 introduces the entire field of quantum computing directly in the language of ZX-diagrams. A Python package PyZX Kissinger and van de Wetering, 2020a provides automatic tools to reason about ZX-diagrams. It is particularly useful to study diagrams that are too large to be analyzed by hand. Many works on ZX-diagrams, including this thesis, use the Tikzit software ${ }^{8}$ to draw ZX-diagrams.

ZX-calculus was used to address many quantum applications. For instance, it was used to establish important results for the measurement-based quantum computing (MBQC) Duncan and Perdrix, 2010 Kissinger and van de Wetering, 2019. We recall that MBQC is an alternative model for quantum algorithms. In MBQC the computation is performed by properly designed measurements of a special quantum state. A detailed introduction to MBQC is available in Jozsa, 2005.

Other than MBQC, ZX-calculus was successfully applied for circuit optimization Cowtan et al., 2020, de Beaudrap et al., 2020, Kissinger and van de Wetering, 2020b, Duncan et al., 2020. The graphical language is also adapted for design and verification of quantum error correction codes Garvie and Duncan, 2018. Moreover, ZX-calculus is particularly well-suited for the analysis and compilation of surface codes Horsman, 2011, de Beaudrap and Horsman, 2020, Hanks et al., 2020. Surface codes are extremely promising for the fault-tolerant quantum computing.

In the context of combinatorial optimization an extension of ZX-calculus called ZH-calculus was used to reason about the respective complexities of constraint satisfaction problem (CSP) and \#CSP (the maximization of a number of satisfied clauses) de Beaudrap et al., 2021. In the work de Beaudrap et al., 2021 both problems CSP and \#CSP on the same input were represented by the same diagram, while with interpretation over different rings. A polynomial

[^5]complexity of 2-SAT is reflected in the simplicity of the graphical reduction of the corresponding ZH-diagram.

Another extension of ZX-calculus called ZXH-calculus was used to study physical properties of 1D AKLT state East et al., 2022. Interestingly, AKLT states are ground states of quantum many-body systems under a multidimensional Heisenberg spin model.

Finally, ZX-calculus also found an application in variational quantum algorithms. In particular, the barren plateau phenomena for parameterized quantum circuits, i.e. the exponential vanishing of the gradient with the number of qubits, can be conveniently analyzed with the graphical tools Zhao and Gao, 2021, Toumi et al., 2021. The work Fontana et al., 2020 use ZX-diagrams to express the symmetries in the parameter landscape.

### 7.1 ZX diagrams



Figure 7.1: Example of a ZX-diagram
As we can see on the figure 7.1, a ZX-diagram is a network made out of nodes and wires. There are three types of nodes: red spiders $\alpha$, green spiders a and Hadamard boxes Red and green spiders sometimes are drawn with a rational number inside. This number is called phase or angle. Both spiders can have an arbitrary number of adjacent wires. The Hadamard box has strictly two wires. A wire can connect two nodes or, alternatively, have free ends that point either up or down. Wires with free ends pointing towards the top or towards the bottom of the diagram are called inputs and outputs respectively. We denote by $D(n, m)$ or $D: n \rightarrow m$ a diagram with $n$ inputs and $m$ outputs. Diagrams with no inputs or no outputs are called states and effects respectively. If a diagram has no input and no output wires it is called a scalar.

In quantum computing ZX-diagrams are used as a graphical representation of linear maps. A diagram with $n$ inputs and $m$ outputs corresponds to a complex matrix with $2^{n}$ columns and $2^{m}$ rows. Using the terminology from category theory we call interpretation functor (or, alternatively, a standard interpretation) the map $\llbracket \rrbracket$ that associates to a ZX-diagram the corresponding matrix:

$$
\begin{equation*}
\llbracket D(n, m) \rrbracket=M_{2^{m} \times 2^{n}}(\mathbb{C}) \tag{7.1}
\end{equation*}
$$

### 7.1.1 Generators

Formally, ZX-diagrams are inductively defined from a set of basic generators combined with sequential and parallel compositions.

The generator $-0 \rightarrow 0$ is an empty ZX-diagram. It correspond to the scalar value 1 . For more complex ZX-diagrams the elementary building blocks are red and green spiders and the hadamard box.

## Spiders

The green spider also called $Z$-spider is a diagram with $n \in \mathbb{N}$ inputs and $m \in \mathbb{N}$ outputs with the interpretation:

$$
\begin{equation*}
\llbracket \overbrace{\underbrace{\overbrace{\cdots}^{\infty}}_{m}}^{n} \|=|\underbrace{0 \ldots 0}_{m}\rangle\langle\underbrace{0 \ldots 0}_{n}|+e^{i \alpha}|\underbrace{1 \ldots 1}_{m}\rangle\langle\underbrace{1 \ldots 1}_{n}| \tag{7.2}
\end{equation*}
$$

The spider is parameterized by a real number $\alpha \in \mathbb{R}$ called phase. As the parameter $\alpha$ is used in a complex power of an exponent, we can restrict the range to $\alpha \in[0,2 \pi]$. We remark that the linear map that corresponds to the Z-spider is a symmetric tensor.

The red spider is alternatively referenced as X-spider. Like the green spider, it can have an arbitrary number of inputs $n \in \mathbb{N}$ and outputs $m \in \mathbb{N}$. The red spider is also decorated by a real value $\alpha$ that we can take in the range $\alpha \in[0,2 \pi]$.

$$
\begin{equation*}
\|\overbrace{\underbrace{\alpha}_{m}}^{n}\|=|\underbrace{+\cdots+}_{m}\rangle\langle\underbrace{+\cdots+}_{n}|+e^{i \alpha}|\underbrace{-\cdots-}_{m}\rangle\langle\underbrace{-\cdots-}_{n}| \tag{7.3}
\end{equation*}
$$

In this work we use the popular convention that the phase $\alpha=0$ can be omitted.
Z-spider can be used to express some well-known states and matrices such as Z-rotation gate and vectors $|+\rangle$ and $|-\rangle$ of the Hadamard basis:

$$
\begin{align*}
& \llbracket \propto|\rrbracket=| 0\rangle\langle 0|+e^{i \alpha}|1\rangle\langle 1|=R_{Z}(\alpha)  \tag{7.4}\\
& \llbracket \bigcirc \rrbracket=|0\rangle+|1\rangle=\sqrt{2}|+\rangle  \tag{7.5}\\
& \llbracket \Pi \pi \rrbracket=|0\rangle-|1\rangle=\sqrt{2}|-\rangle \tag{7.6}
\end{align*}
$$

With X-spider we can draw the diagrams for the vectors from the computational basis $\{|0\rangle,|1\rangle\}$. Red spiders with one input and one output corresponds to the X-rotation with angle $\alpha$.

$$
\begin{align*}
& \llbracket \propto \downarrow \rrbracket]=|+\rangle\langle+|+e^{i \alpha}|-\rangle\langle-|=R_{Z}(\alpha)  \tag{7.7}\\
& \llbracket \bigcirc \rrbracket=|+\rangle+|-\rangle=\sqrt{2}|0\rangle  \tag{7.8}\\
& \llbracket \overparen{\pi} \rrbracket \rrbracket=|+\rangle-|-\rangle=\sqrt{2}|1\rangle \tag{7.9}
\end{align*}
$$

It follows from the definition that the no-input no-output spiders $\pi$ and $\pi$ correspond to the zero scalar.

## Hadamard box

The Hadamard gate $H=\frac{1}{\sqrt{2}}\left(\begin{array}{cc}1 & 1 \\ 1 & -1\end{array}\right)$ is represented by a yellow box diagram:

$$
\begin{equation*}
\llbracket \frac{1}{1} \rrbracket \rrbracket=|+\rangle\langle 0|+|-\rangle\langle 1| \tag{7.10}
\end{equation*}
$$

Later in this section we will see that the linear map $H$ can alternatively be represented by a combination of green and red spiders.

## Wire generators

A number of other useful linear transformations are expressed with wires. A simple wire with nothing on it corresponds to the identity:

$$
\begin{equation*}
\llbracket|\rrbracket=| 0\rangle\langle 0|+|1\rangle\langle 1| \tag{7.11}
\end{equation*}
$$

We remark that empty one-input one-output spiders corresponds to the same transformation:

$$
\begin{equation*}
\bigcirc=\bigcirc=1 \tag{7.12}
\end{equation*}
$$

In the expression 7.12 we use the equality notation $D_{1}=D_{2}$ for diagrams to tell that the diagrams have the same semantics:

$$
\begin{equation*}
D_{1}=D_{2} \Longleftrightarrow \llbracket D_{1} \rrbracket=\llbracket D_{2} \rrbracket \tag{7.13}
\end{equation*}
$$

In the following we will see that there are many cases where two visually different diagrams have the same matrix interpretation. The ZX-calculus provides a set of rules that allows to transform diagrams into semantically equivalent ones entirely in the graphical framework.

Another wire generator are the cup $\cup: 2 \rightarrow 1$, cap $\cap: 0 \rightarrow 2$ and SWAP $>: 2 \rightarrow 2$. These generators are interpreted as:

$$
\begin{align*}
& \llbracket \bigcup \rrbracket=|00\rangle\langle 00|+|10\rangle\langle 01|+|01\rangle\langle 10|+|11\rangle\langle 11|=\mathrm{SWAP}  \tag{7.14}\\
& \llbracket \cup \rrbracket=\langle 00|+\langle 11|  \tag{7.15}\\
& \llbracket \Omega \rrbracket=|00\rangle+|11\rangle \tag{7.16}
\end{align*}
$$

### 7.1.2 Compositions

Elementary generators are used to derive bigger diagrams using sequential and parallel composition. The parallel composition is usually referred as tensor product.

## Tensor product

We can obtain a ZX-diagram by putting two other diagrams aside each other. Formally, for



The term "tensor product" is employed as the interpretation of the new diagram is chosen to be exactly the Kronecker product of the matrices corresponding to $D_{1}$ and $D_{2}$ :

## Sequential composition

As we have observed on our first example (7.1), the output wires of one generator are usually connected to input wires of other generators. This procedure is called sequential composition.
 as
where the composition $A \circ B$ of two matrices $A_{2^{m} \times 2^{n}}$ and $B_{2^{l} \times 2^{m}}$ is matrix multiplication:

$$
\begin{equation*}
A \circ B=A B \tag{7.19}
\end{equation*}
$$

In the terminology of tensor networks each node of the diagram corresponds to a tensor and the wire connecting two nodes represents tensor contraction.

## Only topology matters

Using the sequential composition and the tensor product we can compute the interpretation of complex diagrams such as the one on the figure 7.1. However, our current definition strictly discriminates between input and output wires. Therefore, the nodes of the diagram have to come with a predefined order. It implies, in particular, that we still can't interpret a horizontal wire.

It turns out that in ZX-diagram the respective order of nodes actually doesn't matter. In other words, we can freely move nodes on the 2 D plane. As soon as the order of inputs and outputs is preserved, the interpretation doesn't change. This fantastic property follows from the set of equalities:

that are grouped under the paradigm "Only topology matters". We remark that the second set of formulas 7.21 is also true for the red spider.

These equalities can be verified by a careful computation of the corresponding matrices. Resulting matrix computations are remarkably tedious comparing to a simplicity of corresponding
transformations in ZX-calculus. In fact, in the graphical language we only have to remember that the wires can be bent at will.

By bending wires we can transform any diagram $D: n \rightarrow m$ with $n$ inputs and $m$ outputs to a state $D^{\prime}: 0 \rightarrow n+m$ :


For fixed dimensions $n$ and $m$ such mapping is bijective. The mapping is commonly referred as the Choi-Jamiolkiwski isomorphism Haapasalo, 2019.

### 7.2 Adjoint diagram

If we replace all phases $\alpha$ in a diagram $D$ by their negations the resulting diagrams $D^{\prime}$ will correspond to a matrix that is conjugate of $\llbracket D \rrbracket$. Indeed, for the complex phases $e^{-i \alpha}=\overline{e^{i \alpha}}$

Using cups and caps we can also find the transpose of a diagram:


By combining this two transformations we can obtain for each diagram its adjoint in a simple way. Moreover, as we can bend internal wires at will it actually is sufficient to simply "flip" the diagram upside down and negate the phases.

### 7.3 Universality

ZX-diagrams is a universal language for linear maps on dimensions that are powers of 2 Coecke and Duncan, 2011. The universality states that for each every linear map $M: \mathbb{C}^{2^{n}} \rightarrow \mathbb{C}^{2^{m}}$ there exists a diagram $D: n \rightarrow m$ such that the interpretation of $D$ is exactly equal to $M$.

The universality is relatively easy to prove. Firstly, we observe that every complex number $c \in \mathbb{C}$ can be represented as a scalar diagram $D: 0 \rightarrow 0$. Indeed, from the definition of generators we have:

## Lemma 7.3.1

$$
\begin{equation*}
\llbracket \bigcirc \rrbracket=2 \tag{7.24}
\end{equation*}
$$

$$
\llbracket \propto \rrbracket=1+e^{i \alpha}
$$

$$
\begin{align*}
& \| O=\sqrt{2}  \tag{7.25}\\
& \| O=\frac{1}{\sqrt{2}} \tag{7.26}
\end{align*}
$$

Any complex number can be decomposed as $Z=e^{i \beta} \cos (\alpha) 2^{n}$. The previous equation shows how this product can be obtained by composing scalar diagrams.

The universality of ZX-diagrams for unitary matrices straightforwardly follows from the universality of the gate set $\left\{\mathrm{H}, \mathrm{CNOT}, R_{X}(\theta), R_{Z}(\theta)\right\}$ for the pure qubit quantum computing. For H and $R_{Z}(\theta)$ the ZX-diagrams are already known. By explicitly computing the interpretation, we can show that:

$$
\begin{equation*}
\llbracket-\rrbracket=\sqrt{2} \mathrm{CNOT} \tag{7.27}
\end{equation*}
$$

A universal gate set allows to compute any complex vector of unit norm starting from the trivial state $|0 \ldots 0\rangle$. For other vectors we proceed by multiplication by the scalar factor equal to the norm. For maps with different number of input and outputs we use the Choi-Jamiolkiwski isomorphism (7.22).

An alternative proof of universality of ZX-diagrams was obtained using the normal forms Jeandel et al., 2019.

We remark that while it is quite easy to pass from a circuit to a diagram the backward translation is much more challenging Duncan et al., 2020. In fact, if a diagram doesn't correspond to a unitary matrix such translation is impossible without using ancilla and postselection. In the opposite case, a circuit can be efficiently extracted for a diagram that has a gflow Backens et al., 2021.

## Fragments

Up to know we considered the diagrams without any specific restriction on the phases. In other words, in general ZX-calculus the phases are allowed to take any value in $\mathbb{R} / 2 \pi \mathbb{Z}$. It turns out that in some settings it is meaningful to restrict the values of the angles to a specific sub-group $\mathcal{G}$ of $\mathbb{R} / 2 \pi \mathbb{Z}$. In such case we speak about the fragments of the language, conventionally denoted as $\mathrm{ZX}_{\mathcal{G}}$-calculus.

The interpretation of a ZX-diagram with general angles is a matrix with elements in $\mathbb{C}$. It was proven in Jeandel et al., 2019 that for a specific fragment defined by a group $\mathcal{G}$ the corresponding matrices have elements in the ring $\mathcal{R}_{\mathcal{G}}=\mathbb{Z}\left[\frac{1}{\sqrt{2}}, e^{i \mathcal{G}}\right]$. The notation $\mathbb{Z}\left[\frac{1}{\sqrt{2}}, e^{i \mathcal{G}}\right]$ stands for the smallest ring that contains $\mathbb{Z}, \frac{1}{\sqrt{2}}$ and $\left\{e^{i \alpha} \mid \alpha \in \mathcal{G}\right\}$. It was shown in Jeandel et al., 2019 that if the group $\mathcal{G}$ contains $\frac{\pi}{4}$ the $\mathbf{Z X}_{\mathcal{G}}$-diagrams are universal for the matrices over $\mathbb{Z}\left[\frac{1}{\sqrt{2}}, e^{i \mathcal{G}}\right]$.

A particular interest is attracted to finitely generated fragments. We use the notation $\mathcal{G}=$ $\left\langle x_{1}, \ldots, x_{n}\right\rangle$ to denote the smallest sub-group of $\mathbb{R} / 2 \pi \mathbb{Z}$ that contains all $x_{i}$ from the list. If the set of group generators contain only one element $\alpha$, we will refer to the corresponding fragment as $\alpha$-fragment

One of the most famous fragment is the Clifford fragment with $\mathcal{G}=\langle\pi\rangle$. Diagrams from the Clifford fragment allow to reason about the stabilizer quantum mechanics Backens, 2014].

The stabilizer quantum mechanics studies all transformations that can be obtained using measurements in the computational basis and unitary circuits made out of elements of the Clifford group. A peculiar property of such transformations is that they can be efficiently simulated on a classical computer Aaronson and Gottesman, 2004. A slightly more powerful real Clifford $\frac{\pi}{2}$-fragment was investigated in Duncan and Perdrix, 2014.

Gates from the Clifford group completed with so-called T-gate $T=\left(\begin{array}{cc}1 & 0 \\ 0 & e^{i \pi / 4}\end{array}\right)$ are approximately universal for pure qubit quantum computations. This set of gates can be represented by diagrams from the $\frac{\pi}{4}$-fragment Jeandel et al., 2018a, Ng and Wang, 2018.

Other fragments over rational angles were discussed in Jeandel et al., 2019.

### 7.4 Rewrite rules

As we have seen before, ZX-diagrams can represent any linear map of dimension $2^{m} \times 2^{n}$. However, a diagrammatic representation is far from being unique. We've already mentioned in $\quad 7.12$ ) an example of multiple diagrams for the trivial identity matrix.

A semantical equivalence of two diagrams can be established by explicitly computing corresponding matrices. Such approach is not very inspiring as the matrix representation usually has an exponential size for a given diagram. A slightly more efficient technique proceeds by the tensor network contraction. However, both approaches significantly reduce the interest of compact diagrammatic representation.

Alternatively, ZX-diagrams can be manipulated with semantic-preserving rewrite rules. A rule is an equality between diagrams $D_{1}: n \rightarrow m$ and $D_{2}: n \rightarrow m$ with the same amount of inputs and outputs. A rule (r) is called sound or, equivalently, semantic-preserving if:

$$
\begin{equation*}
D_{1} \underset{(r)}{=} \quad D_{2} \Longleftrightarrow \llbracket D_{1} \rrbracket=\llbracket D_{2} \rrbracket \tag{7.28}
\end{equation*}
$$

In principle, any sound equality $D_{1}=D_{2}$ can be states as a rule. For practical reasons, we are interested in keeping the set of rewrite rules as small as possible. If such set of rules is sufficient to establish any semantically valid equality it is called complete. In the following we present the most important rewrite rules together with the motivation behind them.

### 7.4.1 Meta-rule

According to the inductive definition, a composition of two ZX-diagrams is another diagram. When a sound rewrite rule is used to transform one of parts of the diagram the interpretation of the overall diagram doesn't change. This means that the local application of sound rewrite rules preserves the semantics. We call this fact a meta-rule:

In other words, if it is true that $\stackrel{|\cdots|}{|\cdots|} |$| $D_{0}$ |
| :---: |
| $\|\cdots\|$ |
| $D_{1}$ |
| $\|\cdots\|$ | for two diagrams $D_{1}, D_{2}: n \rightarrow m$ than for all

$D_{2}: k \rightarrow l$ and $D_{3}: m \rightarrow l$ we have:

### 7.4.2 Hadamard rules

The Hadamard gate has two crucial properties. The first property, called Euler decomposition, states that the Hadamard gate can be obtained with a diagram made out of green and red spiders:


In fact, there are multiple ways to represent a Hadamard with a fully spider diagram van de Wetering, 2020. We selected this particular one as it doesn't contain globals scalars. In this representation it is also obvious that the Hadamard gate is equal to its own transpose. Moreover, together with some rewrite rules that we will see later (EU) allows to diagrammatically prove that the Hadamard gate is equal to its own inverse, i.e.:

Lemma 7.4.1 [van de Wetering, 2020

$$
q=
$$

The second significant property of the Hadamard gate maps is that it maps the computational basis to the Hadamard basis:

$$
\begin{align*}
& \mathrm{H}|0\rangle=|+\rangle  \tag{7.32}\\
& \mathrm{H}|1\rangle=|-\rangle \tag{7.33}
\end{align*}
$$

This property translates to the following diagrammatic equality:


This equality is often referred as Hadamard conjugation. With the Hadamard conjugation rule we can flip colors of all red spiders and transform the diagram into a graph state form Duncan and Perdrix, 2009. Moreover, combined with the lemma 7.4.1 the Hadamard conjugation implies that any equality between two diagrams remains true with flipped colors of spiders.

It was shown in Duncan and Perdrix, 2009 that the rule (EU) can't be derived from a simpler set of rules. On the other hand, the second rule (H) is particularly important to manipulate nodes of different colors. For this reasons we select (EU) and (H) as basic rewrite rules.

### 7.4.3 Spider fusion

The special structure of green and red spider implies multiple useful diagrammatic equalities. One of them, the identity removal, was already presented in 7.12 . The identity removal states that we can remove an empty spider from a wire.

We have also seen that one can freely bend the legs of a spider (7.21). In fact, the symmetric structure of spiders leads to a much more powerful family of equalities called spider fusion:

where $" \ddots$. " means that there is at least one wire between spiders.
According to this rule we can fuse any two spiders of the same color that are connected by at least one wire. The phase of the resulting spider is equal to the sum of phases of individual spiders.

### 7.4.4 Bialgebra rules

In the previous section we have seen that two connected spiders of the same color can be fused. The specific interaction of two spiders of different colors follows from the fact that the computational basis $\{|0\rangle,|1\rangle\}$ and the Hadamard basis $\{|+\rangle,|-\rangle\}$ correspond to complementary observables $Z$ and $X$. In quantum mechanics complementarity means that we can't simultaneously measure the value of two observables in a quantum state Kiukas et al., 2019.

We remark that the interpretation of Z-spider implies that the spider


$$
\begin{equation*}
\llbracket \wp \rrbracket=|00\rangle\langle 0|+|11\rangle\langle 1|=\mathrm{COPY} \tag{7.36}
\end{equation*}
$$

copies the states from the computational basis, i.e. for a vector $|b\rangle, b \in\{0,1\}$ we have COPY $|b\rangle=|b b\rangle$. As the vector $|0\rangle$ is (up to a scalar) a one-legged red spider (7.8) we can write this equality as:


Using this rule together with (K), (S1) we can derive the following lemmas:

## Lemma 7.4.2



Lemma 7.4.3

where $k$ is a natural number.
 outputs their sum in $F_{2}:\left|b_{1} \oplus b_{2}\right\rangle$ :

$$
\begin{equation*}
\oiiint 』=|0\rangle\langle 00|+|0\rangle\langle 11|+|1\rangle\langle 01|+|1\rangle\langle 10|=\mathrm{XOR} \tag{7.40}
\end{equation*}
$$

An important rule about interaction of two spiders is a consequence of the relation of XOR and COPY operators:


Roughly speaking, this rule states that when we apply XOR to two entries and then copy the output we obtain the same result as when we start by copying each input and then XORing each pair separately.

In mathematical terms the rule (B2) states that the spider $\wp$ and form a bialgebra. A fundamental consequence of (B1) and (B2) (and fusion rules) is the Hopf rule:

## Lemma 7.4.4



The Hopf rule is precisely the expression of complementarity of $Z$ and $X$ in the language of ZX-diagrams.

Finally, by iteratively applying the bialgebra rule we can obtain an equation for arbitrary number of interacting spiders:

Lemma 7.4.5 Coecke and Kissinger, 2017, theorem 9.71


### 7.4.5 Completeness

Previously, we have introduced some important graphical rewrite rules that are sound:

$$
\begin{equation*}
D_{1} \underset{(r)}{\overline{=}} D_{2} \Longrightarrow \llbracket D_{1} \rrbracket=\llbracket D_{2} \rrbracket \tag{7.44}
\end{equation*}
$$

Using these rules we can perform computations diagrammatically, i.e. without passing by the matrix representation of the tensor network contraction. A set of rules $\mathcal{A}$ is called complete is it allows to demonstrate that any two diagrams with the same semantics are equivalent, i.e.

$$
\begin{equation*}
\llbracket D_{1} \rrbracket=\llbracket D_{2} \rrbracket \Longrightarrow \exists\left(r_{1}\right), \ldots,\left(r_{n}\right) \in \mathcal{A} \text { s.t. } \quad D_{1} \underset{\left(r_{1}\right)}{=} \cdots \underset{\left(r_{n}\right)}{=} D_{2} \tag{7.45}
\end{equation*}
$$

We use the term axioms to speak about the rules from a complete set $\mathcal{A}$.
For a long time the question about the existence of a compact set of complete rules was open. The first demonstration of completeness appeared for the Clifford fragment [Backens, 2014]. It was followed by a result for the real Clifford fragment Duncan and Perdrix, 2014.

First complete set of rules for the approximately universal Clifford +T fragment was presented in Jeandel et al., 2018a. It was followed by two completeness results for the general ZX-calculus Hadzihasanovic et al., 2018, Jeandel et al., 2018b.

On the figure 7.2 we show a compact set of axioms presented in Jeandel et al., 2019. This set of rules is complete for $\frac{\pi}{4}$-fragment. In Vilmart, 2019 it was shown that only one additional axioms suffices to make such set of rules complete for general ZX-diagrams. This axiom is:

where the angles $\alpha_{i=\{1,2,3\}}, \beta_{j=\{1,2,3\}}$ and $\gamma$ are related by a non-linear side condition. The additional axiom essentially states the equivalence of Euler decompositions of one-qubit unitary transformation.

Completeness is one of the major advantages of the ZX-calculus. It implies that we can prove any sound diagrammatic equality by a series of local graphical transformations from a compact set of axioms. We remark that in the set of rules on the Figure 7.2 we explicitly keep the scalar factors that are usually ignored in papers about ZX-calculus. The reason for it will become clear later when we will present the addition of diagrams. To easily manipulate scalars we will extensively use the lemma representing the equality $\frac{1}{\sqrt{2}} \times \sqrt{2}=1$ :

Lemma 7.4.6 (|Jeandel et al., 2019|)



Figure 7.2: Axioms for ZX. All rules remain true flipped upside down and with inverted colors. Families of equations are given using 'dots': ... means any number of wires, ${ }^{`} \cdot$ means at least one wire.

## Triangle

Initially, completeness proofs used the translation from ZX-calculus to another language called ZW-calculus for which a complete set of rules already existed. More modern results Vilmart, 2019, Jeandel et al., 2019 follow a different approach that pass by normal forms. More precisely, in such approach any diagram is transformed to a specific normal form. The reduction is carried out in a finite number of steps. The normal form is exclusively determined by the underlying matrix. Thus, two diagrams with the same semantics are reduced to the same normal form.

The construction for the normal form extensively uses a syntactic sugar called triangle orig-
inally introduced in Jeandel et al., 2018a:


The triangle corresponds to a non-unitary transformation:

$$
\llbracket \boldsymbol{\Perp} \rrbracket=|0\rangle\langle 0|+|1\rangle\langle 0|+|1\rangle\langle 1|=\left(\begin{array}{ll}
1 & 1  \tag{7.48}\\
0 & 1
\end{array}\right)
$$

that maps $|0\rangle$ to $\sqrt{2}|+\rangle$ and $|1\rangle$ to $|1\rangle$.
The work [Jeandel et al., 2019] proves multiple useful lemmas involving the triangle:

## Lemma 7.4.7



Lemma 7.4.8


Lemma 7.4.9


Lemma 7.4.10


Lemma 7.4.11


Lemma 7.4.12


Lemma 7.4.13


Lemma 7.4.14


Lemma 7.4.15


Lemma 7.4.16


Lemma 7.4.17


## Lemma 7.4.18



Lemma 7.4.19


Lemma 7.4.20

and


## Lemma 7.4.21



Lemma 7.4.22
O$=$

## Lemma 7.4.23

## Lemma 7.4.24





We also introduce some new lemmas that will be useful in the following chapters. The proofs serve as illustration of the graphical computation process.

Lemma 7.4.26



Lemma 7.4.27


Proof: [Proof (lemma 7.4.25]]


Proof: [Proof (lemma 7.4.27)] The left hand side is:


For the right hand side we get:


Proof: [Proof (lemma 7.4.26]]
To prove the lemma we adopt a slightly different technique that a usual sequence of rewrites. The technique consists in verification that the diagrams from the left hand side and the right hand side lead to the same result when applied on a set of basis states $\{|0\rangle,|1\rangle\}$ corresponding to diagrams \& and $\mathcal{O}_{8}^{\pi}$. As was pointed out in Coecke and Kissinger, 2017 for two matrices $M_{1}, M_{2}: \mathbb{C}^{2^{n}} \rightarrow \mathbb{C}^{2^{m}}$ and a basis $\left\{\left|b_{1}\right\rangle, \ldots,\left|b_{2^{n}}\right\rangle\right\}$ in $\mathbb{C}^{2^{n}}$ the equality on basis input is equivalent to the equality of matrices, i.e.

$$
\begin{equation*}
\forall i \in\left[1, \ldots, 2^{n}\right]: M_{1}\left|b_{i}\right\rangle=M_{2}\left|b_{i}\right\rangle \Longleftrightarrow M_{1} \equiv M_{2} \tag{7.52}
\end{equation*}
$$

 lish that $\left.\llbracket \begin{array}{|c|}\hline l h s \\ \frac{1}{1 . .1}\end{array} \rrbracket=\llbracket \begin{array}{c}1 \\ r h s \\ \ldots . .1 \\ \hline\end{array}\right]$. Finally, the lemma condition follows from the completeness of ZX-calculus.



### 7.5 Linear diagrams

In ZX-diagrams the angles inside spiders are real numbers or elements of some group $\mathcal{G}$. However, in many equations such as (S1), (K), (H), (SUP) and (C) we used symbols $\alpha$ and $\beta$ instead of numbers. The symbolic notation was extremely handy to define families of equalities, i.e.
equalities that hold for any real values of angles. By explicitly keeping symbols inside spiders we can draw parameterized diagrams. A parameterized diagram with $n$ different symbols can be interpreted as a function that associates to each element of the space $\mathbb{R}^{n}$ a well-defined ZXdiagram.

Between parametrized diagrams, we distinguish a family of linear diagrams denoted ZX $(\boldsymbol{\beta})$ :
[Linear diagrams Jeandel et al., 2018b] A ZX-diagram is linear in $\beta_{1}, \ldots, \beta_{k}$ with constants
 composition with $\alpha$ of the form $\sum_{i} n_{i} \beta_{i}+c$ with $n_{i} \in \mathbb{Z}$ and $c \in L$.

It was shown in Jeandel et al., 2018b that for $L=\left\{\frac{n \pi}{4}\right\}_{n \in \mathbb{Z}}$ the Clifford+T axiomatization (7.2) is complete for linear diagrams.

We say that a diagram $D$ is parametrized by $\beta_{1}, \ldots, \beta_{k}$ if its angles are some functions on $\beta_{1}, \ldots, \beta_{k}$. We denote such a diagram by $D\left(\beta_{1}, \ldots, \beta_{k}\right)$.

It is possible to evaluate a parametrized diagram $D(\boldsymbol{\beta}), \boldsymbol{\beta} \in \mathbb{R}^{k}$ in a point $\boldsymbol{\beta}^{0} \in \mathbb{R}^{k}$ by replacing every occurrence of $\beta_{i}$ with the respective value $\beta_{i}^{0}$. The result of evaluation is a diagram $D\left(\boldsymbol{\beta}_{0}\right)$ from $\mathrm{ZX}_{\mathbb{R}}$.

The family of linear diagrams may appear restricted compared to ZX-diagrams that allow angles from a more general class of functions. It is, however, sufficient for applications in variational quantum algorithms as they use circuits where parameters appear in a linear fashion Preskill, 2018.

## Chapter 8

## ZX-diagrams for variational quantum algorithms

Due to the relative simplicity of graphical manipulations comparing to matrix multiplications ZX-calculus is widely used to reason about applications from the quantum computing. Recently, ZX-calculus was applied to analyze parameterized quantum circuits Zhao and Gao, 2021, Stollenwerk and Hadfield, 2022, Fontana et al., 2020.

Parameterized quantum circuits (PQC) are essential building blocks of variational quantum algorithms (VQA). In the variational algorithms the parameters of the circuit are trained in a classical loop. The classical loop, in turn, may use the quantum computer to evaluate the loss function. The loss function is usually expressed as an expectation of some observable in a parameterized quantum state:

$$
\begin{equation*}
L(\boldsymbol{\theta})=\langle 0 \ldots 0| U^{\dagger}(\theta) O U(\theta)|0 \ldots 0\rangle \tag{8.1}
\end{equation*}
$$

### 8.1 The diagram for QAOA circuit

In this work we concentrate on the Quantum Approximate optimization Algorithm (QAOA). We recall that QAOA operates the circuit:

$$
\begin{equation*}
\left|\psi_{p}(\boldsymbol{\beta}, \gamma)\right\rangle=U_{M}\left(\beta_{m-1}\right) U_{P}\left(\gamma_{m-1}\right) \ldots U_{M}\left(\beta_{0}\right) U_{P}\left(\gamma_{0}\right)\left|\psi_{0}\right\rangle \tag{8.2}
\end{equation*}
$$

where $U_{M}(\beta)=e^{i \beta \sum_{i} X_{i}}$ and $U_{P}(\beta)=e^{i \gamma C}$ for a diagonal operator $C$. We will see in this section that the parameterized unitary 8.2 can be decomposed in terms of $R_{X}(\beta), R_{Z}(\gamma), \mathrm{H}$ and CNOT gates and easily translated to a ZX-diagram.

We recall that in previously used notations the time in a ZX-diagram flows vertically from top to the bottom. In this chapter in order to keep the parallel with the circuit notation we will adopt the horizontal time line. Therefore, the input of the diagram is on the left and the output on the right. As most of the diagrams in this section are scalars, the particular time direction actually doesn't matter at all.

## Phase gadgets

When QAOA is used to solve a combinatorial optimization problem the operator $C$ encodes the objective function. For constraint satisfaction optimization problems the objective function
$f:\{0,1\}^{n} \rightarrow \mathbb{R}$ is usually given as:

$$
\begin{equation*}
f(\boldsymbol{x})=\sum_{S \subset[1, \ldots, n]} c_{S} \bigoplus_{i \in S} x_{i} \tag{8.3}
\end{equation*}
$$

For example, for the MaxCut problem we have:

$$
\begin{equation*}
f(\boldsymbol{x})=\sum_{(u, v) \in E} e_{u, v} x_{u} \oplus x_{v} \tag{8.4}
\end{equation*}
$$

A unitary $U_{\mathcal{X}}(\gamma)=e^{i \gamma x_{i_{1}} \oplus \cdots \oplus x_{i_{l}}}$ acting on qubits from the set $\mathcal{X}$ corresponds (up to a scalar $\lambda$ ) to a phase gadget. Phase gadgets are ZX-diagrams of the form:

up to a scalar factor $\lambda$.
The work Cowtan et al., 2020 is dedicated to the properties of phase gadgets and their generalization called phase polynomials. We provide a small example showing how obtain a phase polynomial diagram representing $e^{i \gamma C}$ for an Ising Hamiltonian $C$ :

$$
\begin{equation*}
C=\sum_{i=1}^{n} h_{i} Z_{i}+\sum_{1 \leq i<j \leq n} h_{i, j} Z_{i} Z_{j} \tag{8.6}
\end{equation*}
$$

A ZX-diagram for the term $e^{i \gamma C}$ can be composed out of diagrams for individual and pair terms $e^{i \gamma h_{i} Z_{i}}$ and $e^{i \gamma h_{i, j} Z_{i} Z_{j}}$. For the one-qubit term $e^{i \gamma h_{i} Z_{i}}$ the diagram is simply-2 $2 h_{i} \gamma-$ up to the phase scalar $e^{-i \gamma h_{i}}$.

For the pair interaction we know that

$$
\begin{equation*}
e^{i \gamma Z_{i} Z_{j}}=e^{i \gamma} \operatorname{CNOT} R_{Z}(-2 \gamma) \text { CNOT } \tag{8.7}
\end{equation*}
$$

Up to a scalar that doesn't depend on $\gamma$, the diagrams for the number $e^{i \gamma}$ and for the gate CNOT are $\underset{\gamma}{\pi}$ and
 respectively. Therefore, we can obtain a diagram for $e^{i \gamma Z_{i} Z_{j}}$ using the composition:


We introduce the shortcuts $\gamma_{i}=2 \gamma h_{i}$ and $\gamma_{i, j}=2 \gamma h_{i, j}$. Using phase gadgets we can easily draw a ZX-diagram for the parameterized $e^{i \gamma C}$ :


In this diagram we dropped the global complex phase as it is not meaningful in our context.
The uniform superposition $\frac{1}{\sqrt{2^{n}}} \sum_{x \in\{0,1\}^{n}}|x\rangle=|+\rangle \otimes \cdots \otimes|+\rangle$ that is used as the initial state can be represented by the diagram $n\left\{\begin{array}{l}\vdots \\ \vdots\end{array}\right.$. The diagram for the mixing operator $U_{M}(\beta)=\prod_{1 \leq i \leq n} e^{i \beta X_{i}}$ is also straightforward:


By combining the building blocks we obtain a diagram for the one-depth QAOA state $|\psi(\beta, \gamma)\rangle=e^{i \beta \sum_{i} X_{i}} e^{i \gamma C}|+\cdots+\rangle:$


### 8.2 Loss function

The classical loop of QAOA optimizes the loss function. The loss function is given by the expectation of the operator $C$ in the quantum state:

$$
\begin{equation*}
\langle\psi(\beta, \gamma)| C|\psi(\beta, \gamma)\rangle \tag{8.11}
\end{equation*}
$$

For the Ising Hamiltonian the energy expectation can be computed as the weighted sum of individual contributions $\left\langle Z_{v}\right\rangle_{\beta, \gamma}$ and $\left\langle Z_{u} Z_{v}\right\rangle_{\beta, \gamma}$ :

$$
\begin{equation*}
\langle C\rangle_{\beta, \gamma}=\sum_{i=1}^{n} h_{i}\left\langle Z_{i}\right\rangle_{\beta, \gamma}+\sum_{1 \leq i<j \leq n} h_{i, j}\left\langle Z_{i} Z_{j}\right\rangle_{\beta, \gamma} \tag{8.12}
\end{equation*}
$$

In the next chapter we will see how to represent such sum as a ZX-diagram. For now, we will consider only the individual terms $\left\langle Z_{v}\right\rangle_{\beta, \gamma}$ and $\left\langle Z_{u} Z_{v}\right\rangle_{\beta, \gamma}$. In particular, we will show how to derive the analytical expression for these values with the means of ZX-calculus.

We consider QAOA states of depth $p=1$. Originally, an analytical expression for the mean values in QAOA states was presented in Wang et al., 2017. The formula was derived for instances of unweighted MaxCut problem. In Dalyac et al., 2021 we extended this result to the weighted MaxCut (see the formula 3.37).

In this section we will demonstrate how to use ZX-calculus to derive an analytical formula for general Ising models. This result was derived as a part of the thesis project and is presented in the current manuscript for the first time. We remark that a similar objective was pursued in [Stollenwerk and Hadfield, 2022] where ZX-diagrams are used to re-derive the formula for unweighted MaxCut. Contrary to Stollenwerk and Hadfield, 2022, we don't proceed by an explicit extension of the traditional framework. Instead, we locally use a decomposition of the identity matrix.

We will also extensively use the following lemmas that relate to the supplementarity of quantum states Perdrix and Wang, 2016:

Lemma 8.2.1


## Lemma 8.2.2



Lemma 8.2.3


Lemma 8.2.4


## Lemma 8.2.5



## Lemma 8.2.6



Proof: 8.2.2



The equality 8.2 .3 simply follows from $-\pi-\pi-\pi \overline{\overline{\mathrm{S} 1}}$
Proof: 8.2.4


Proof: 8.2.5


Proof: 8.2.6


We use the ZX rules to make as much simplifications as possible in the graphical framework. At the end we aim to get a formula that corresponds to scalar diagrams representing $\left\langle Z_{i}\right\rangle_{\beta, \gamma}$ and $\left\langle Z_{i} Z_{j}\right\rangle_{\beta, \gamma}$. To obtain the final analytic expression we will use the following result:


In order to have diagrams of reasonable size, we introduce a new syntactic sugar. For each finite set of diagrams $D_{1}, \ldots, D_{s}$ with the same number of input and output wires we denote:


We remark that this notation is by no mean necessary for our proofs. It just make the computations easier to follow. In Kissinger et al., 2014 and Carette and Lemonnier, 2022 a similar notation was introduced under the name !-box for general string diagrams and ZH diagrams respectively. We will also sometimes use red rectangle to emphasize the part of diagram that is transformed.

### 8.2.1 The mean value of $\left\langle Z_{i}\right\rangle$

The diagram for the operator $Z$ is simply $-\pi-$.
We can evaluate the expectation $\langle\psi(\beta, \gamma)| Z|\psi(\beta, \gamma)\rangle$ by explicitly composing diagrams for the operator $Z$, the state $|\psi(\beta, \gamma)\rangle$ and its adjoint. We recall that the diagram for the adjoint $\langle\psi(\beta, \gamma)|$ is precisely the diagram for $|\psi(-\beta,-\gamma)\rangle$ flipped from left to right.

The diagram for $\langle\psi(\beta, \gamma)| Z|\psi(\beta, \gamma)\rangle$ is:


We remark that all phase gadgets connecting qubits $i$ and $j$ that are not acted by the Pauli operator cancel while leaving a constant factor:


Such simplifications lead to the diagram:


We use (K) to move the $\pi$-spider to the left and (S1) to fuse and unfuse spiders of the same
colors. We obtain the diagram:

where $d$ is the degree of the node 1 . We assume that the connectivity graph of the Hamiltonian $C$ is full. A missing edge is simply an edge of zero weight. Therefore, in our settings $d=n-1$.

We can apply the generalized bialgebra rule (7.4.5 to spiders. After the rule application, the diagram becomes:


Then we apply the lemma 8.2.3 to obtain:


Finally, the lemmas 7.4 .2 and 8.2.1 lead us to a diagram that is a simple multiplication of
known cosines and sinus scalars 8.18):


Using the result 8.18, we can directly compute the scalar value of the final diagram. This value is:

$$
\begin{equation*}
\langle\psi(\beta, \gamma)| Z|\psi(\beta, \gamma)\rangle \underset{\square}{\overline{[8.18}} \quad \sin \beta \sin \gamma_{1} \prod_{i=2}^{n} \cos \gamma_{1, i} \tag{8.28}
\end{equation*}
$$

### 8.2.2 The mean value of $\left\langle Z_{i} Z_{j}\right\rangle$

In the previous section we managed to compute $\langle\psi(\beta, \gamma)| Z_{i}|\psi(\beta, \gamma)\rangle$ using only graphical rewritings for internal transformations. For the two-qubit operator $Z_{i} Z_{j}$ the situation is more complicated. Firstly, we start with the diagram:


By performing spider fusion and commuting the $\pi$-spider with (K) rule we get a slightly easier representation.


At this step it is unclear how to proceed further. Indeed, linear rules don't seem to simplify much the diagram. Moreover, we recall that the diagram equals to a complex formula for the mean value in QAOA state (3.37). The formula contains sums and can't be decomposed in a product of cosines and sinuses with angles that are linear on parameters $\beta$ and $\gamma$. Therefore, the diagram 8.30 is probably the simplest one that express the desired mean value.

At this step the work Stollenwerk and Hadfield, 2022 suggests to explicitly decompose Pauli rotations. The decomposition is integrated in the language with a new sum-gadget that represents a linear combination of diagrams:

where $a$ and $b$ are the coefficients of the decomposition. The diagrams $: A A_{B}$ and $: B$ have the same number of input and output wires. The new gadget comes with product pull rule and diagram pull rule. The product pull rule allows to compose two subsequent sum-gadgets. The diagram pull rule essentially embodies the distributivity of addition with respect to tensor product and matrix multiplication.

In Stollenwerk and Hadfield, 2022 the sum-gadget is used to commute $Z$-spiders and $X$ spiders in non-trivial cases when angles are not 0 or $\pi$. For the mean values appearing in QAOA the work extensively uses the decomposition of the spiders corresponding to Pauli rotations:


We remark that explicit linear combinations of diagrams, alternatively called bags of diagrams were already used in Toumi et al., 2021 in the context of differentiation. In fact, until the recent paper Yeung et al., and our contribution JJeandel et al., 2022] it was the only way to deal with sums. In the next chapter we will present how to add diagrams entirely diagrammatically.

To derive the analytical formula, we suggest to use a more restrictive set of tools than the one presented in Stollenwerk and Hadfield, 2022. Actually, the only thing we need is the trivial decomposition of the identity matrix:

$$
\begin{equation*}
\llbracket-\rrbracket=\frac{1}{2}(\llbracket-\bigcirc \bigcirc-\rrbracket+\llbracket-\pi \circlearrowleft-\rrbracket) \tag{8.33}
\end{equation*}
$$

For linear maps of same dimensions the addition can be naturally defined. Trivially, a sum of two matrices can be obtained with entry-wise addition of their elements. A formal definition for sum operator is given in Coecke and Kissinger, 2017 (chapter 5). According to this definition, the sum operator is distributive with respect to tensor product and composition. The distributive property implies that we can apply the decomposition (8.33) locally for an internal wire of the diagram:

Using this decomposition, we can prove the following lemma:

## Lemma 8.2.7



## Proof:



We apply the decomposition to the wire surrounded by red rectangle. The first term becomes:


The second term may be computed similarly:





$$
\begin{equation*}
\stackrel{\boxed{8.18}}{\longrightarrow} 2 \cos \left(\alpha_{0}-\beta_{0}\right) \prod_{i=1}^{s} \cos \left(\alpha_{i}-\beta_{i}\right) \tag{8.39}
\end{equation*}
$$

We will use the identity decomposition to compute the value of the diagram 8.30). In a manner, our graphical computations will be assisted with the semantics equalities.

Firstly, we transform the part of the diagram 8.30 in the red frame.

We now reduce overall diagram for each part of the decomposition 8.40). For the moment, we will ignore the common scalar factor

The scalar factor will be added in the final steps of the computation.
For the first part corresponding to $A$ in (8.40) after the application of (S2) and (S1) we obtain:


With the blue frame the diagram can be written in a more compact way.


We apply the lemma 8.2 .4 to the part in red frame.


With the lemma 7.4 .8 that allows to push red $\pi$-spider through the greed spider and the rule (S1) we get:

(2 ${ }^{(2)}$


$$
i=[3, \ldots, n]
$$

Multiplying this diagram by the scalar factor we obtain with the aid of lemmas 7.4 .6 and 8.18 the formula:

$$
\begin{equation*}
\frac{1}{2} \llbracket A \rrbracket=\sin \beta \cos \beta \sin \gamma_{1,2} \cos \gamma_{2} \prod_{i=3}^{n} \cos \gamma_{2, i} \tag{8.47}
\end{equation*}
$$

For the part that correspond to $B$ in the decomposition we have:


We apply one more time the decomposition 8.40 for the bottom spider. We obtain two diagrams $A_{1}$ and $B_{1}$ :

$$
\begin{equation*}
\llbracket B \rrbracket=\frac{1}{2}\left(\llbracket A_{1} \rrbracket+\llbracket B_{1} \rrbracket\right) \tag{8.49}
\end{equation*}
$$

The first diagram $A_{1}$ is


We simplify $A_{1}$ firstly with the lemmas 8.2 .5 and 8.2 .4 .


Then the application of (S2) and (B1) leads to:




$$
i=[3, \ldots, n]
$$

We use the lemma 8.2.3 and the rules (S1), (S2) to find the result:


We remark this diagram is almost the same as 8.46 with interchanged indexes 1 and 2 . Therefore, by multiplying this diagram by the factor with $\frac{1}{2}$ resulting from the 8.49 we obtain an expression similar to 8.47):

$$
\begin{equation*}
\frac{1}{2} \llbracket s A_{1} \rrbracket=\sin \beta \cos \beta \sin \gamma_{1,2} \cos \gamma_{1} \prod_{i=3}^{n} \cos \gamma_{1, i} \tag{8.54}
\end{equation*}
$$

The second term corresponding to $B_{1}$ is reduced as follows:


Firstly we apply the lemma 8.2.5.


Then we simply apply the (S1) to reproduce the condition of the lemma 8.2.6.


Applying 8.2.6 we obtain:


We combine scalar phases with 7.4.7.


At the last step we use the generalized bialgebra rule 7.4 .5 ,


We've already seen a similar diagram in the lemma 8.2.7. After the multiplication by the scalar $s^{\prime}=s / 2$ we get the last component of the mean value formula:

$$
\begin{align*}
\frac{1}{2} \llbracket s B_{1} \rrbracket & =-\sin ^{2} \beta(-1)\left(\cos \left(\gamma_{1}+\gamma_{2}+\pi\right) \prod_{i=3}^{n} \cos \left(\gamma_{1, i}+\gamma_{2, i}\right)+\cos \left(\gamma_{1}-\gamma_{2}\right) \prod_{i=3}^{n} \cos \left(\gamma_{1, i}-\gamma_{2, i}\right)\right) \\
& =\sin ^{2} \beta\left(\cos \left(\gamma_{1}-\gamma_{2}\right) \prod_{i=3}^{n} \cos \left(\gamma_{1, i}-\gamma_{2, i}\right)-\cos \left(\gamma_{1}+\gamma_{2}\right) \prod_{i=3}^{n} \cos \left(\gamma_{1, i}+\gamma_{2, i}\right)\right) \tag{8.61}
\end{align*}
$$

Combining the terms (8.47), (8.54), (8.61) we obtain:

$$
\begin{align*}
\langle\psi(\beta, \gamma)| Z_{1} Z_{2}|\psi(\beta, \gamma)\rangle= & \frac{\sin 2 \beta}{2} \sin \gamma_{1,2}\left(\cos \gamma_{1} \prod_{i=3}^{n} \cos \gamma_{1, i}+\cos \gamma_{2} \prod_{i=3}^{n} \cos \gamma_{2, i}\right)+ \\
& +\sin ^{2} \beta\left(\cos \left(\gamma_{1}-\gamma_{2}\right) \prod_{i=3}^{n} \cos \left(\gamma_{1, i}-\gamma_{2, i}\right)-\cos \left(\gamma_{1}+\gamma_{2}\right) \prod_{i=3}^{n} \cos \left(\gamma_{1, i}+\gamma_{2, i}\right)\right) \tag{8.62}
\end{align*}
$$

### 8.3 Discussion

In the previous section we used the rules of ZX-calculus to derive formulas (8.28) and (8.62) for the loss function in QAOA of depth $p=1$. Our result was derived for a general Ising model. Our analytical formula are extremely helpful for numerical simulations. For instance, the efficient classical simulation of RQAOA with depth $p=1$ used to obtain the results presented in sections 5.3 .3 and 6.4 .3 is owed to this formula. In addition, our original demonstration performed in the framework of ZX-calculus allows to connect two important domains of the quantum computing: ZX-calculus and variational algorithms.

In this particular example we can hardly claim that the diagrammatic representation made the computation easier. Moreover, sometimes we had to use semantic equalities such that 8.33) to be able to continue the simplifications. Nevertheless, our computation pointed out several interesting results.

Firstly, we demonstrated how to express the routines of QAOA in the framework of ZXcalculus. Such diagrammatic representation can have its own merits. Possible application is the study of the parameter landscape. For instance, in the context of quantum machine learning ZX-diagrams were used to reason about the barren plateau phenomena Zhao and Gao, 2021.

Secondly, we have shown how to perform the computation for a difficult quantum algorithm by manipulating the diagrams with the rewrite rules. We also suggested a combined approach that outsources only a part of computation to the graphical rewriting. This combination passed by the semantical decomposition of identity on a sum of two projectors 8.33. In Stollenwerk and Hadfield, 2022 a slightly different technique was used. In this alternative technique the spiders corresponding to Pauli gates were decomposed in a linear combination.

Finally, we have observed that despite multiple advantages of the graphical representation, the framework was still lacking tools that are important to reason about variational algorithms. One such desirable tool is addition. Indeed, in the traditional ZX-calculus there is no an efficient way to extract a diagram from the expression corresponding the Ising Hamiltonian 8.6. Obviously, as ZX-diagrams are universal such diagram exists. Moreover, it can be constructed from the explicit matrix representation while the complexity of such procedure is exponential.

The other missing tool is differentiation. It is particularly important to reason about the training of the variational circuit. In the next chapter we will present our original technique for the graphical addition and differentiation of ZX-diagrams.

## Chapter 9

## Addition and Differentiation of ZX-diagrams

This chapter is based on our results presented in Jeandel et al., 2022.

### 9.1 Motivation

Due to its flexibility, ZX-calculus is widely used to address different problems of quantum computing. However, its application to the rapidly growing field of variational algorithms Cerezo et al., 2021a such as QAOA Farhi et al., 2014a, VQE Peruzzo et al., 2014 and variational quantum machine learning are so far limited. Nevertheless as variational algorithms do not require heavy resource for error-correction, the incoming emergence of $N I S Q$ devices makes from them an object of particular attention Preskill, 2018.

In the previous chapter we have shown how to derive the analytical expression for the loss function in a particular case of QAOA with depth $p=1$. We remark that at some point we had to decompose the diagram on a sum of two terms. The resulting computations were quite tedious. The advantage of diagrammatic representation was insignificant.

However, we believe that the study of variational algorithms with the means of ZX-calculus is a promising research direction. For instance, the diagrammatic representation is beneficial in the analysis of the barren plateau phenomena Zhao and Gao, 2021. Yet the work Zhao and Gao, 2021 pointed out a crucial obstacle limiting the application of ZX-calculus to variational algorithms. This obstacle turns out to be the absence of convenient tools for addition and differentiation of parametrized diagrams.

Indeed, basic building blocks of variational algorithms are parametrized circuits and the search of optimal parameter values is a crucial part of these algorithms. The search is usually done by classical numerical optimization methods and most of them use derivatives Guerreschi and Smelyanskiy, 2

Addition is closely related to differentiation via product rules. Product rules are crucial to evaluate the derivative of the composition and the tensor product.

Furthermore, in variational algorithms the loss function is typically given as the expectation of a local Hamiltonian in a parameterized quantum state. Addition of diagrams may be used to obtain graphical representations of such Hamiltonians. We recall that Hamiltonians are nonunitary Hermitian operators corresponding to the energy of the system. In practical applications

Hamiltonians are usually given as a sum of local terms:

$$
\begin{equation*}
H=\sum_{i=1}^{k} H_{i} \tag{9.1}
\end{equation*}
$$

where each $H_{i}$ acts only on a limited number of qubits. Usually, individual terms $H_{i}$ can be relatively easily represented as diagrams. For instance, for the Ising Hamiltonian $H_{i}=Z_{u} Z_{v}$ and the diagram for individual terms is trivial: $-\pi-\pi$. However, in the traditional framework there is no graphical procedure that constructs the overall sum out of such elementary blocks.

### 9.1.1 Addition in quantum computing

In pure qubit quantum mechanics computations are performed by unitary transformations. Unitary transformations form a group with respect to multiplication. On the other side, a sum of two unitary transformations $U_{a}$ and $U_{b}$ is not unitary. However, in some context it is meaningful to consider transformations of the form $U_{a}+U_{b}$.

For instance, the simulation of an evolution under some Hamiltonian may be done more efficiently using linear combination of unitary operators Childs and Wiebe, 2012. The Hamiltonian simulation problem that is central problem for quantum chemistry. The traditional solution for the problem uses Lie-Trotter-Suzuki product formulas. As was shown in Childs and Wiebe, 2012 this standard approach may outperformed by using linear combination of unitaries.

The work Childs and Wiebe, 2012 suggests the following protocol for the addition of unitaries:

where $q_{a}$ is an ancilla qubit initialized in a zero-state, and unitary gates $U_{a}$ and $U_{b}$ are conditioned on an ancilla.

The protocol $\sqrt{9.2}$ can be easily verified:

$$
\begin{align*}
|0\rangle|\psi\rangle \rightarrow \frac{|0\rangle+|1\rangle}{\sqrt{2}}|\psi\rangle & \rightarrow \frac{1}{\sqrt{2}}\left(|0\rangle \otimes U_{b}|\psi\rangle+|1\rangle \otimes U_{a}|\psi\rangle\right)  \tag{9.3}\\
& \rightarrow \frac{1}{2}\left(|0\rangle \otimes\left(U_{b}+U_{a}\right)|\psi\rangle+|1\rangle \otimes\left(U_{b}-U_{a}\right)|\psi\rangle\right) \tag{9.4}
\end{align*}
$$

If the ancilla qubit is measured in $|0\rangle$ state, the input state $|\psi\rangle$ was successfully transformed to $\frac{1}{c}\left(U_{a}+U_{b}\right)|\psi\rangle$ where $c$ is the normalization constant.

We remark that the protocol 9.2 use an ancilla qubit and post-selection on the measurement output. Therefore, it exceeds the framework of pure qubit computations. We also emphasize that the circuit involves controlled versions of unitary transformations. In order to be executed on a quantum computer, the controlled unitaries have to be compiled to a sequence of elementary gates. An efficient compilation is non-trivial even for a basic controlled-CNOT gate usually called Toffoli gate. Interestingly, a graphical representation of controlled gates is particularly natural in ZH-calculus Backens and Kissinger, 2019. The graphical representation may enhance the compilation of controlled gates Kuijpers et al., 2019.

Traditional ZX-calculus is not equipped with a protocol allowing the linear combinations of arbitrary ZX-diagrams. Fundamentally, this deficiency is due to the absence of a proper physical interpretation for a sum of diagrams. Indeed, ZX-calculus is a language for process theory representation of quantum computing Coecke and Kissinger, 2017. In other words, each diagram is a process. The composition of processes can be interpreted as their sequential application. The tensor product corresponds to the parallel application. In the context of process theory, the addition doesn't have a meaningful physical interpretation.

Several attempts were made to address this issue. The first option selected in Toumi et al., 2021, Stollenwerk and Hadfield, 2022 consists in an extension of the formalism from diagrams to bags of diagrams. More precisely, this approach suggest to consider formal sums of diagrams as an object of study. Unfortunately, there is no rules to manipulate sums of ZX-diagrams. Therefore, the interest of graphical computations is significantly reduced. In the previous chapter we have seen an example application of this approach confirming our doubts.

In Jeandel et al., 2022 we have introduced an original technique that allows to add diagrams entirely graphically. In our approach a sum of diagrams is another diagram. Interestingly, we derived the rule for diagrammatic addition using controlled versions of diagrams. We suggest an inductive procedure that allows to construct a controlled version for arbitrary diagram. Our result with be presented in the section 9.2 .

An independent work Yeung et al., suggests an alternative protocol for diagrammatic addition for slightly more expressive algebraic ZXW-calculus. This result is also based on a translation to a controlled form. However, their definition of a controlled form is not equivalent to ours. In addition, instead of inductive procedure, they rely on a decomposition of a diagram on elementary matrices that have an exponential complexity in the worst-case. Later in this chapter we will provide a more detailed comparison of two approaches.

## Representing Hamiltonians as ZX-diagram

The addition of diagrams naturally appears in the representation of Hamiltonians. Usually, Hamiltonians are given as a weighted sum of Pauli tensors. Pauli tensors have a particularly simple decomposition on elementary matrices. Therefore, the procedure described in Yeung et al., is well-suited for the construction of diagrammatic representation for such Hamiltonians.

In addition, the work Yeung et al., also suggests a way to exponentiate diagrams representing Hamiltonians. The exponentiate of a Hamiltonian $H$ is a unitary $e^{i \theta H}$. The exponentiation is done with the aid of the Cayley-Hamilton theorem. The Cayley-Hamilton theorem allows to decompose the exponent $e^{i \theta H}$ on a linear combination of $H^{k}=\underbrace{H \times \cdots \times H}_{k}$ for $k \in\left[1, \ldots, 2^{n}\right]$.

We remark that in a particular case when $H=\sum_{i} H_{i}$ is a sum of commuting terms it is usually straightforward to obtain $e^{i \theta H}$ as a product of individual $e^{i \theta H_{i}}$. However, if some local terms in $H$ don't commute the product decomposition doesn't apply. For instance, this is the case for modified mixer Hamiltonians used in the extension of QAOA to constrained optimization problems Hadfield et al., 2019.

In our work Jeandel et al., 2022 we considered an inverse problem. Using the Stone's theorem, we have demonstrated how to obtain a diagram for $H$ given the derivative of the diagram for $e^{i \theta H}$. We illustrate this procedure on an example of Ising Hamiltonian. Our result in presented the Section 9.4 .

The Stone's theorem provides a link between differentiation and addition. With this theorem some specific kind of sums may be straightforwardly expressed as diagrams if we can compute derivatives. On the other directions, the ability to perform addition directly leads to a procedure
for differentiation. Indeed, the major obstruction for diagrammatic differentiation comes from the product rules that involve sums. In the Section 9.3.1 we will show how an inductive definition of the derivative is obtained by explicit diagrammatic representation of the product rules.

### 9.1.2 Differentiation in variational algorithms

In the previous chapter we have seen that the loss function in variational algorithms can be represented as a parameterized scalar ZX-diagram. Many state-of-the art methods for the optimization of the loss function such as Quantum Analytic Descent Koczor and Benjamin, 2022 and meta-learning oprimizers Wilson et al., 2021 use derivatives.

For an efficient training it is important to detect such deficiencies in the parameter landscape as barren plateau. In the context of variational algorithms, it was demonstrated that a specific structure of the cost function may lead to such phenomena Cerezo et al., 2021b. Alternatively, the structure of the ansatz itself may lead to the exponentially vanishing gradient McClean et al., 2018.

For the first time a graphical study of barren plateau was presented in Zhao and Gao, 2021. This work analyzes the diagrams with only two occurrences of the parameter. In Toumi et al., 2021 an arbitrary number of parameter occurrences is tackled with explicitly written product rules leading to bags of diagrams.

Alternatively, diagrammatic representation can be used to design parameter shift rules.

## Parameter shift rules

Parameter shift rules appeared as a solution to the problem of gradient evaluation. Indeed, as we have pointed out in the Section 3.5 .7 the gradient of the loss function $\frac{\partial\left\langle\psi_{0}\right| U^{\dagger}(\boldsymbol{\alpha}) O U(\boldsymbol{\alpha})\left|\psi_{0}\right\rangle}{\partial \alpha_{i}}$ in general can't be directly evaluated on a quantum hardware. However sometimes we can express the gradient as a linear combination of the loss function values in different points:

$$
\begin{equation*}
\frac{\partial L}{\partial \alpha_{i^{*}}}(\hat{\boldsymbol{\alpha}})=\sum_{k=1}^{m} \epsilon_{k} L\left(\boldsymbol{\alpha}^{\boldsymbol{k}}\right) \tag{9.5}
\end{equation*}
$$

where $\alpha_{i}^{k}=\left\{\begin{array}{ll}\hat{\alpha}_{i}, & \forall i \neq i^{*} \\ \hat{\alpha}_{i}+\phi_{i}, & i=i^{*}\end{array}\right.$ The values $L\left(\boldsymbol{\alpha}^{\boldsymbol{k}}\right)$ can be computed with stochastic approximation.

Decompositions of the form (9.5) are called parameter shift rules.
In most cases such rules consider the parameterized operators of the form $U(\alpha)=e^{i \alpha H}$ for a Hermitian operator $H$ called generator. For instance, the first shift rule was derived for the generators $H$ with two eigenvalues +1 and -1 Mitarai et al., 2018. It was further extended to generators with arbitrary symmetric eigenvalues $+r$ and $-r$ Schuld et al., 2019:

$$
\begin{equation*}
\frac{\partial L}{\partial \alpha}(\alpha)=\frac{r}{4}\left(L\left(\alpha+\frac{\pi}{4 r}\right)-L\left(\alpha-\frac{\pi}{4 r}\right)\right) \tag{9.6}
\end{equation*}
$$

An arbitrary parameterized transformation can be decomposed on a sequence of elementary gates $U(\alpha)=E_{1}(\alpha) \times \cdots \times E_{M}(\alpha)$ where some of the gates depend on the parameter $\alpha$. Typically, parameterized elementary gates are one-qubit Pauli rotations with two eigenvalues. Using the decomposition on elementary gates we can express gradients for small gates $E_{i}(\alpha)$ with the
parameter shift rule. The gradient of the overall unitary $U(\alpha)$ is computed with the product rule Crooks, 2019:

$$
\begin{equation*}
\frac{\partial U(\alpha)}{\partial \alpha}=\frac{\partial E_{1}(\alpha)}{\partial \alpha} \times \cdots \times E_{M}(\alpha)+\cdots+E_{1}(\alpha) \times \cdots \times \frac{\partial E_{M}(\alpha)}{\partial \alpha} \tag{9.7}
\end{equation*}
$$

In the decomposition approach the number of shifts that require evaluation scales polynomially with the number of parameter occurrences. Alternatively, a spectral decomposition of $H$ can be used to design a rule with reduced number of shifts. This approach was developed in Wierichs et al., 2022, Kyriienko and Elfving, 2021. Two different algebraic approaches were presented in Izmaylov et al., 2021.

Such generalized rules are a priori more economical than the ones resulting from the decomposition on elementary gates. More precisely, the generalized rules lead to a smaller number of elements in the linear combination (9.5). As each value $L\left(\boldsymbol{\alpha}^{\boldsymbol{k}}\right)$ is typically evaluated on a quantum computer the number of terms in the linear combination is an important cost criteria. Remarkably, for the second-order derivatives the potential savings of generalized rules compared to decomposition approach are of order $O(n)$ Wierichs et al., 2022.

We remark that the parameter shift rules lead to an unbiased estimator for the gradient compared to the finite difference approximation $\frac{\partial L}{\partial \alpha}\left(\alpha_{0}\right)=\frac{L\left(\alpha_{0}+\delta\right)-L\left(\alpha_{0}\right)}{\delta}$. Both approaches were compared analytically and numerically in Mari et al., 2021. It was observed that parameter shift rules lead to more correct gradient estimations which, in turn, enhance the parameter optimization process.

We remark that even if we have a diagram for the derivative, we need to further extend the toolbox of ZX-calculus to suggest new parameter-shift rules. For instance, we may need a procedure that compiles a diagram into a linear combination of other diagrams corresponding to valid quantum expectations. In the diagrammatic analysis we don't have to systematically recover the structure of the initial circuit. As a consequence, we may end up with rules in a more general form:

$$
\begin{equation*}
\frac{\partial L}{\partial \alpha}(\alpha)=\left\langle\psi_{1}(\alpha)\right| O_{1}\left|\psi_{1}(\alpha)\right\rangle+\cdots+\left\langle\psi_{M}(\alpha)\right| O_{1}\left|\psi_{M}(\alpha)\right\rangle \tag{9.8}
\end{equation*}
$$

## Diagrammatic differentiation

The works Zhao and Gao, 2021, Toumi et al., 2021 use explicit sums of diagrams in order to represent the derivative of diagrams with multiple occurrences of the parameter. In contrast, in our approach the derivative of a parametrized ZX-diagram is another ZX-diagram. Hence we avoid the extension of the signature with formal sums. In order to tackle sums we use our original technique for the diagrammatic addition. An inductive definition of the derivative is obtained by explicit diagrammatic representation of the product rules. This inductive definition is presented in the Section 9.3.1.

In addition, in an attempt to give a ready-to-use toolbox for differentiation, we provide easy and convenient formulas to compute the derivative for the family of linear diagrams $\mathrm{ZX}(\boldsymbol{\beta})$. These formulas are non-inductive but extremely fruitful for practical manipulations. Most of circuits for variational algorithms belong to $\operatorname{ZX}(\boldsymbol{\beta})$ and we believe that our formulas will make the analysis of them much simpler.

A similar result was obtained in an independent work Wang and Yeung, 2022. This work uses W-spiders to handle product rules. In contrast to our result, the diagrammatic differentiation presented in Wang and Yeung, 2022 associates to a ZX-diagram a diagram from another language called algebraic ZX-calculus Wang, 2020. The algebraic ZX-calculus is convenient to
represent arbitrary complex numbers, therefore their differentiation procedure can handle more general families of diagrams than $\operatorname{ZX}(\boldsymbol{\beta})$. This advantage comes with the cost of abandoning the legacy of the vanilla ZX-calculus which is by far the most popular graphical calculus for quantum computing.

### 9.2 Addition of ZX-diagrams

The addition is a natural operation on matrices from a Hilbert space. In quantum mechanics it can be interpreted as the superposition phenomenon. However, the addition is not a physical process, hence it is not reflected in the standard ZX-calculus Coecke and Kissinger, 2017.

On the other hand, for any two diagrams $D_{1}, D_{2}: n \rightarrow m$, the universality of the ZX-calculus guarantees that there exists a diagram $D: n \rightarrow m$ such that $\llbracket D \rrbracket=\llbracket D_{1} \rrbracket+\llbracket D_{2} \rrbracket$. We provide in this section a general construction for such a diagram.

As pointed out in Jeandel et al., 2019 for the definition of normal forms in the ZX-calculus, one can inductively define the addition on 'controlled' versions of the diagrams. A controlled version of a diagram $D_{0}$ is roughly speaking a diagram with an extra input such that when this extra input is set to $|1\rangle$ the diagram behaves as $D_{0}$ and when it is set to $|0\rangle$ the diagram behaves as a neutral diagram. To build a controlled version of a diagram our construction pass by controlled states originally introduced in Jeandel et al., 2019. In a nutshell, our procedure for addition consists in a way to represent every ZX-diagram by such a state.

### 9.2.1 Controlled states

Definition 9.2.1 (Controlled state Jeandel et al., 2019) A ZX-diagram $C: 1 \rightarrow n$ is a controlled state if $\llbracket C \rrbracket|0\rangle=\sum_{x \in\{0,1\}^{n}}|x\rangle=\llbracket \underbrace{\bigcirc \ldots \bigcirc}_{n} \rrbracket$.

Intuitively, a controlled state is a way to encode the state $\llbracket C \rrbracket|1\rangle$. A controlled state with no outputs is called controlled scalar. For instance, the scalar 0 is encoded by a controlled scalar $C_{0}=\oiint$. Indeed, we can explicitly check that $\llbracket C_{0} \rrbracket|0\rangle$ is a uniform superposition:

$$
\| \begin{array}{ll}
\otimes 2  \tag{9.9}\\
1 & 0 \\
\hline 7.4 .6
\end{array} \quad\left[\begin{array}{l}
=-1 \\
-1
\end{array}\right]=1
$$

Verifying $\llbracket C_{0} \rrbracket|1\rangle=0$ completes the demonstration:

$$
\begin{equation*}
\llbracket Q^{\otimes 2} \quad \bigcirc_{\Omega}^{\pi} \rrbracket \rrbracket=\llbracket 8^{\otimes 2} \rrbracket \times \llbracket \pi \rrbracket=\llbracket Q^{\otimes 2} \rrbracket \times\left(1+e^{i \pi}\right)=0 \tag{9.10}
\end{equation*}
$$

We remark that it is practical to use graphical proofs instead of matrix computation to manipulate controlled states. Strictly speaking, we verify the condition $\llbracket C \rrbracket|0\rangle=\sum_{x \in\{0,1\}^{n}}|x\rangle$ by
 We will also extensively use the controlled scalars for the values $\frac{1}{\sqrt{2}}$ and 2 .

Lemma 9.2.2 (Jeandel et al., 2019|) Diagrams $C_{1 / \sqrt{2}}=\bigcirc_{-\frac{\pi}{4}}^{4}$ and $C_{2}=\boldsymbol{\wedge}$ are controlled states such that


Controlled states have nice properties that allows to perform element-wise addition and tensor product of corresponding vectors:

Lemma 9.2.3 (Sum and tensor product Jeandel et al., 2019) For any controlled states $C_{1}, C_{2}: 1 \rightarrow n$ and $C_{3}: 1 \rightarrow m$ the diagrams:

are controlled states such that $\llbracket C_{+} \rrbracket|1\rangle=\llbracket C_{1} \rrbracket|1\rangle+\llbracket C_{2} \rrbracket|1\rangle$ and $\llbracket C_{\otimes} \rrbracket|1\rangle=\llbracket C_{1} \rrbracket|1\rangle \otimes \llbracket C_{3} \rrbracket|1\rangle$.

We note that the lemma 7.4 .14 states precisely the commutativity property of the gadget


With the lemma 9.2 .3 we can add two controlled states. More precisely, if one of the controlled states encodes a state diagram $D_{1}: 0 \rightarrow n$ and the second controlled state encodes $D_{2}: 0 \rightarrow n$ the diagram $\underset{\cdots}{D_{+}}=\frac{\sigma_{+}}{C_{+}}$( returns a controlled state we can directly proceed to the addition of new terms. For $N$ diagrams $D_{1}, \ldots, D_{N}$ encoded by the controlled versions $C_{1}, \ldots, C_{N}$ the diagram

correspond to the matrix $\llbracket \boldsymbol{D} \rrbracket=\llbracket D_{1} \rrbracket+\cdots+\llbracket D_{N} \rrbracket$.
We remark that for each number of output wires $n$ the zero-element $C_{0}^{n}: 1 \rightarrow n$ is a controlled state $\wp \bigcirc \underbrace{\bigcirc \ldots \bigcirc}_{n}$. Indeed, we can prove the following lemma:
Lemma 9.2.4 For any controlled state $C: 1 \rightarrow n$


Proof: We start by proving the first equality. The left-hand side is:


The top part of the diagram can be transformed as follows:


The statement follows from (S2) and (S1).
For the second equality we have:


The lemma follows from the definition of the controlled state.
The work Yeung et al., suggests another definition of the controlled state. Essentially, the difference is the output of the zero-case. In their definition a controlled state $\tilde{C}: 1 \rightarrow n$,
when applied to the zero-input is transformed to $8^{\otimes n}$, superposition $\underbrace{\bigcirc \ldots \text {. }}_{n}$. Diagrammatically, the condition may be written as:

$$
\begin{equation*}
\frac{B Q}{\tilde{C}^{B} \cdot B^{\otimes n} \underbrace{\bigcirc \cdots \bigcirc}_{n}} \tag{9.18}
\end{equation*}
$$

### 9.2.2 Controlizer

The lemma 9.2 .3 provides a way to obtain a sum of two diagrams in a controlled state form. In order to extend the addition to arbitrary diagrams we introduce controlizers - maps that associate diagrams with the corresponding controlled states. In what follows we denote by ZX $[n, m]$ a family of ZX-diagrams with $n$ inputs and $m$ outputs.

The formal definition of the controlizer is:

Definition 9.2.5 (Controlizer) We say that a map $C: \mathrm{ZX}[n, m] \rightarrow \mathrm{ZX}[1, n+m]$ that associates to every diagram $D: n \rightarrow m$ a diagram $C(D): 1 \rightarrow n+m$ is controlizer if the following conditions hold for any ZX-diagram $D$ :

- $C(D)$ is a controlled state


There are several maps that satisfy the definition of controlizer. For instance, the map presented in Jeandel et al., 2019 that associates to each diagram its normal form is in fact a controlizer. For the alternatively defined controlled state 9.18), the work Yeung et al., shows that for each $D: 0 \rightarrow n$ the corresponding controlled state $C(D): 1 \rightarrow n$ can be constructed from $2^{n}$ coefficient of the algebraic normal form.

In our work Jeandel et al., 2022 we suggested an approach that doesn't rely on normal forms. In our approach the controlizer is defined by induction over the composition and the tensor product. Importantly, the computational complexity of our inductive procedure is polynomial on the number of elementary terms in the decomposition of the diagram on composition and tensor products of elementary generators. As a result, the final diagram has polynomial size on the number of generators and wires in the input diagram.

Definition 9.2.6 (Inductive controlizer) We define the map $C: \operatorname{ZX}[n, m] \rightarrow \mathrm{ZX}[1, n+m]$ that associates to each diagram $D: n \rightarrow m$ a diagram $C(D): 1 \rightarrow n+m$ :

- For the generators ${ }^{(\beta)}, \uparrow, \downarrow, \cap, \cup, \aleph:$

- Generators @ and @ can be decomposed as follows using the above generators:
- For tensor product $D_{\otimes}=D_{2} \otimes D_{1}$ and composition $D_{\circ}=D_{3} \circ D_{1}$ where $D_{1}: n \rightarrow m$, $D_{2}: k \rightarrow l$ and $D_{3}: m \rightarrow k:$


We remark that the output of the inductive controlizer is not unique as it depends on the decomposition order.

Lemma 9.2.7 The inductive controlizer from the definition 9.2.6) satisfies the definition 9.2.5) of controlizer.
 and both compositions. The case of $\bigodot$ and $\bigodot$ is a direct consequence of the axiom (S1).

First, we inductively prove that for all diagrams $D$ the diagram $C(D)$ is controlled state.

- Generators:



- Tensor product: it directly follows from the lemma 9.2 .3$)$ that if $C\left(D_{1}\right)$ and $C\left(D_{2}\right)$ are controlled states, then the diagram $C\left(D_{1} \otimes D_{2}\right)$ from 9.2.6) is controlled state.
- Composition: let's assume that $C\left(D_{1}\right)$ and $C\left(D_{2}\right)$ are controlled states. We show that $C\left(D_{1} \circ D_{2}\right)$ is also a controlled state:

$$
C\left(D_{1} \circ D_{2}\right) \circ 8 \mathrm{~B}=\underbrace{\left\lvert\, \begin{array}{ll}
\mid C\left(D_{1}\right) \\
\cdots & \cdots \\
C\left(D_{3}\right) \\
|r| & g^{82 m}
\end{array}\right.}_{n}
$$



$$
\begin{equation*}
\underset{\boxed{7.4 .6}}{=} \underbrace{\bigcirc \ldots \bigcirc}_{n} \underbrace{\bigcirc \ldots \uparrow}_{k} \tag{9.27}
\end{equation*}
$$

Now we show that the map $C$ satisfy the property A.33, i.e. it correctly encodes the input diagram.

- Generators:






We obtain the property A.33) for generators by swapping directions of outputs and multiplying by $\overbrace{}^{\otimes n+m}$

- Composition: Let's assume that the property A.33) holds for diagrams $D_{1}: n \rightarrow m$ and $D_{2}: m \rightarrow k$.


- Tensor product: Let's assume that the property A.33) holds for diagrams $D_{1}: n \rightarrow m$ and $D_{2}: k \rightarrow l$.


$$
\begin{equation*}
=\frac{\overbrace{|\cdots|}^{D_{2} \mid}}{\overbrace{|\cdots|}^{|\cdots|}} \frac{\overbrace{|\cdots|}^{D_{2}}}{\frac{n}{D_{1}}} \frac{\overbrace{m}^{|\cdots|}}{\frac{\mid \cdots}{|c|}} \tag{9.34}
\end{equation*}
$$

To give a flavor of our computation process we show how to obtain a controlled version of $\stackrel{\pi}{\pi}$

$\overline{=}$



## Addition of arbitrary diagrams

The controlizer allows to map every two ZX-diagrams $D_{1}: n \rightarrow m$ and $D_{2}: n \rightarrow m$ to controlled states $C_{1}: 1 \rightarrow n+m$ and $C_{2}: 1 \rightarrow n+m$. The lemma 9.2 .3 provides a way to obtain a sum of resulting controlled states. By combining this two results we can recover the diagram for addition of $D_{1}$ and $D_{2}$. The addition protocol is resumed in the following theorem:

Theorem 9.2.8 For diagrams $D_{1}: n \rightarrow m$ and $D_{2}: n \rightarrow m$ the diagram

is such that $\llbracket D_{+} \rrbracket=\llbracket D_{1} \rrbracket+\llbracket D_{2} \rrbracket$.
Proof: [Proof (theorem 9.2.8]] The theorem follows from the definition of controlizer and 9.2.3.
We illustrate the diagrammatic addition with a simple example.
Example 9.2.9 Using 9.2.8, we construct a diagram $D$ corresponding to the addition of $\cap$ and $\overbrace{}^{\pi}$.


In this example we simplified the diagram with graphical rewriting to obtain a known result:

$$
\begin{equation*}
\llbracket \curvearrowright \rrbracket+\llbracket \overbrace{}^{\pi}|\rrbracket=(|00\rangle+|11\rangle)+(|01\rangle+|10\rangle)=2|++\rangle=\llbracket \bigcirc \bigcirc \rrbracket \tag{9.38}
\end{equation*}
$$

### 9.2.3 Relation to other results

As a matter of fact, our inductive procedure leads to large diagrams, even when the initial diagrams are fairly simple. Moreover, resulting diagrams don't resemble the initial ones. This defect is however consistent with the intuition. Indeed, contrary to sequential and parallel compositions, the addition is not a physical operation, hence it is not surprising that can't be easily incorporated in the framework of ZX-diagrams. On the positive side, we can rely on the powerful equation theory of the ZX-calculus to simplify, when it is possible, the diagrams representing the sum of diagrams.

An alternative approach for the addition in the algebraic ZXW-calculus was presented in Yeung et al., It also relies on the controlled froms of the diagram. We remark that the circuit protocol (9.2) for addition of unitaries presented in Childs and Wiebe, 2012 also requires controlled versions of unitary transformation.

The controlled forms used in Yeung et al., are different from ours. Notably, contrary to our approach that accepts diagrams with an arbitrary number of input and output wires, the controlled forms in Yeung et al., are defined only for two families of diagrams.

The first family consists of state diagrams $\mathrm{ZX}[0, n]$. The work Yeung et al., suggests to construct a controlled state from the normal form Wang, 2020. Importantly, for a state diagram with $n$ outputs the normal form has an exponential size on $n$.

The second family of diagrams treated in Yeung et al., , includes diagrams ZX $[m, m]$ that correspond to square matrices. A procedure suggested in Yeung et al., for controlled matrices requires the decomposition of the initial diagram $D$ on a product of diagrams for so-called elementary matrices Wang and Yeung, 2021. As far as we have understood, in the general case the decomposition on elementary matrices requires an explicit computation of the matrix corresponding to the diagram. Therefore, although for given controlled forms the addition from Yeung et al., is done diagrammatically with the help of $W$-spider, the computation of the controlled forms pass by the semantics. In contrast, in our inductive procedure the entire process is maintained in the graphical framework.

However, in some special cases of practical importance the approach suggested in Yeung et al., is beneficial. For instance, as was shown in Yeung et al., for local Hamiltonians made out of Pauli tensors the decomposition on elementary matrices is straightforward. As a consequence, the controlled versions are easy to compute and the overall procedure leads to nice-looking diagrams. Moreover, considering the diagrams from more expressive algebraic ZXW-calculus leads to a simpler representation of linear combinations of diagrams with arbitrary complex coefficients.

### 9.3 Differentiation of ZX-diagrams

A basic definition of the derivative is the derivative for a smooth function $f: \mathbb{R} \rightarrow \mathbb{R}$ :

$$
\begin{equation*}
\frac{\partial f}{\partial x}\left(x_{0}\right)=\lim _{\delta \rightarrow 0} \frac{f\left(x_{0}+\delta\right)-f\left(x_{0}\right)}{\delta} \tag{9.39}
\end{equation*}
$$

For a function on multiple variables this definition is extended to a gradient. In the traditional calculus the derivative usually is computed from the decomposition of the function $f: \mathbb{R} \rightarrow \mathbb{R}$ on a sum, product and composition of elementary functions. For elementary functions the derivatives are explicitly computed from the definition 9.39 . Derivatives for complex functions are computed from the linearity and the product rules. The linearity and the product rules that are derived from the definition.

A parameterized ZX-diagram $D(\beta)$ can be understood as a function $f: \mathbb{R} \rightarrow$ ZX that associated to each real value $\beta \in \mathbb{R}$ a ZX-diagram. In the context of variational algorithms, we consider the derivatives over the parameters. We remark that as the sums and subtractions are unnatural for ZX-diagrams we avoid to define the derivative as a limit of a finite difference 9.39.

In order to keep the notations simple in what follows we restrict our attention to one-variable diagrams. Moreover, we will consider only the family of linear diagrams 7.5 . In our definition of linear diagrams the phases are affine functions with integer coefficients before parameters. We denote the corresponding matrices by $\mathcal{M}(\beta)$. The elements of matrices in $\mathcal{M}(\beta)$ belong to the ring generated by complex numbers and the smooth functions $e^{i \mathcal{A}}$ where $\mathcal{A}=\{k \beta \mid k \in \mathbb{Z}\}$ is the group of affine functions with integer coefficients.

The sum for the matrices in $\mathcal{M}(\beta)$ is naturally defined by entrywise addition. Therefore, we can define the derivative in $\mathcal{M}(\beta)$ by specifying the general definition for monoidal categories with sums given in Toumi et al., 2021:

Definition 9.3.1 (Derivative in $\mathcal{M}(\beta))$ A derivative $\partial_{M}$ in $\mathcal{M}(\beta)$ is a linear unary operator on $\mathcal{M}(\beta)$ that associated to a matrix $L \in \mathcal{M}(\beta), L: n \rightarrow m$ another matrix $L^{\prime} \in \mathcal{M}(\beta), L^{\prime}:$ $n \rightarrow m$ of the same dimensions. It satisfies the following axioms:
○-product rule : $\partial_{M}[A \circ B]=\partial_{M}[A] \circ B+A \circ \partial_{M}[B]$
$\otimes$-product rule : $\partial_{M}[A \otimes B]=\partial_{M}[A] \otimes B+A \otimes \partial_{M}[B]$
Rather than defining the derivative as a limit of finite difference and deriving the product rules, the work Toumi et al., 2021 suggests an axiomatic definition that proclaims the product rules as the fundamental property of the derivative operator. Furthermore, the work Toumi et al., 2021 shows how their general definition can be used to define derivatives for bags of $Z X$-diagrams.

Although our theorem 9.2 .8 provides a fully diagrammatic way for addition of diagrams, we avoid to extend the language with a formal introduction of a sum operator in order to circumvent unnecessary complications. Therefore, the general definition from Toumi et al., 2021 doesn't apply. Instead, we suggest an alternative semantics for the derivative in the ZX-calculus. In place of product rules we require derivative to satisfy the property of diagrammatic differentiation. This property states the coherence between the derivative of a diagram and the derivative of the corresponding linear map:

Definition 9.3.2 (ZX-derivative) A derivative $\partial_{\mathrm{ZX}}: \mathrm{ZX}[n, m] \rightarrow \mathrm{ZX}[n, m]$ is an unary operator that commutes with the standard interpretation:

$$
\begin{equation*}
\llbracket \partial_{Z X} D \rrbracket=\partial_{M} \llbracket D \rrbracket \tag{9.40}
\end{equation*}
$$

where $\partial_{M}$ is the derivative 9.3.1) in $\mathcal{M}(\beta)$.
We remark that even if we restricted the consideration to diagrams over one variable, all results may be easily extended to the case of partial derivatives $\partial_{\beta_{i}}$ for linear diagrams with an arbitrary number of variables.

### 9.3.1 Diagrammatic differentiation with controlizers

The derivative in $\mathcal{M}(\beta)$ is defined through product rules that involve sums. In this section we show how to incorporate these sums in the diagrammatic framework using controlizers.

In what follows we denote by $C$ any map that satisfies the definition 9.2 .5 of controlizer.

Definition 9.3.3 We call C-derivative a map $\Delta: \operatorname{ZX}(\beta) \rightarrow \operatorname{ZX}(\beta)$ that associates to a parameterized diagram $D: n \rightarrow m$ another diagram $\Delta(D): 1 \rightarrow n+m$ defined as follows:

- Generators: For parametrized spiders:

For all generators $g: n \rightarrow m$ that are independent on $\beta \Delta(g)=B \underbrace{\ldots}_{n+m}$.

- Tensor product: for $D_{1}: n \rightarrow m$ and $D_{2}: k \rightarrow l$ the diagram $\Delta\left(D_{1} \otimes D_{2}\right)$ is:

- Composition: for $D_{1}: n \rightarrow m$ and $D_{2}: m \rightarrow k$ the diagram $\Delta\left(D_{2} \circ D_{1}\right)$ is:


A simple verification shows that for any generator $G$ the $C$-derivative $\Delta(G)$ is a controlled state. From the lemma 9.2 .3 it follows that for every diagram $D: n \rightarrow m$ the $C$-derivative $\Delta(D): 1 \rightarrow n+m$ is a controlled state. Moreover, as was proven in the lemma 9.2 .4 the $C$-derivative $\underbrace{\underbrace{\ldots}_{n+m}}_{n+m}$ of generators that are independent on $\beta$ is precisely the zero-element $C_{0}^{n+m}$ for the addition of controlled states. The same lemma implies, in particular, that the $C$-derivative of any diagram $D: n \rightarrow m$ that is independent on $\beta$ is $\delta \underbrace{\ldots \ldots}_{n+m}$.

Similarly to controlizers, a step-by-step application of the map $\Delta$ may lead to different diagrams depending on the order of decomposition on tensor products and compositions. However,
all possible outputs are semantically equivalent and by completeness of ZX -calculus are equivalent as diagrams.

By combining $C$-derivatives and controlizers we can reproduce the semantics of product rules in a graphical representation.

Definition 9.3.4 Given a $C$-derivative $\Delta$, let $\partial_{C}: \operatorname{ZX}(\beta) \rightarrow \operatorname{ZX}(\beta)$ be the unary operator such that for any diagram $D: n \rightarrow m$

Theorem 9.3.5 The operator $\partial_{C}$ satisfies the definition 9.40) of diagrammatic differentiation.
Proof: [Proof (theorem 9.3.5)] We give the proof by induction over tensor product and composition.

- Generators: Firstly, we show that $\left.\partial_{M}[[\boxed{\beta}]]\right]=i e^{i \beta}|-\rangle$ and $\left.\partial_{M}[\llbracket(-\beta)]\right]=-i e^{-i \beta}|-\rangle$. Indeed, we have:

$$
\begin{array}{r}
\partial_{M} \llbracket[(\beta)]=\partial_{M}\left(|+\rangle+e^{i \beta}|-\rangle\right)=i e^{i \beta}|-\rangle \\
\partial_{M} \llbracket-\frac{-\beta}{-} \|=\partial_{M}\left(|+\rangle+e^{-i \beta}|-\rangle\right)=-i e^{-i \beta}|-\rangle \tag{9.47}
\end{array}
$$

For the diagrams $\partial_{C}(ß)$ and $\partial_{C}\left(\Theta^{(\beta)}\right)$ we verify the property of diagrammatic differentiation:

For a generator $g: n \rightarrow m$ that doesn't depend on $\beta$ :

- Tensor product: Let's assume that $\llbracket \partial_{C}\left(D_{1}\right) \rrbracket=\partial_{M} \llbracket D_{1} \rrbracket$ and $\llbracket \partial_{C}\left(D_{2}\right) \rrbracket=\partial_{M} \llbracket D_{2} \rrbracket$ for $D_{1}: n \rightarrow m$ and $D_{2}: l \rightarrow k$. For $D=D_{1} \otimes D_{2}$ :


By applying the definition 9.3 .3 of the $C$-derivative for the tensor product we obtain the diagram


We denote by

| $T$   <br> $\underbrace{}_{n+m} \cdots$ $\mid \cdots$  |
| :---: | :---: | the part of the diagram surrounded by the red frame. Ac-

$$
n+m \quad k+l
$$

cording to the definition, $C$-derivatives $\Delta\left(D_{1}\right), \Delta\left(D_{2}\right)$ and controlizer outputs $C\left(D_{1}\right)$, $C\left(D_{2}\right)$ are controlled states. Therefore, from the lemma 9.2 .3 we get:

Meanwhile, by bending wires we see that the resting part of the diagram 9.52 is:


The sum in $\mathcal{M}(\beta)$ is distributive over the composition. Therefore,



The claim follows from the induction assumption and the definition of controlizer.

- Composition: Following the same reasoning as for the tensor product we obtain for $D_{1}: n \rightarrow m$ and $D_{2}: m \rightarrow l$

where the diagram $T$ is exactly the same as in 9.53 ). On more time, by distributivity of the sum in $\mathcal{M}(\beta)$ over the composition we obtain

The claim follows from the definition of controlizer an the induction assumption.

We consider a simple example to illustrate how to perform the inductive procedure.
Example 9.3.6 We apply the definition 9 9.3.4) to the simple scalar diagram $\underbrace{\beta}_{-\pi}={ }_{(\pi)}^{\beta}$

already found in 9.35 . We obtain a diagram for $C$ (



(9.61)



### 9.3.2 Formula for derivatives in $\mathrm{ZX}(\beta)$

Although perfectly correct, the differentiation procedure described above leads to very puzzling output even for small diagrams (see the example 9.3.6). In this section we provide a simpler approach to obtain the derivative of a diagram in $\mathrm{ZX}(\beta)$ written in a special form. We formalize it in definitions $\partial_{\mathrm{ZX}}$ and $\partial_{P}$ of unary operators that satisfy the property of diagrammatic differentiation 9.40 .

Firstly, we remark that even in the inductive procedure the effort can be significantly reduced if we separate the parts that depend on the parameter from the rest of the diagram that is constant on $\beta$. Indeed, for the constant diagrams $A: k \rightarrow l$ and $B: m \rightarrow k$ the $C$-derivative is the controlled version of the zero diagram $C_{0}^{k+l}$ and $C_{0}^{m+k}$. Therefore, for a parameterized diagram $D(\beta): n \rightarrow m$ the lemma 9.2.4 implies :


This expressions are significantly simple that the general case 9.3 .3 where both parts of the compositions depend on $\beta$.

We want to further explore this property to find simplified formulas for the derivative. Let's denote by $X_{\beta}(n, m)$ diagrams $\underbrace{\beta, \ldots{ }^{\beta} \underbrace{-\beta}_{m} \ldots)^{-\beta}}_{n}$ from $\mathrm{ZX}(\beta)$.

From the (S1) and (H) rules and the paradigm Only topology matters follows:
Claim 9.3.7 Using the rules of $Z X$ calculus, each diagram $D(\beta)$ from $Z X(\beta)$ may be transformed into the form

where $n, m$ are some integer numbers and $D_{1}, D_{2}$ are constant with respect to $\beta$. We call diagrams in this form $\beta$-factored.

A rigorous demonstration of the claim 9.3 .7 may be found in Jeandel et al., 2018b.
The diagram $X_{\beta}(n, m)$ has a particularly regular form, so we expect it to have a nice-looking derivative. We would like to develop the intuition from the semantics $\partial_{M} \llbracket X_{\beta}(n, m) \rrbracket$.

## Intuition

We denote by $|\beta\rangle=|0\rangle+e^{i \beta}|1\rangle$ the state $\llbracket\left(\begin{array}{|}(\beta) \\ \uparrow\end{array}\right]$ and by $|-\beta\rangle=|0\rangle+e^{-i \beta}|1\rangle$ the state $\llbracket-\beta$.
We also introduce the notation:

$$
\begin{equation*}
|\beta(n, m)\rangle=|\beta\rangle^{\otimes n} \otimes|-\beta\rangle^{\otimes m}=\sum_{x \in\{0,1\}^{n+m}} e^{i \beta\left(\left|x^{+}\right|-\left|x^{-}\right|\right)}\left|x^{+}\right\rangle \otimes\left|x^{-}\right\rangle \tag{9.65}
\end{equation*}
$$

where $x^{+} \in\{0,1\}^{n}$ corresponds to first $n$ qubits that are in the state $|\beta\rangle^{\otimes n}$ and $x^{-} \in\{0,1\}^{m}$ corresponds to the last $m$ qubits that are in the state $|-\beta\rangle^{\otimes m}$. The notation $|x|:\{0,1\}^{k} \rightarrow \mathbb{Z}_{+}$ stands for the Hamming weight of the binary string $x$.

By applying entrywise differentiation to the vector $|\beta(n, m)\rangle$, we obtain the following result:

$$
\begin{equation*}
\partial_{M}(|\beta(n, m)\rangle)=\sum_{x \in\{0,1\}^{n+m}} i\left(\left|x^{+}\right|-\left|x^{-}\right|\right) e^{i \beta\left(\left|x^{+}\right|-\left|x^{-}\right|\right)}\left|x^{+}\right\rangle \otimes\left|x^{-}\right\rangle \tag{9.66}
\end{equation*}
$$

We observe that $\partial_{M}(|\beta(n, m)\rangle)=i M_{\Delta}|\beta(n, m)\rangle$ where $M_{\Delta}: 2^{n+m} \rightarrow 2^{n+m}$ is a linear map such that $M_{\Delta}:|x\rangle \rightarrow\left(\left|x^{+}\right|-\left|x^{-}\right|\right)|x\rangle$.

In order to find a diagram $D_{M}$ for the matrix $M_{\Delta}$ we introduce the controlled-triangle $D_{\lambda}: 2 \rightarrow 2$. Semantically, we want the diagram $D_{\lambda}$ to apply the triangle $\boldsymbol{A}$ to the second input if the first one is in the $|1\rangle$ state and do nothing otherwise. We denote the matrix corresponding
to $D_{\lambda}$ by $M_{\lambda}$. Crucially, when the matrix $M_{\lambda}$ is applied to a vector $|b\rangle \otimes(k|0\rangle+|1\rangle)$ for $b \in\{0,1\}$ the coefficient $k$ is modified depending on the value of $b$ :

$$
\begin{equation*}
|b\rangle \otimes(k|0\rangle+|1\rangle) \xrightarrow{M_{\lambda}}|b\rangle \otimes((k+b)|0\rangle+|1\rangle) \tag{9.67}
\end{equation*}
$$

Similarly, when the input vector is $|b\rangle \otimes(k|0\rangle-|1\rangle)$ the value of $b$ is subtracted from $k$ :

$$
\begin{equation*}
|b\rangle \otimes(k|0\rangle+|1\rangle) \xrightarrow{M_{\lambda}}|b\rangle \otimes((k-b)|0\rangle-|1\rangle) \tag{9.68}
\end{equation*}
$$

We remark the matrix $M_{\lambda}$ is not norm-preserving, so it doesn't correspond to a valid quantum gate. However, it doesn't matter as ZX-diagrams are universal for general linear maps of dimensions $2^{n} \times 2^{m}$.

We prove that $D_{\lambda}=Q_{-}$( is in $|0\rangle$ and in $|1\rangle$ states.

For the $|0\rangle$-case we have:


When the first input is in the $|1\rangle$ state we get:


It turn out that the matrix $M_{\Delta}$ can be written as an evolution in $n+m+1$ space with one ancilla input followed by a projection on $n+m$ dimensional space. The ancilla input is used to accumulate the coefficient $\alpha_{x}$ of the basis state $|x\rangle$. We suggest a construction in which the coefficient $\alpha_{x}$ is recovered from the projection $\left\langle 0 \mid \psi_{a}\right\rangle$ of the state $\left|\psi_{a}\right\rangle$ of the ancilla input.

To be exact, we design an iterative procedure that multiplies the input vector $\left|b_{0}, \ldots, b_{m+1}\right\rangle$ by the matrix $M_{\Delta}$. Initially, we set $\left|\psi_{a}\right\rangle=|1\rangle$. Then, for $i \in[1, \ldots, n]$ we apply the transformation $M_{\lambda}$ to $\left|b_{i}\right\rangle\left|\psi_{a}\right\rangle$ controlled by the input $i$. After $n$ applications the state $\left|\psi_{a}\right\rangle$ becomes $\left|\psi_{a}\right\rangle=\left(\operatorname{sum}_{i=1}^{n} b_{i}\right)|0\rangle+|1\rangle$. To proceed further we apply $Z$ matrix to $\left|\psi_{a}\right\rangle$ and obtain $\left|\psi_{a}\right\rangle=$ $\left(\operatorname{sum}_{i=1}^{n} b_{i}\right)|0\rangle-|1\rangle$. Then the state $\left|\psi_{a}\right\rangle$ is transformed to $\left|\psi_{a}\right\rangle=\left(\sum_{i=1}^{n} b_{i}-\sum_{i=n+1}^{m} b_{i}\right)|0\rangle-|1\rangle$ by iterative application of $M_{\lambda}$ to $\left|b_{i}\right\rangle\left|\psi_{a}\right\rangle$. Finally, the coefficient $\alpha_{x}$ is:

$$
\begin{equation*}
\alpha_{x}=\left\langle 0 \mid \psi_{a}\right\rangle=\left(\sum_{i=1}^{n} b_{i}-\sum_{i=n+1}^{m} b_{i}\right) \tag{9.71}
\end{equation*}
$$

As the original input wires are used only in control, the output of such procedure is $\alpha_{x}\left|b_{0}, \ldots, b_{n+m}\right\rangle$.
This procedure can be directly translated in a diagram $D_{\lambda}$;


The diagram for the derivative of $X_{\beta}(n, m)=\underbrace{\beta} \ldots \underbrace{\beta-\beta} \ldots$ can be obtained by conjugating every input with Hadamard boxed:

By applying the rule (H) to the final we obtain

## Diagrammatic differentiation of $\beta$-factored forms

Motivated by our the simplicity of the result 9.75 , we suggest an alternative non-inductive definition for derivatives of diagrams in $\beta$-factored forms:

Definition 9.3.8 Given a diagram $D(\beta)=D_{2} \circ\left(D_{1} \otimes X_{\beta}(n, m)\right)$ in $\beta$-factored form, let $\partial_{\mathrm{ZX}}[D]=$ $D_{2} \circ\left(D_{1} \otimes \partial_{\mathrm{ZX}}\left[X_{\beta}(n, m)\right]\right)$ where

$$
\begin{align*}
& \partial_{\mathrm{ZX}}\left[X_{\beta}(n, m)\right]=\partial_{\mathrm{ZX}}[\underbrace{(\beta) \ldots \Theta^{\beta}}_{n} \underbrace{-\beta) \ldots \Theta^{-\beta}}_{m}] \tag{9.76}
\end{align*}
$$

Although our definition is inspired from the intuitive reflection presented above, we provide a rigorous diagrammatic demonstration of the soundness of our definition. More precisely, we prove the following theorem:

Theorem 9.3.9 The operator $\partial_{\mathrm{ZX}}[-]$ defined at 9.3.8) satisfies the property of diagrammatic differentiation: for any diagram $D(\beta) \in \mathrm{ZX}(\beta)$ in $\beta$-factored form,

$$
\begin{equation*}
\llbracket \partial_{\mathrm{ZX}} D(\beta) \rrbracket=\partial_{M} \llbracket D(\beta) \rrbracket \tag{9.78}
\end{equation*}
$$

We remark that according to the the definition 9.3 .2 the derivative $D^{\prime}: n \rightarrow m$ of a diagram $D: n \rightarrow m$ that is constant on $\beta$ is such that $\llbracket D^{\prime} \rrbracket=\partial_{M} \llbracket D \rrbracket=(0)_{n \times m}$. Therefore, the theorem 9.3 .9 is a direct consequence of the following lemma:

Lemma 9.3.10 For any $n, m$ :

$$
\begin{equation*}
\llbracket \partial_{\mathrm{ZX}} X_{\beta}(n, m) \rrbracket=\partial_{M} \llbracket X_{\beta}(n, m) \rrbracket \tag{9.79}
\end{equation*}
$$

Proof: [Proof (lemma 9.3.10]] We prove the lemma by induction. The demonstration is done for $n=n+1$, the proof for $m=m+1$ is directly obtainable in the same way.

Base: We already know that $\partial_{M}\left[\llbracket(\beta) \rrbracket=i e^{i \beta}|-\rangle\right.$. Thus, for the case $n=1, m=0$ the lemma statement follows from:

Step: By induction, we assume that the equation (9.79) holds for some $n$ and $m$. We show that under this assumption $\llbracket \partial_{\mathrm{ZX}} X_{\beta}(n+1, m) \rrbracket=\partial_{M} \llbracket X_{\beta}(n+1, m) \rrbracket$.

In order to prove the induction we will proceed in three steps. Firstly we apply the product rule to the matrix $\partial_{M} \llbracket X_{\beta}(n+1, m) \rrbracket$ :

$$
\begin{align*}
\partial_{M} \llbracket X_{\beta}(n+1, m) \rrbracket & =\partial_{M} \llbracket X_{\beta}(1,0) \rrbracket \otimes \llbracket X_{\beta}(n, m) \rrbracket+\llbracket X_{\beta}(1,0) \rrbracket \otimes \partial_{M} \llbracket X_{\beta}(n, m) \rrbracket \\
& =\llbracket \partial_{\mathrm{ZX}}\left[X_{\beta}(1,0) \rrbracket \otimes X_{\beta}(n, m) \rrbracket+\llbracket X_{\beta}(1,0) \otimes \partial_{\mathrm{ZX}}\left[X_{\beta}(n, m) \rrbracket \rrbracket\right.\right. \tag{9.81}
\end{align*}
$$

where the second equality follows from the inductive assumption.
Secondly, we show the last sum in the equation (9.81) can be represented diagrammatically with a controlled state, i.e. the following claim holds:
Claim 9.3.11 We can find a controlled state $\tilde{X}: 1 \rightarrow n+1+m$ and a constant scalar $c \in \mathrm{ZX}_{\frac{\pi}{2}}$ such that


Finally, we will show that we can transform the diagram $\partial_{\mathrm{ZX}} X_{\beta}(n+1, m)$ to
 with a sequence of graphical rewrites:

## Claim 9.3.12



As the rewrite rules preserve the semantics, it follows from the claims above that:

$$
\begin{align*}
& \underset{\text { 9.3.11 }}{=} \llbracket \partial_{\mathrm{ZX}}\left[X_{\beta}(1,0)\right] \otimes X_{\beta}(n, m) \rrbracket+\llbracket X_{\beta}(1,0) \otimes \partial_{\mathrm{ZX}}\left[X_{\beta}(n, m)\right] \rrbracket \\
& \underset{9.81}{=} \partial_{M} \llbracket X_{\beta}(n+1, m) \rrbracket \tag{9.84}
\end{align*}
$$

## Proofs

While proving the claims 9.3 .11 and 9.3 .12 we will repeatedly use the following lemmas:
Lemma 9.3.13


Proof: [lemma 9.3.13] The right hand side transformation:






$\stackrel{=}{7.4 .24}$




Lemma 9.3.14 Diagrams:
are controlled states such that

$$
\begin{equation*}
\llbracket C_{ \pm \beta} \rrbracket|1\rangle=\llbracket \wp^{ \pm \beta} \rrbracket, \quad \llbracket C_{\beta}^{(n, m)} \rrbracket|1\rangle=\llbracket \Theta_{-\frac{\pi}{2}}^{(\pi)} \bigotimes^{\otimes(n+m-1)} \partial_{\mathrm{ZX}}[X(n, m)] \rrbracket \tag{9.87}
\end{equation*}
$$

Proof: [lemma 9.3.14] For the diagram $C_{\pi \beta}$ the claim follows directly from the definition. For the diagram $C_{\beta}^{(n, m)}$ we check two cases:
-



$8^{(\pi)}$
$\left\lvert\, \begin{gathered}C_{\beta}^{(n, m)} \\ \mid \cdots\end{gathered}\right.$



Proof: [claim 9.3.11
We introduce the diagram:

where $C_{ \pm \beta}$ and $C_{\beta}^{(n, m)}$ are defined in the previous lemma 9.3 .14 . From these definitions and the lemma 9 9.2.3) follows that

$$
\begin{align*}
& \llbracket \tilde{X} \rrbracket|1\rangle=\llbracket C_{\beta}^{(1,0)} \rrbracket|1\rangle \otimes \llbracket C_{+\beta} \rrbracket|1\rangle \otimes \cdots \otimes \llbracket C_{-\beta} \rrbracket|1\rangle+\llbracket C_{+\beta} \rrbracket|1\rangle \otimes \llbracket C_{\beta}^{(n, m)} \rrbracket|1\rangle \\
& =\llbracket\left(\begin{array} { c } 
{ \pi } \\
{ - \frac { \pi } { 2 } 2 } \\
{ Q ^ { 8 } }
\end{array} \mathrm { B } ^ { \otimes ( n + m ) } \rrbracket \left(\llbracket \partial_{\mathrm{ZX}}\left[X_{\beta}(1,0)\right] \otimes X_{\beta}(n, m) \rrbracket+\llbracket X_{\beta}(1,0) \otimes \partial_{\mathrm{ZX}}\left[X_{\beta}(n, m) \rrbracket \rrbracket\right)\right.\right. \tag{9.89}
\end{align*}
$$

Therefore the claim holds for the controlled state 9.88 and the scalar $c=8_{8}^{8(n+m-2)}$
Proof: [claim 9.3.12 Firstly, we simplify the expression $\tilde{X} \circ \frac{8}{T_{1}}$ :




We apply the lemma 9.3 .13 to parts in the red frame and obtain:


The claim holds for the scalar $c=\underbrace{(\pi)} 8^{\left(\frac{1}{2}\right)} 0^{(n+m-2)}$

### 9.3.3 Simplified formula for paired spiders

Variational quantum algorithms use gradients in the search for optimal parameter values. The objective minimized by these algorithms can be expressed as $\langle\psi(\beta)| H|\psi(\beta)\rangle$ where the diagram for $\langle\psi(\beta)|=(|\psi(\beta)\rangle)^{\dagger}$ is obtained out of the diagram for $|\psi(\beta)\rangle$ by flipping up side down followed by the change of signs in spiders. Therefore, parameters in the diagram for $\langle\psi(\beta)| H|\psi(\beta)\rangle$ appear in pairs

We suggest a more compact formula for diagrams in what we call pair-factored form: $D_{2}$ 。 $\left(D_{1} \otimes Y(n)\right)$. In this expression $Y_{\beta}(n)=\underbrace{\left(\begin{array}{ll}-\beta & \beta \\ \hline & \\ \hline\end{array}\right) \ldots\left(\begin{array}{ll}-\beta & B\end{array}\right)}_{n}$.

Lemma 9.3.15 The diagram:

satisfies $\llbracket \partial_{P}\left(Y_{\beta}(n)\right) \rrbracket=\partial_{M} \llbracket Y_{\beta}(n) \rrbracket$.

We prove the lemma 9.3 .15 by applying the same approach as in the proof of of the formula 9.3 .10 for individual spiders. We can then replace by (9.92) the expression 9.77) in 9.3 .8 and obtain the derivative for diagrams in pair-factored form.
Proof: [lemma 9.3.15
Base: We have $\partial_{M}\left[\begin{array}{|c|}-\beta \\ \underset{|c|}{\beta} \\ \hline\end{array}\right]=\partial_{M}\left(|++\rangle+e^{i \beta}|+-\rangle+e^{-i \beta}|-+\rangle+|--\rangle\right)=i\left(e^{i \beta}|+-\rangle-e^{-i \beta}|-+\rangle\right)$.
Simplifying the diagram 9.92 for $n=1$ we obtain:


From the definition of the standard interpretation $\llbracket \overbrace{\sim}^{T}|\rrbracket=| 00\rangle-|11\rangle=|+-\rangle+|-+\rangle$ and $\left[R_{X}(\pi-\beta) \otimes R_{X}(\beta)\right] \circ[|+-\rangle+|-+\rangle]=e^{i \beta}|+-\rangle+e^{-i \beta}|-+\rangle$. Therefore, the lemma is true for $n=1$.

Step: By induction, we assume that $\llbracket \partial_{P} Y_{\beta}(n) \rrbracket=\partial_{M} \llbracket Y_{\beta}(n) \rrbracket$. We show that under this assumption the lemma holds for $n=n+1$ by proving following claims:

Claim 9.3.16

$$
\begin{equation*}
\partial_{M} \llbracket Y_{\beta}(n+1) \rrbracket=\llbracket \partial_{P} Y_{\beta}(1) \otimes Y(n) \rrbracket+\llbracket Y_{\beta}(1) \times \partial_{P} Y(n) \rrbracket \tag{9.94}
\end{equation*}
$$

Claim 9.3.17 There exist a controlled state $\tilde{Y}$ and a scalar c such that

$$
\begin{equation*}
\llbracket c \otimes\left(\tilde{Y} \circ \mathcal{B}^{\pi}\right) \rrbracket=\llbracket \partial_{P} Y_{\beta}(1) \otimes Y_{\beta}(n) \rrbracket+\llbracket Y_{\beta}(1) \times \partial_{P} Y_{\beta}(n) \rrbracket \tag{9.95}
\end{equation*}
$$

Claim 9.3.18 $Z X \vdash \partial_{P}\left(Y_{\beta}(n+1)\right)=c \otimes(\tilde{Y} \circ \overbrace{T}^{\pi})$ where $C$ and $c$ are respectively the controlled state and scalar from the previous claim.

The claim 9.3.16) follows from the product rule for $\partial_{M}$ and from the induction assumption. Proof: [Proof (claim 9.3.17]] We can explicitely check that for each $n \in \mathbb{N}$ the following
 $\llbracket \sum_{(\pi)}^{\left(-\frac{\pi}{2}\right.} 88^{\otimes 2 n-2} \partial_{p}\left[Y_{\beta}(n)\right] \rrbracket:$

Indeed,




 $C_{\beta}$ and $C_{-\beta}$ were defined in the lemma 9.3 .14 . Due to the lemma 9.2.3),

$$
\begin{align*}
& \left.+\llbracket \otimes^{\otimes 2}(-\beta) \mathbb{\beta}\right) \rrbracket \otimes\left[\begin{array}{l}
(\pi) \\
-\frac{\pi}{2} \\
\hline
\end{array}\right) 88^{\otimes 2 n-2} \partial_{p}[Y(n)] \rrbracket \\
& =\llbracket\left(\begin{array}{c}
\pi \\
-\frac{\pi}{2} \\
2
\end{array}\right), O^{\otimes 2 n} \rrbracket\left[\llbracket \partial_{P} Y(1) \otimes Y(n) \rrbracket+\llbracket Y(1) \times \partial_{P} Y(n) \rrbracket\right] \tag{9.96}
\end{align*}
$$

Therefore, the claim holds for $C=\hat{C}$ and $c=\stackrel{\pi}{\pi} 90^{2 n}$
Proof: [Proof (claim 9.3.18)] Firstly, we simplify the diagrams $\tilde{Y} \circ \frac{\pi}{\pi}$ :



We subsequently apply the lemma 9.3 .13 ) to parts in the red frames:


The last expression multiplied by the constant $c=\frac{\pi}{\substack{(2) \\ 2}} 80^{2 n}$ is equal to $\partial_{P} Y(n+1)$. Induction is proven.

It is possible to extend 9.3 .15 to find the derivative for $X_{\beta}(n, m)$ when $n \neq m$. Indeed, using the fact that $O+ \pm=-\quad$ we can balance the number of $\beta$ and $-\beta$. For instance, if $n>m$ : $\partial_{P}\left(X_{\beta}(n, m)\right)=\sigma \circ\left[\bigotimes^{\otimes n-m} \begin{array}{l}\wp_{n-m}^{\partial_{P}\left(Y_{\beta}(n)\right)} \\ \wp_{2 m}^{|\cdots|}\end{array}\right]$ where $\sigma$ is some wire permutation and


Example 9.3.19 We apply 9.3.15 to the same diagram as in 9.3.6:

We observe that using the formula for the pair-factored form 9.92 we obtain a much more compact result than with the inductive procedure (see the example 9.3.6). Even if both approaches lead to diagrams with the same semantics, in practice the diagrams obtained with the formulas 9.92 and 9.77 ) are less verbose. For this reason they are easier to manipulate.

### 9.3.4 Relation to other results

Contrary to derivatives defined in Toumi et al., 2021 and Zhao and Gao, 2021 in our approaches a derivative of a ZX-diagram is another ZX-diagram. In other words, we circumvent the necessity of introducing formal sums. In the inductive approach we proceed by integration of the product rules using controlizers.

We also observe that as the derivative of a constant diagram is trivial, it is beneficial to 'factor-out' the part that depends on the parameter. We adopted this approach in formulas for diagrams in $\beta$-factored form and pair-factored forms. To derive the formulas, we got the intuition in the desired semantics of the diagrams. Thenceforth, our formulas were rigorously proven by induction.

A result similar to our simplified formula for $\beta$-factored forms was independently derived in Wang and Yeung, 2022. The major difference between our formula and the method shown in Wang and Yeung, 2022 is the considered language. Indeed, in our work we operate ZX-diagrams while Wang and Yeung consider diagrams from more expressive algebraic ZX-calculus.

The difficulty to represent derivatives for non-linear diagrams follows from the fact that there is no a simple way to represent real numbers in the vanilla ZX-calculus. In the algebraic ZX-calculus this restriction is removed. As a consequence, when an algebraic ZX-diagram is parameterized by an arbitrary derivable function $f(x)$, the differentiated algebraic ZX-diagram is parametrised by $f^{\prime}(x)$.

### 9.4 Diagrammatic representation of Ising Hamiltonians

Traditionally, in the variational algorithms the optimization objective is encoded by a Hamiltonian. For combinatorial optimization problems the Hamiltonian is typically an Ising model:

$$
\begin{equation*}
H=\sum_{i=1}^{n} h_{i} Z_{i}+\sum_{1 \leq i<j \leq n} h_{i, j} Z_{i} Z_{j} \tag{9.101}
\end{equation*}
$$

There is no direct way to transform the definition of the Hamiltonian (9.101) to a ZX-diagram. Indeed, Hamiltonian is a non-unitary matrix equal to a sum of Pauli gates that is inherently difficult to represent as a diagram.

The work Yeung et al., suggests a convenient way to represent Hamiltonians that are linear combinations of Pauli tensors in algebraic ZXW-calculus. In Jeandel et al., 2022 we adopted a different approach that use our diagrammatic differentiation together with the Stone's theorem. The Stone theorem relating the derivative of a unitary group $e^{i \gamma H}$ to its generator $H$ was used before in Toumi et al., 2021. However, in Toumi et al., 2021 the derivatives were defined as formal sums of diagrams. As a consequence, the work Toumi et al., 2021 didn't suggest a single diagram representation for a Hamiltonian but rather a representation as a sum of diagrams.

We remark that for an Ising Hamiltonians $H$ the diagram $D_{U}(\beta)$ of the linear map $U(\beta)=$ $e^{i \beta H}$ is easy to find. For Hamiltonians with integer coefficients the matrix $U(\beta)=e^{i \beta H}$ belongs to $\mathcal{M}(\beta)$. It satisfies the definition of strongly continuous one-parameter unitary group:

Definition 9.4.1 (Unitary group Toumi et al., 2021]) A one-parameter unitary group is a unitary matrix $U: n \rightarrow n$ in $\mathcal{M}(\beta)$ with $U(0)=i d_{n}$ and $U(\beta) U\left(\beta^{\prime}\right)=U\left(\beta+\beta^{\prime}\right)$ for all $\beta, \beta^{\prime} \in \mathbb{R}$. It is strongly continuous when $\lim _{\beta \rightarrow \beta_{0}} U(\beta)=U\left(\beta_{0}\right)$ for all $\beta_{0} \in \mathbb{R}$.

Theorem 9.4.2 (Stone (|Stone, 1932|)) There is a one-to-one correspondence between strongly continuous one-parameter unitary groups $U: n \rightarrow n$ in $\mathcal{M}(\beta)$ and self-adjoint matrices $H: n \rightarrow$ $n$ in $\mathcal{M}$. The bijection is given explicitly by $U(\beta)=e^{i \beta H}$ and $H=-i\left(\partial_{M} U\right)(0)$.

We use the bijection from the Stone's theorem to find the diagram $h \in Z X_{\mathbb{R}}$ such that $\llbracket h \rrbracket=H$. Using the property $U(0)=\mathrm{I}_{n}$ we obtain:

$$
\begin{align*}
H & =-i\left[\partial_{M} U(\beta)\right](0)=-i \otimes \llbracket\left[\partial_{\mathrm{ZX}} D_{U}\right](\beta) \rrbracket(0)=-i \otimes \llbracket\left[\partial_{\mathrm{ZX}} D_{U}\right](0) \rrbracket \\
& =\llbracket \overbrace{-\frac{\pi}{2}}^{\pi} \bigcirc\left[\partial_{\mathrm{ZX}} D_{U}\right](0) \rrbracket=\llbracket h \rrbracket \tag{9.102}
\end{align*}
$$

where the third equality is due to the fact that the evaluation commutes with the standard interpretation.

We give an example of diagram for an Ising Hamiltonian obtained via our approach.
Example 9.4.3 Let $H: 2 \rightarrow 2, H=Z_{1}-Z_{2}+Z_{1} Z_{2}$. The diagram $D_{U}(\beta)$ for $U(\beta)=e^{i \beta H}$ is:

$$
\begin{equation*}
D_{U}(\beta)=\overbrace{\pi}^{\beta}-(\beta-\beta--2 \beta)=e_{0}^{\beta} \tag{9.103}
\end{equation*}
$$

Using the formula 9.77 we find the derivative of $D_{U(\beta)}$ :

$$
\begin{equation*}
\partial_{\mathrm{ZX}} D_{U(\beta)}=\underbrace{\pi}_{-\frac{\pi}{2}} \tag{9.104}
\end{equation*}
$$

## Conclusion

In this thesis we explored the potential of quantum algorithms for combinatorial optimization on realistic industrial applications. On the theoretical side, we have developed a set of tools that allows to reason about variational algorithms in ZX-calculus.

We have considered two usecases issued from the field of smart charging. For both usecases our presentation contains i) a detailed complexity analysis of the natural formulations ii) a modelization in a quantum-amenable formulation iii) a description of approximation properties of the resulting models. We provide experimental protocols for QAOA and RQAOA for this applications. In particular, we specify selected QUBO encodings and parameter optimization routines. Finally, we numerically evaluate the performance of our routines. We compare our evaluations to the performance of classical algorithms. We carefully justify the choice of classical counterparts based on approximability properties of the considered problems.

We believe that our presentation positively contributes in the establishment of a general methodology for analysis of quantum algorithms for optimization problems. Completed with a mapping to hardware specifications, it constitutes a fairly exhaustive protocol for the examination of quantum heuristics applied to a considered practical problem.

Quantum algorithms for combinatorial optimization that are analyzed in this work are heuristics. Therefore, their practical performance in general heavily depends on the chosen application, or, more precisely, on considered instances of the optimization problem. For this reason it is extremely important to select right usecases. After working on this thesis, we may claim that a well-suited usecase for quantum algorithms should be at the same moment simple enough to be easily modeled as unconstrained problem over binary variables but hard enough to be difficult for classical algorithms. The model simplicity is crucial to avoid the overconsumption of limited resources available on NISQ devices. The classical intractability is necessary to leave the room for quantum advantage.

For an optimization problem the classical intractability may be interpreted in different ways. In this work we highlighted the importance of considering approximation algorithms and heuristics rather than exponential exact procedures for a fair evaluation of the difficulty of NP-hard optimization problems. In addition, we have observed that the instances issued from practical applications may differ in approximation complexity with the general model selected to represent them.

From the algorithmic side we have detected the importance of the fine-tuning of the experimental protocol. For instance, we have confirmed the observation stating that the genuine bottleneck of QAOA is the search of optimal parameters. This problem can be addressed by exploring the properties of the parameter landscape. In our numerical experiments, we have tested several existing approaches for the parameter optimization and confirmed the intuition behind them.

Our numerical results are reported in the sections 4.4 and 5.3 .3 proved the interest of the
application of quantum algorithms for both considered usecases. We have observed that QAOA and RQAOA have the performance that is comparable with pertinent classical competitors. We remark that due to the lack of resources our numerical analysis was restricted to versions of QAOA with low depth. We believe that the increase of depth may improve the performance of QAOA making it a powerful heuristic for practical applications.

In order to extend the range of applications that can be addressed by quantum algorithms we considered hybrid quantum-classical approaches. We believe that hybridization is essential to leverage the power of quantum computers for industrial applications. In particular, we suggested to outsource a difficult part of computation in the classical Branch \& Price to a quantum routine. The resulting method, called Quantum-assisted Branch \& Price, can be used to address integer linear programs with many variables. We have observed that quantum routines can enhance the runtime performance of the original classical scheme but the actual improvement is not the same across different families of considered instances (see section 6.4.3). We expect that on a real hardware the performance dependence on the instance properties (in particular on the connectivity) will be even more significant Harrigan et al., 2021, Pagano et al., 2020.

A further research in this direction requires more computational resources, in particular a real hardware of plausible size.

## Future work

Our work points out several promising directions for the future research.
In the close-term perspective we can enhance our experimental protocols by improving their performance and reducing the runtime. In particular, we can leverage the parameter concentration phenomena Brandao et al., 2018 and integrate learning-based techniques in the parameter optimization routine. We can start by considering the approaches presented in Khairy et al., 2020, Moussa et al., 2022.

In the long-term we are planning to validate our results on the quantum hardware and integrate hardware specifications in the design of experimental protocols.

Indeed, a fair evaluation of the performance of quantum algorithms requires hardware experimentations. In this direction we are planning to evaluate the impact of noise and the limited connectivity on the performance of quantum algorithms. In addition, we intend to explore the error mitigation techniques to handle the imperfections of the NISQ hardware.

On the application side, it is particularly compelling to qualify our hybrid Quantumassisted Branch \& Price algorithm on other practical optimization problems. For instance, in our future work we intend to apply this approach to the Frequency Containment Reserve problem - an industrial problem issued from the field of energy management where pricing subproblems take the form of Shortest Path to be found in graphs representing the charge/discharge of electrical vehicles connected to the grid.

Finally, we aim to continue the theoretical exploration of variational algorithms by using diagrammatic reasoning. For this purpose we lean on our tools for addition and differentiation of ZX-diagrams. For instance, we can use diagrammatic representation to analyze such properties of the parameter landscape as symmetries and barren plateau. We believe that with some extra work we can improve the protocol design for variational variational algorithms, for example by suggesting new parameter-shift rules. In addition, we plan to explore how such fundamental properties as $Q A O A$ supremacy can be understood in the diagrammatic terms.

## Appendix A

## Linear programming

Linear program is an optimization problem over continuous variables such that the objective function and the constraints are linear on input variables. While not strictly belonging to the field of combinatorial optimization, linear programs are ubiquitous in subroutines of exact and approximate algorithms.

Formally, a linear program in a standard form is:

$$
\begin{array}{ll}
\min & \sum_{i=1}^{n} c_{i} x_{i} \\
\text { s.t. } \sum_{i=1}^{n} a_{j, i} x_{i} \geq b_{j}, & j \in[1, \ldots, m] \\
x_{i} \geq 0, & i \in[1, \ldots, m] \tag{A.3}
\end{array}
$$

We remark that any linear program can be reformulated in a standard form.
Although relatively simple to formulate, linear program is an extremely powerful model for the operational research. For instance, it incorporates many resource allocation problems. A seminal example of a linear program is the diet problem where on have to built a meal of minimal cost satisfying nutritional requirements.

From the mathematical point of view a bounded linear program is precisely the optimization over a convex polyhedron defined by the set of constraints. Crucially, if the instance is feasible and bounded there exists at least one extreme point of the feasible region that corresponds to the optimal solution. We remark that in general case the optimal solution isn't unique.

## A. 1 Duality

For each linear program there exists another linear program called dual that results from an attempt to provide a lower bound on the minimum value of the initial problem. In the context of duality the initial linear program is usually referred as primal problem or primal formulation.

To formulate the dual problem we introduce one dual variable $\lambda_{j}$ per inequality constraint A.2 from the primal formulation. We want to derive bounds on the variables $\lambda_{j}$ such that for each feasible solution of the primal problem we have:

$$
\begin{equation*}
\sum_{j=1}^{m} b_{j} \lambda_{j} \leq \sum_{i=1}^{n} c_{i} x_{i} \tag{A.4}
\end{equation*}
$$

Firstly, we impose the non-negativity on the variables $\lambda_{j}$. The non-negativity of $\lambda_{j}$ allows to state that:

$$
\begin{equation*}
\lambda_{j} \sum_{i=1}^{n} a_{j, i} x_{i} \geq b_{j} \lambda_{j} \tag{A.5}
\end{equation*}
$$

Taking the sum over all constraints we obtain the condition:

$$
\begin{equation*}
\sum_{j=1}^{m} \lambda_{j} \sum_{i=1}^{n} a_{j, i} x_{i} \geq \sum_{j=1}^{m} b_{j} \lambda_{j} \tag{A.6}
\end{equation*}
$$

or, equivalently,

$$
\begin{equation*}
\sum_{i=1}^{n}\left(\sum_{j=1}^{m} a_{i, j} \lambda_{j}\right) x_{i} \geq \sum_{j=1}^{m} b_{j} \lambda_{j} \tag{A.7}
\end{equation*}
$$

We remark that for non-negative primal variables $x_{i} \geq 0$ the condition A.4 is satisfied if for each variable $x_{i} \geq 0$ the respective coefficients on the left sides of A.7) are smaller or equal to the coefficients on the right side, i.e.

$$
\begin{equation*}
\sum_{j=1}^{m} a_{i, j} \lambda_{j} \leq c_{i} \quad \Longrightarrow \quad \sum_{j=1}^{m} b_{j} \lambda_{j} \leq \sum_{i=1}^{n}\left(\sum_{j=1}^{m} a_{i, j} \lambda_{j}\right) x_{i} \leq \sum_{i=1}^{n} c_{i} x_{i} \tag{A.8}
\end{equation*}
$$

Therefore, if dual variables $\lambda_{j}$ satisfy the constraints

$$
\begin{array}{lr}
\sum_{j=1}^{m} a_{i, j} \lambda_{j} \leq c_{i}, & \forall i \in[1, \ldots, n] \\
\lambda_{j} \geq 0, & \forall j \in[1, \ldots, m]
\end{array}
$$

then $\sum_{j=1}^{m} b_{j} \lambda_{j}$ provides a lower bound for the objective value of the primal problem: $\sum_{i=1}^{n} c_{i} x_{i}$.
Roughly speaking, the dual problem consists in finding the best lower bound derivable in such way. Formally, we find the optimal bound by solving the following linear program:

$$
\begin{array}{ll}
\max & \sum_{j=1}^{m} b_{j} \lambda_{j} \\
\text { s.t. } & \sum_{j=1}^{m} a_{i, j} \lambda_{j} \leq c_{i},
\end{array} \quad \forall i \in[1, \ldots, n] \begin{array}{ll} 
& \\
\lambda_{j} \geq 0, & \forall j \in[1, \ldots, m] \tag{A.13}
\end{array}
$$

There is one dual variable for each constraint in the primal problem and every variable in the primal problem generates a constraint in the dual formulation. By construction, the optimal value of the dual formulation provides a lower bound on the optimum for the initial problem.

In fact, this latter observation about the respective optimal values of dual and primal formulations can be further strengthened with the duality theorem Gass and Harris, 2001. The duality theorem states in particular that if both problems are feasible and bounded then they have the same optimal value, i.e.

$$
\begin{equation*}
\sum_{j=1}^{m} b_{j} \lambda_{j}^{*}=\sum_{i=1}^{n} c_{i} x_{i}^{*} \tag{A.14}
\end{equation*}
$$

## A.1.1 Complementary slackness

Another important relation of dual and primal formulations is resumed in the theorem of complimentary slackness. Roughly speaking, the theorem states that in an optimal solution $\left(\boldsymbol{x}^{*}, \boldsymbol{\lambda}^{*}\right)$ either $x_{i}^{*}=0$ or the corresponding constraint in the dual formulation is binding (and vice-versa).

This theorem is a consequence of the duality theorem. Firstly, we recall that by definition the optimal values of both formulations satisfy the inequality:

$$
\begin{equation*}
\sum_{j=1}^{m} b_{j} \lambda_{j}^{*} \leq \sum_{i=1}^{n}\left(\sum_{j=1}^{m} a_{i, j} \lambda_{j}^{*}\right) x_{i}^{*} \leq \sum_{i=1}^{n} c_{i} x_{i}^{*} \tag{A.15}
\end{equation*}
$$

Actually, for feasible and bounded problems it follows from the duality theorem that:

$$
\begin{equation*}
\sum_{i=1}^{n}\left(\sum_{j=1}^{m} a_{i, j} \lambda_{j}^{*}\right) x_{i}^{*}=\sum_{i=1}^{n} c_{i} x_{i}^{*} \tag{A.16}
\end{equation*}
$$

or, equivalently,

$$
\begin{equation*}
\sum_{i=1}^{n}\left(\sum_{j=1}^{m} a_{i, j} \lambda_{j}^{*}-c_{i}\right) x_{i}^{*}=0 \tag{A.17}
\end{equation*}
$$

According to the definition a feasible solution for the primal satisfies $x_{i}^{*} \geq 0$ while for the feasible solution of the dual we have $\sum_{j=1}^{m} a_{i, j} \lambda_{j}^{*} \leq c_{i}$. Therefore, an optimal pair $\boldsymbol{x}^{*}, \boldsymbol{\lambda}^{*}$ is feasible if

$$
\begin{equation*}
\left(\sum_{j=1}^{m} a_{i, j} \lambda_{j}^{*}-c_{i}\right) x_{i}^{*} \leq 0 \tag{A.18}
\end{equation*}
$$

The theorem can be proved by contradiction. Indeed, if the complementary slackness doesn't hold, then from A.17 and A.18 follows that there exists at least one $v \in[1, \ldots, n]$ such that $\left(\sum_{j=1}^{m} a_{v, j} \lambda_{j}^{*}-c_{v}\right) x_{v}<0$. As the sum A.17 is strictly equal to zero, there should be at least one $u \in[1, \ldots, n]$ such that $\left(\sum_{j=1}^{m} a_{u, j} \lambda_{j}^{*}-c_{u}\right) x_{u}>0$. Such value $x_{u}$ violates the feasibility condition A.18).

An equivalent result for the dual solution can be proven following the same reasoning.

## A. 2 Simplex method

The simplex method was originally suggested in 1951 by Dantzig. A detailed description can be found in Nash, 2000. An introduction and a comprehensive review of important theoretical and practical aspect of the simplex method can be found in Vanderbei, 2014.

In a nutshell, the simplex algorithm iteratively considers extreme points of the polyhedron of feasible solutions. At each step it moves to an improving point. Usually, when the assignment $x_{1}=\cdots=x_{n} \equiv 0$ is feasible it is taken as a staring point. In the rest of the section we will consider the simplest case where the instance is bounded and a feasible initial solution is known.

More precisely, we begin by introducing slack variables $s_{1}, \ldots, s_{m}$ that transform inequality constraints A.2 to equalities. For the simplex method it actually doesn't matter is the variable
is present in the initial formulation or if it was introduced to capture the slack differences. Therefore, we consider the initial variables and the slack variables together and we denote them as $x_{1}, \ldots, x_{n+m}$.

With the slack variables the problem is reformulated as follows:

$$
\begin{array}{lr}
\min & \chi\left(x_{1}, \ldots, x_{n+m}\right)=\sum_{i=1}^{n} c_{i} x_{i} \\
\text { s.t. } & \sum_{i=1}^{n} a_{j, i} x_{i}-x_{n+j}=b_{j},
\end{array} \quad j \in[1, \ldots, m]
$$

We remark that by fixing values of any $n$ variables we can derive a unique feasible solution from the equality constraints A.20). This is the core idea of the simplex algorithm. In a nutshell, at each iteration the simplex methods chose a subset $\mathcal{B}$ of $n$ variables that will be set to zero. The set $\mathcal{B}$ is cased basis. We denote by $\mathcal{N}$ the set of $m$ non-basis variables. For fixed values of basis variables the values of variables in $\mathcal{N}$ are computed from the equality constraints:

$$
\begin{equation*}
x_{i}=\alpha_{i}+\sum_{j \in \mathcal{B}} \omega_{j, e} x_{j} \tag{A.22}
\end{equation*}
$$

Likewise, by replacing the variables $x_{i}, i \in \mathcal{N}$ with the expressions A.22 the objective function can be written as a linear combination of basis variables:

$$
\begin{equation*}
\chi_{\mathcal{B}}(x)=\chi_{0}+\sum_{i \in \mathcal{B}} \chi_{i} x_{i} \tag{A.23}
\end{equation*}
$$

The simplex method searches the set $\mathcal{B}$ for an improving variable. A variable $x_{i} \in \mathcal{B}$ is improving if we can diminish the value of the objective function by increasing the value of $x_{i}$ from zero. Clearly, a variable is improving if $\chi_{i}<0$.

In bounded problems the linear relationships A.20 determine an upper bound on the potential increase of every improving variable that doesn't violate the constraints $x_{i} \geq 0, \forall i \in \mathcal{N}$. The strongest bound correspond to so-called entering variable $x_{e} \in \mathcal{N}$. In other words, for an improving variable $x_{i}$ the entering variable is:

$$
\begin{equation*}
x_{e}=\alpha_{e}-\omega_{i, e} x_{i}+\sum_{\substack{j \in \mathcal{B} \\ j \neq i}} \omega_{j, e} x_{j} \tag{A.24}
\end{equation*}
$$

such that the ratio $u_{i, e}=\alpha_{e} / \omega_{i, e}$ is the smallest one for $x_{e} \in \mathcal{N}$. If we increase $x_{i}$ for more than $u_{i, e}$ units the constraint $x_{e} \geq 0$ becomes violated.

We remark that there may be several improving variables. Usually, we select the one that leads to the most important decrease in the value of the objective function. Precisely, we chose the improving variable $i$ with the value $\chi_{i} u_{i, e}<0$ of largest magnitude.

Once the method has identified one improving and one entering variable it performs a pivot. During the pivot procedure the improving variable $x_{i}$ leaves the basis replaced by the entering variable $x_{e}$. In other words, we rewrite the objective function A.23) and the constraints A.22) by replacing $x_{i}$ with

$$
\begin{equation*}
x_{i}=\frac{1}{\omega_{i, e}}\left(\alpha_{e}-x_{e}+\sum_{\substack{j \in \mathcal{B} \\ j \neq i}} \omega_{j, e} x_{j}\right) \tag{A.25}
\end{equation*}
$$

The new solution is computed by fixing $x_{e}=0$. The variable $x_{i}$ gets the value $u_{i, e}$ and the value of the objective function id decreased by $\left|\chi_{i} u_{i, e}\right|$. The procedure continues until there is no more improving variables.

## A.2.1 Graphical interpretation

The simplex procedure has a particularly nice graphical interpretation.
Firstly, we remark that setting $x_{j}=0$ for a slack variable $(j \in[n+1, \ldots, m])$ implies that the constraint $j$ in the original formulation is binding, i.e. the solution lays exactly on the polyhedron edge defined by $\sum_{i=1}^{n} a_{j, i}=b_{j}$. Similarly, if $x_{i}=0$ for a variable from the initial formulation the solution belongs to the edge defined by the constraint $x_{i} \geq 0$. Therefore, the basis $\mathcal{B}$ defines an intersection of $n$ polyhedron edges. For a polyhedron in $n$-dimensional space such intersection is precisely one of the extreme points. The graphical representation of the simlpex method for the two-variable case is shown on the Figure A.1.


Figure A.1: At each iteration the simplex method moves to an extreme point corresponding to a better solution. Levels of the objective function $\chi\left(x_{1}, x_{2}\right)=x_{1}+x_{2}$ are denoted by dashed lines.

We remark that if we add the integrality requirement $\boldsymbol{x} i n \mathbb{Z}^{n}$ on the variables (as in Integer Linear Programs (ILP) studied in the Section 6.1) the optimum solutions are in general not in the extreme point of the polyhedron. Indeed, if an extreme points has fractional coordinates it is not even a feasible solution for such formulations. This is the major reason why the optimization of linear programs with integer variables is in general much harder then the optimization over real values.

For bounded problems the polyhedron has $\binom{n+m}{n}$ extreme points. Therefore, in the worst case the simplex method has an exponential runtime. However, such situation rarely happens in practice and the simplex remains one of the most efficient methods for the linear programs. Alternatively, linear program can be solved in guaranteed polynomial runtime with so-called interior-point methods Karmarkar, 1984, Khachiyan, 1980.

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#### Abstract

The domain of energy management involves many combinatorial optimization problems known to be computationally hard. The emergence of quantum computers suggests new approaches for these problems. For near-future machines particularly promising are variational quantum heuristics such as QAOA that can leverage the computational power of the imperfect quantum hardware.

We explore the potential of variational quantum algorithms for optimization problems issued from the field of "smart charging" of electrical vehicles. We consider two problems inspired by real-world usecases. In the first problem, modeled as Max-K-Cut, we search to schedule a set of prioritized charges on several stations while minimizing the weighted completion time. In the second problem, modeled as Maximum Independent Set, we aim to maximize the number of satisfied charge demands on a single station while respecting the conflicts between demands. For both problems we develop an experimental protocol specifying the encoding step and the parameter optimization routine. Our numerical experiments confirm the interest of quantum heuristics for these problems as well as the quality of our experimental protocol.

In order to extend the applicability of quantum heuristics we introduce a new hybrid approach that integrates quantum routines in the classical Branch \& Price algorithm for large integer linear programs. We test this approach on a smart charging problem that is modeled as graph coloring problem. Our computational results affirm the potential of the hybrid approach while revealing the considerable dependence of the performance gain on the particular instance of the problem.

Important components of variational algorithms can be represented as ZX-diagrams. We demonstrate how the rewriting rules of ZX-calculus can be used to derive the analytical formula for the mean energy of a general Ising model in a $\mathrm{QAOA}_{1}$ state. Furthermore, we contribute to the theoretical exploration of variational algorithms by extending the ZX-calculus with addition and differentiation of ZX-diagrams. Our inductive procedure for the addition is fully diagrammatic. For the differentiation we suggest two approaches. The first approach is inductive, it leverages our procedure for addition to explicitly represent the product rules. The second approach is resumed in two formulas that are derived from the factored form of parameterized diagrams.


Keywords: Combinatorial optimization, quantum computing, energy management, smart charging, ZX-calculus

## Résumé

Le domaine du management de l'énergie implique de nombreux problèmes d'optimisation combinatoire connus pour être difficiles. L'émergence des ordinateurs quantiques suggère de nouvelles approches pour ces problèmes. Pour les machines du futur proche, les heuristiques quantiques variationnelles telles que QAOA, qui peuvent tirer parti de la puissance de calcul des ordinateurs quantiques imparfaits, sont particulièrement prometteuses.

Nous explorons le potentiel des algorithmes quantiques variationnels pour des problèmes d'optimisation issus du domaine de "smart charging" des véhicules électriques. Nous considérons deux problèmes inspirés de cas d'utilisation réels. Dans le premier problème, modélisé par Max-$K$-Cut, nous cherchons à planifier un ensemble de charges avec priorités sur plusieurs stations tout en minimisant le temps de completion pondéré. Dans le deuxième problème, modélisé par Maximum Independent Set, nous cherchons à maximiser le nombre de demandes de charge satisfaites sur une seule station tout en respectant les conflits entre les demandes. Pour les deux problèmes, nous développons un protocole expérimental spécifiant l'encodage et la routine d'optimisation des paramètres. Nos expériences numériques confirment l'intérêt des heuristiques quantiques pour ces problèmes ainsi que la qualité de notre protocole expérimental.

Afin d'étendre l'applicabilité des heuristiques quantiques nous introduisons une nouvelle approche hybride qui intègre des routines quantiques dans l'algorithme classique de Branch \& Price pour les programmes linéaires en nombres entiers de grande taille. Nous testons cette approche sur un problème de smart charging qui se modélise comme un problème de coloration de graphe. Nos résultats numériques affirment le potentiel de l'approche hybride tout en révélant la variabilité du gain de performance en fonction de l'instance particulière du problème.

Les parties importants des algorithmes variationnels peuvent être représentées sous forme de diagrammes ZX. Nous démontrons comment les règles de réécriture du ZX-calculus peuvent être utilisées pour dériver la formule analytique de l'énergie moyenne d'un modèle d'Ising dans un état QAOA de profondeur 1. De plus, nous contribuons à l'exploration théorique des algorithmes variationnels en étendant le ZX-calcul avec addition et différenciation de ZX-diagrammes. Notre procédure inductive pour l'addition est entièrement graphique. Pour la différenciation, nous suggérons deux approches. La première approche est inductive, elle s'appuie sur notre procédure d'addition pour représenter explicitement les règles du produit. La deuxième approche est résumée dans deux formules qui sont dérivées de la forme factorisée des diagrammes paramétrés.

Mots-clés: Optimisation combinatoire, algorithmes quantiques, management de l'energie, smart charging, ZX-calculus

## A. 3 Résumé étendu (en français)

Le domaine de la gestion de l'énergie traite de nombreux problèmes d'optimisation qui sont difficiles. Trouver de meilleures solutions à ces problèmes peut réduire considérablement les coûts, améliorer la qualité du service ou résoudre des problèmes opérationnels émergents tels que l'intégration de sources d'énergie renouvelables. Cette thèse est dédiée à l'application d'algorithmes quantiques à des problèmes d'optimisation combinatoire issus du domaine de la "recharge intelligente" - un domaine qui couvre les problèmes liés à la recharge des véhicules électriques.

De nombreux problèmes de charge intelligents sont modélisés comme des problèmes d'optimisation combinatoire NP-difficiles. Cela implique que (dans le pire des cas) la recherche de la solution optimale prend un temps exponentiel. Par conséquent, en pratique de tels problèmes sont résolus de façon approché. Une solution approchée peut être calculée par deux types d'algorithmes : ceux avec une qualité d'approximation théoriquement démontrée (appelés algorithmes d'approximation) et ceux qui performent bien en pratique (appelés heuristiques).

Pour ces problèmes is est probable que l'utilisation des algorithmes quantiques donne un avantage soit en termes du temps d'écécution, soit en thèrmes de la qualité de la solution obtenue.

L'idée d'utiliser des effets quantiques pour améliorer les calculs des propriétés des systèmes physiques (et les calculs en général) a été suggérée indépendamment par Richard Feynman Feynman, 1986, Yuri Manin Manin, 1980] et Paul Benioff Benioff, 1982. . Des machines qui exploitent les principes de la mécanique quantique comme l'intrication et la superposition sont manifestement plus puissantes que les machines de Turing classiques. Pour le moment, les deux résultats les plus impressionnants sont l'algorithme de Grover pour une recherche dans une base de données non structurée Grover, 1996 et l'algorithme de factorisation des entiers découvert par Shor Shor, 1995.

D'un point de vue très général la programmation d'un ordinateur quantique consiste à combiner petites briques pour préparer une distribution de probabilitéqui est pertinent pour le problème en question. Il n'est pas évident de concevoir des algorithmes dans un tel framework. En partie à cause de cela, le domaine de l'informatique quantique se concentre beaucoup sur les algorithmes variationnels qui ne proposent pas des programmes fixes mais plutôt des modèles entraînables qui sont optimisés pour trouver de bonnes solutions Preskill, 2018.

Les algorithmes variationnels sont bien adaptés au ainsi applélé NISQ-hardware (Noisy Intermediate Scale Quantum) qui propose entre $10^{2}$ et $10^{3}$ qubits. En raison de la taille limitée et du niveau de bruit élevé, nous ne nous attendons pas à ce que les ordinateurs NISQ exécutent les célèbres algorithmes de Grover et Shor sur des problèmes de tailles pratiquement intéressantes Babbush et al., 2021. En plus, il est probable que l'accéleration polynomiale des algorithmes Montanaro, 2020, Ambainis, 2005, Cerf et al., 2000, van Apeldoorn et al., 2020 sera anéanti par les shémas de correction du bruit Babbush et al., 2021.

Nous suggerons que le meilleur chemin pour avoir un avantage avec les machines NISQ est de combiner less algorithmes variationels avec les routines classiques dans les procédures hybrides.

L'AA initialise le système dans un état fondamental facile à préparer de certains hamiltoniens H initialiser, qui évolue sous l'hamiltonien dépendant du temps qui passe progressivement deHinitialiser auHfinaloù l'état fondamental deHfinalcorrespond à la solution du problème d'optimisation.

Algorithmes quantiques pour l'optimisation combinatoire: Dans ce travail, nous considérons le soi-disant algorithme adiabatique (qui renvoie une solution exacte pour un problème
d'optimisation) et ses proches cousins - le recuit quantique et QAOA donné. L'algorithme adiabatique The initialise le système dans un état fondamental d'un certains Hamiltonien $H_{\text {init }}$ facile à préparer (appelés Hamiltonian de mélange), laisse évoluer le système sous l'Hamiltonien qui passe progressivement de $H_{\text {init }}$ à $H_{\text {final }}$ où l'état fondamental de $H_{\text {final }}$ correspond à la solution du problème d'optimisation.

Le théorème adiabatique postule que si l'évolution est suffisamment lente (par rapport à la différence entre l'énergie de l'état fondamental et l'énergie du premier état excité habituellement appelé écart spectral), le système reste proche de l'état fondamental instantané pendant tout le processus Jansen et al., 2007, Amin, 2009a. Par conséquent, avec une forte probabilité, le résultat de mesure de l'état final sera l'état fondamental de $H_{\text {final }}$ - précisément la solution exacte du problème initial.

Comme la condition d'être "suffisamment lent" est qénéralement difficile à satisfaire en pratique (ou même à évaluer), l'Algorithme Adiabatique reste essentiellement une construction théorique. Il a cependant inspiré plusieurs heuristiques polynomial telles que le recuit quantique et le QAOA Farhi et al., 2014a

Le recuit quantique modifie l'hamiltonien de la même manière que l'algorithme adiabatique mais beaucoup plus rapidement que ne le permet la condition adiabatique. L'évolution est effectuée sur les machines spécialles appelées quantum annealers, et le temps d'exécution est généralement limité par les caractéristiques de la machine. Le recuit quantique a été appliqué à de nombreux problèmes d'optimisation tels que le voyageur de commerce et la coloration de graphes Rieffel et al., 2014, reconnaissance d'images Venturelli et al., 2015 et la recherche de l'état fondamental du modèle de Sherington-Kirkpatrik [Venturelliet coll., 2015]. Essentiellement, l'idée derrière le recuit quantique consiste à utiliser les fluctuations quantiques afin d'explorer l'espace des solutions de façon efficace.

Contrairement au recuit quantique qui nécessite une machine analogique spécifique, le QAOA adapte les idées de l'algorithme adiabatique aux ordinateurs quantiques universels basés sur les circuits [Farhi et al., 2014a. La dynamique QAOA n'est pas régie par une interpolation continue enrte deux hamiltoniens, mais par une séquence de pulsions ("bangs") qui saute entre $H_{\text {initial }}$ et $H_{\text {final }}$.Le nombre total des pulsions $p \in \mathbb{N}$, appelé profondeur, est prédéfini par l'utilisateur tandis que les durées de bangs sont optimisées directement à l'intérieur de l'algorithme par une routine classique. C'est pour cela que l'algorithme est dit variationnel. Il est connu que pour le grand nombre de couches $p \rightarrow$ inf l'algorithme atteint l'optimum exact. De plus, contrairement au recuit quantique Crosson et al., 2014, la croissance de la qualité de la solution en temps est monotone.

Récemment, il y a eu un progrès remarquable dans la mise en œuvre experimental des quantum annealers : 5000 les qubits sont accessibles sur DWave Advantage system 6.1 Systems, 2022 ] par rapport à des 127 qubits sur l'ordinateur quantique universelle IBM Eagle Dial, 2022. Par conséquent, le recuit quantique se prête mieux à l'analyse expérimentale alors que en raison du manque d'outils, l'analyse théorique de l'algorithme est nettement plus difficile Braida et al., 2022, Moosavian et al., 2022.

D'autre part, pour QAOA de nombreux résultats théoriques ont été obtenus. Entre ces résultats, les plus importants sont la suprématie [Farhi and Harrow, 2016] et l'universalité Lloyd, 2018]. Certaines garanties de performances ont été démontrées pour des instances spéciales de MaxCut [Farhi et al., 2014a, Hadfield, 2018, Wurtz and Love, 2021, Marwaha, 2021, Basso et al., 2022] et E3LIN2 Farhi et al., 2014b ainson que les résultats limitants Bravyi et al., 2020, Farhi et al., 2020, Chou et al., 2021, bien qu'aucun avantage quantique n'ait été prouvé.

Le dernier algorithme considéré est le QAOA récursif (RQAOA). Le RQAOA a été initialement conçu pour traiter le problème delocalité qui a souvent été mentionné comme la condi-
tion limitante pour la performance QAOA Bravyi et al., 2020, Farhi et al., 2020, Hastings, 2019, Chou et al., 2021. RQAOA procède par une élimination récursive de variables, où à chaque étape la variable à éliminer est sélectionnée en utilisant l'état optimal de QAOA pour le problème restreint. La promesse derrière est que, contrairement au QAOA à profondeur constante, le RQAOA finit par voir la structure globale du problème. Du point de vue expérimental, un fait particulièrement intéressant à propos du RQAOA est que pour $p=1$ il peut être efficacement simulé sur un ordinateur classique Bravyi et al., 2020, Egger et al., 2021.


Figure A.2: Schema d'interaction entre la partie classique et la partie quantique dans QAOA

La formulation QUBO: I Afin d'appliquer des algorithmes quantiques à des problèmes d'optimisation combinatoire, nous devons d'abord associer l'espace des solutions $S$ à l'ensemble $|x\rangle, x \in\{0,1\}^{n}$ des états de base d'une système à $n$ qubits. Ensuite, nous devons définir un Hamiltonien diagonal $H$, telle que pour chaque solution réalisable $x \in\{0,1\}^{n}$, la valeur propre de $H$ corresponde à la valeur de la fonction à optimiser:

$$
\begin{equation*}
H|x\rangle=f(x)|x\rangle, \quad|x\rangle \in S_{b} \tag{A.26}
\end{equation*}
$$

Ce mapping est évident pour les problems quadratiques sur les variables binaires et sans contraintes (applelés QUBO) et relativement facile pour beaucoup de problèmes d'optimisation NP-difficiles Lucas, 2014, Alidaee et al., 1994, Kochenberger et al., 2013, Alidaee et al., 2008].

Lorsqu'une formulation du problème d'optimisation dépasse le cadre QUBO (elle a les contraintes ou les variables non-binaires), un traitement spécial est nécessaire pour capturer dans un modèle QUBO ces éléments. En informatique quantique, il existe deux approches pour intédrer les contraintes. Le premier est la pénalisation, soit la modification de la fonction objectif d'une manière qui rend les solutions irréalisables sous-optimales. Pour une contrainte $x \in S$ la pénalité est une fonction $g(x):\{0,1\}^{n} \rightarrow \mathbb{R}^{+}$tel que: $g(x)=\left\{\begin{array}{ll}0, & x \in S \\ \geq 1, & x \notin S\end{array}\right.$.

La deuxième technique pour gérer les contraintes consiste à remplacer l'Hamiltonian de mélange pour maintenir l'évolution dans le sous-espace réalisable Hen and Spedalieri, 2016, Hen and Sarandy, 2016, Hadfield et al., 2019.

Le recharge intelligente avec les algorithmes quantiques Nous avons sélectionné deux cas d'usage réels pour évaluer les approches quantiques. La première concerne l'ordonancement: le but est d'attribuer un ensemble de charges à différentes bornes de recharge. Dans le deuxième problème, il faut sélectionner un sous-ensemble de demandes de charge à satisfaire tout en respectant les conflit entre les demandes.

Dans le problème d'ordonnancement, on avons un ensemble de demandes (non-préemptif) de charge $J,|J|=n$. Chaque charge $j \in J$ à un temps d'éxécution fixe $t_{j} \in \mathbb{R}^{+}$et une priorité $w_{j} \in \mathbb{R}^{+}$. L'objectif est de distribuer les charges entre un ensemble $M$ de $m$ bornes de recharge identiques. Dans le cas trivial avec une borne un algorithme polynomial découle directement de la rè̀le de Smith [Smith, 1956]. Cependant, pour des instances avec les priorités positives générales le problème est NP-difficile au sens ordinaire quand $m \geq 2$ Skutella and Woeginger, 2000 et NP-difficile au sens fort quand $m$ fait partie de l'entrée du problème Garey and Johnson, 1990. Néamons, du point de vue de l'approximation, le problème est simple - pour un nombre fixe de stations $m$ le problème est pseudopolynomial et admet un FPTAS [Sahni, 1976, Woeginger, 2001].

Le problème peut être modélisé comme la recherche de Maximum-m-Cut dans un graph pondéré Skutella, 1998, et pour le cas ou $m=2$ cette formulation est facilement transformable dans la forme QUBO Alidaee et al., 1994. Le problème Maximum-2-Cut est NPdifficile, en plus il est NP-difficile à approximer la solution avec un facteur superieur à $0.941 \ldots$ Arora et al., 1998. Le meilleur algorithme d'approximation pour Maximum-2-Cut a le ratio $\alpha_{G W} \geq 0.87856$ Goemans and Williamson, 1995. De plus, sous la Unique Game Conjecture Khot, 2002 le ratio 0.87856 est optimal Khot et al., 2007.

Une formulation naturelle du problème Max-m-Cut utilise des variables entières $x_{i} \in[1, \ldots, m]$. correspondants aux couleurs attribuées aux nœuds:

$$
\begin{equation*}
\max _{x \in[1, \ldots, m]^{n}} \sum_{(u, v) \in E} e_{u, v} I\left(x_{u} \neq x_{v}\right) \tag{A.27}
\end{equation*}
$$

Pour mapper la formulation A.27à QUBO, nous avons suggéré d'utiliser un encodage binaire. Pour $m=2^{l}$ chaque variable entière $x_{i} \in[0, \ldots, m-1]$ est mappé à une séquence de $l=\log _{2} m$ variables binaires. Une chaîne de bits $\boldsymbol{b}_{\boldsymbol{i}}=\left[b_{0}, \ldots, b_{l-1}\right]$ correspond à l'entier $x_{i}=\sum_{0 \leq j \leq l-1} 2^{j} b_{j}$.

Nous avons analysé numériquement les performances de QAOA sur des graphes correspondant au problème de recharge intelligente. Les instances ont été générées à partir d'un ensemble de données réels contenant environ 2250 charges réalisées au cours du mois de mai 2017 sur des points de charge du réseau Belib's Bel, 2017. Nous avons utilisé le phénomène de concentration pour prédire de bons points d'initialisation pour différentes taille des instances $n$. Pour l'optimisation des paramètres, nous avons comparé différentes méthodes d'optimisation locales; il s'est avéré que le Nelder-Mead est le meilleur choix à la fois en termes de qualité de la solution et de temps d'exécution. Comme prévu, nous avons observé que la performance QAOA augmente avec la profondeur $p$. De manière surprenante, nous avons également remarqué que le rapport d'approximation s'améliore avec la taille $n$ de l'instance.

Dans le problème de sélection des demandes de charge à satisfaire, chaque demande (non préemptive) doit être éxécuté dans un interval prédéfini $\left[T_{j}^{0}, T_{j}^{0}+t_{j}^{0}\right]$. Si accepté, un demande $j$ apporte un profit $p_{j} \in \mathbb{R}^{+}$. Le décideur vise à sélectionner une planification valide avec le profit le plus élevé. Une planification est valide si les tâches sélectionnées sont simultanément satisfaisables, i.e. leurs intervals ne se chevauchent pas $\left[T_{i}^{0}, T_{i}^{0}+t_{i}\right] \cap\left[T_{j}^{0}, T_{j}^{0}+t_{j}\right] \neq \emptyset$. Nous considérons ici ce problème avec des contraintes de groupe en plus qui indiquent que de demande de même groupe de ne peuvent pas être satisfait ensembe.

Ce problème équivaut à la recherche d'un stable maximal (MIS) dans un graphe. Le MIS est difficile à approximer : selon Håstad, 1999, Zuckerman, 2007) pour toute $\epsilon$ il n'y a pas
d'algorithme polynomial retournant pour chaque instance un stable $S^{*}$ tel que $\left|S_{\text {max }}\right| /\left|S^{*}\right| \leq n^{1-\epsilon}$ sauf si $P=$ NP.

La formulation QUBO pour le problèmde de stable maximal est:

$$
\begin{align*}
& \max \sum_{v \in V} x_{v}-\lambda \sum_{(u, v) \in E} x_{u} x_{v}  \tag{A.28}\\
& \quad x \in\{0,1\}^{|V|} \tag{A.29}
\end{align*}
$$

where $\lambda=2$.
Afin d'évaluer les performances de l'heuristique quantique sur le problème MIS, nous avons créé un benchmark qui contient des graphes réguliers aléatoires avec des degrés dans $[3,5,7,9]$ et des tailles allant de 100 à 1000 noeuds. Nous avons également inclus des graphes du benchmark DIMACS et des instances réelles issues d'un simulateur de scénarios de demande de recharge développés en interne par EDF. Comme RQAOA en profondeur 1 peut être simulé efficacement, nous l'avons utilisé pour évaluer les performances sur de grandes instances.

Nous avons comparé le RQAOA à l'algorithme d'approximation Boppana, 1990 sur 121 graphes réguliers aléatoires avec jusqu'à 1000 sommets. Nous avons observé que sur tous sauf 7 instances RQAOA renvoie une solution au moins aussi bonne que celle trouvée par la routine classique. Sur le benchmark DIMACS, bien que nous n'ayons pas été en mesure de surpasser définitivement les heuristiques classiques, les solutions obtenues étaient encore compétitives dans de nombreux cas. De plus, nous pensons que les performances de l'algorithme peuvent être considérablement augmentées si l'on considère la profondeur plus élevée $p>1$. En revanche, nous avons observé que la qualité de la solution n'est pas robuste d'une instance à l'autre. Par conséquent, la sélection des bons cas d'utilisation qui conduisent à des instances bien adaptées à l'heuristique quantique reste un défi extrêmement important pour les algorithmes quantique du futur proche.

Approches hybrides: Selon les estimations actuelles, le matériel quantique du futur proche aura un nombre relativement faible de qubits avec une connectivité limitée [Preskill, 2018]. Par conséquent, pour la plupart des applications matérielles quantiques, il est nécessaire de combiner des routines classiques et quantiques dans des approches hybrides. Les routines classiques peuvent être utilisées en prétraitement Choi, 2008, en posttraitement ainsi qu'être directement intégrées dans de telsalgorithmes variationnels comme QAOA et VQE Braine et al., 2021] ou les utiliser comme sous-programme Yonaga et al., 2020, Ohzeki, 2020.

Outre la gestion des contraintes matérielles, l'hybridation permet de résoudre les formulations plus compliquées qu'un simple QUBO. Dans cette thèse notre objectif est de concevoir un approche hybride quantique-classique pour les programmes linéaires en nombre entiers sur un grand nombre des variables. Pour cela, nous combinons des routines quantiques avec la technique de Branche Price qui s'attaque à un grand nombre de variables avec la génération de colonnes Vanderbeck, 1994. Nous testons notre approche sur la formulation étendue du problème de coloration de graphe (problème maître) Mehrotra and Trick, 1996:

$$
\begin{array}{cc}
\text { (GC-s) : } \min & \sum_{\substack{s \in I(G)}} \lambda_{s} \\
& \sum_{\substack{s \in I(G) \\
v \in s}} \lambda_{s} \geq 1, \\
& \lambda_{s} \in\{0,1\},
\end{array} \forall v \in V
$$

où $I(G)$ est l'ensemble de tous les stables dans le graphe. Suivant l'approche suggérée dans Mehrotra and Trick, 1996, nous ne considérons que le sous-ensemble restreint $I^{\prime} \subsetneq I(G)$ dev variables et ensuite ajoutons des variables «prometteuses » avec la routine de pricing.

Dans Svensson et al., 2021 QAOA est utilisé comme solveur heuristique pour le problème maître restreint. Au lieu de cela, nous suggérons un approche différent qui intègre l'heuristique quantique dans l'étape de pricing de la génération des colonnes. Pour lae problème de coloration, le sous-problème de pricing équivaut à la recherche d'un stable maximal pondéré (MWIS). Dans notre approche hybride, l'instance de MWIS est abordée avec une procédure en trois volets qui combine l'heuristique classique de Held et al., 2012, RQAOA et l'algorithme exact.

Une expérience numérique sur des instances aléatoires de densités différentes montre qu'il est effectivement profitable dans certaines cas d'utiliser la méthode hybride au lieu de l'approche classique traditionnelle. Malheureusement, pour les instances issues du problème de sélection des tâches de charge, nous avons observé qu'il n'y a pas besoin de RQAOA car l'heuristique classique trouve toujours la solution du problème de pricing quand une telle solution existe.

ZX-calculus: ZX-calcul, initialement introduit dans Coecke and Duncan, 2011, est un langage graphique qui permet de raisonner sur l'informatique quantique. Dans ce langage, les calculs complexes sur les qubits sont représentés par diagrammes fabriqué à partir de générateurs élémentaires. Chaque diagramme correspond à une transformation linéaire entre des espaces de Hilbert d'états des qubits. Un ensemble compact de règles de réécriture permet de transformer des diagrammes en diagrammes équivalents. L'avantage notable du calcul ZX par rapport à d'autres représentations est que dans ce langage, les calculs peuvent être effectués entièrement avec les transformations graphiques Hadzihasanovic et al., 2018, Jeandel et al., 2018b.

Grâce à sa flexibilité, ZX-calculus est largement utilisé pour addresser différents problèmes d'informatique quantique. Par exemple, le calcul ZX a permis la dérivation de résultats importants dans le domaine de Measurement-based Quantum Computing (MBQC) Duncan and Perdrix, 2010 , Kissinger and van de Wetering, 2019 ainsi que pour l'optimisation des circuits Cowtan et al., 2020, de Beaudrap et al., 2020, Kissinger and van de Wetering, 2020b, Duncan et al., 2020 et pour la conception des codes correcteurs d'erreurs [Garvie etDuncan, 2018], ainsi que l'analyse et la compilation decodes de surface Garvie and Duncan, 2018, Horsman, 2011, de Beaudrap and Horsman, 2020, Hanks et al., 2020.

Cependant, les applications du calcul ZX auw algorithmes variationnels sont jusqu'à présent limités. Nous pensons que la raison pour laquelle ils sont encore inexplorés avec les moyens du calcul ZX est l'absence d'un moyen pratique de différencier les diagrammes paramétrés. En effet, les blocs de construction de base des algorithmes variationnels sont des circuits paramétrés, et la recherche de valeurs de paramètres optimales est une partie cruciale de ces algorithmes. En théorie, trouver des paramètres optimaux peut être NP-difficile Bittel and Kliesch, 2021. En pratique, la recherche est généralement effectuée par des méthodes d'optimisation numériques classiques et la plupart d'entre elles utilisent des dérivées Guerreschi and Smelyanskiy, 2017.

La principale difficulté de différenciation des ZX-diagrams proviennent du règle du produit: $\partial f g=\partial f g+f \partial g$ gqui consiste à ajouter deux termes. Cette règle est cruciale pour évaluer la dérivée de la composition séquentielle et du produit tensoriel. Comme les diagrammes ZX sont définis de manière inductive avec ces deux compositions, la dérivée d'un diagramme complexe pourrait être calculée à partir des dérivées de ses parties si nous étions autorisés à utiliser la règle du produit, et donc les diagrammes de somme.

Les œuvres Zhao and Gao, 2021, Toumi et al., 2021 utilisent des sommes explicites de diagrammes pour représenter la dérivée des diagrammes avec plusieurs occurrences du paramètre.

L'inconvénient majeur de cette approche est qu'il n'y a pas de règles pour manipuler des sommes formelles des ZX-diagrammes. Par conséquent, nous ne pouvons pas exploiter pleinement la puissance du calcul graphique. Dans cette thèse, nous proposons une approche oùla dérivée d'un paramètreZX-le schéma est un autre ZX-diagramme. On évite ainsi l'extension de la signature avec des sommes formelles. Afin d'aborder les sommes qui apparaissent dans la règle du produit, nous introduisons une technique originale pour effectuer l'addition des diagrammes entièrement dans laZX-calcul. Pour cela, on utilise des diagrammes spéciaux appelés états contrôlés Jeandel et al., 2019:

Definition A.3.1 (Etats contrôlé Jeandel et al., 2019]) Une diagramme ZX avec une entrée et $n$ sorties est un état contrôlé si $\llbracket C \rrbracket|0\rangle=\sum_{x \in\{0,1\}^{n}}|x\rangle=\llbracket \underbrace{\bigcirc \ldots}_{n} \downarrow \rrbracket$.

Nous proposons une voie inductive (nommée controlizer) pour représenter chaque ZX-diagramme par un tel état.

Definition A.3.2 (Controlizer) On dit qu'un mapping $C$ qui associe à chaque diagramme avec $n$ entrées et $m$ sorties une autre diagramme avec une entrée et $n+m$ sorties est un controlizer si pour toute diagramme ZX-diagram $D$ :

- $C(D)$ est un état contrôlé
- L'état correspond à la diagramme $D$ dans le sense que:

Comme on sait sommer des états contrôlés Jeandel et al., 2019 l'addition pour les diagrammes arbitraires suit. Malheureusement, notre démarche inductive conduit à de grands diagrammes, même lorsque les diagrammes initiaux sont assez simples. Ce défaut est cependant conforme à l'intuition. En effet, contrairement aux compositions séquentielles et parallèles, l'addition n'est pas une opération physique, il n'est donc pas surprenant qu'elle ne puisse pas être facilement intégrée dans le cadre des schémas ZX. Du côté positif, on peut s'appuyer sur la puissante théorie des équations du ZX-calcul pour simplifier, quand c'est possible, les diagrammes représentant la somme de diagrammes.

Une approche alternative pour l'addition dans le calcul ZXW algébrique a été présentée dans Yeung et al., . Il s'appuie également sur les formes contrôlées du diagramme, mais leurs formes contrôlées sont différentes des nôtres.

Une définition inductive de la dérivée est obtenue par une représentation explicite des règles du produit.

Dans un but de donner une technique prête à l'emploi pour la différenciation, nous proposons un moyen simple et pratique de calculer la dérivée pour la famille de diagrammes linéaires $\mathrm{ZX}(\beta)$ Jeandel et al., 2019, c'est-à-dire des diagrammes où les angles ne peuvent dépendre que linéairement (avec des coefficients entiers) d'un paramètre $\beta$. La plupart des circuits utilisés pour
les algorithmes variationnels appartiennent $\operatorname{ZX}(\beta)$ et nous pensons que nos formules rendront leur analyse beaucoup plus simple.

Tout d'abord, nous observons que chaque diagramme dans $\operatorname{ZX}(\beta)$ peut être transformé sous la forme

où $n, m$ sont les nombres entièrs et $D_{1}, D_{2}$ ne contiennent pas $\beta$.
Definition A.3.3 Etant donné une diagramme $D(\beta)=D_{2} \circ\left(D_{1} \otimes X_{\beta}(n, m)\right)$ dans la forme (A.34), on a

$$
\begin{align*}
& \partial_{\mathrm{ZX}}\left[X_{\beta}(n, m)\right]=\partial_{\mathrm{ZX}}[\underbrace{(\beta) \cdots \underbrace{\beta}}_{n} \underbrace{-\beta-\Theta^{-\beta}}_{m}] \tag{A.35}
\end{align*}
$$

Une définition d'une dérivée similaire à nos formules a été obtenue dans le travail indépendant Wang and Yeung, 2022. Ce travail utilise W-spiders pour gérer les règles du produit. Contrairement à notre résultat, la différenciation diagrammatique présentée dans Wang and Yeung, 2022 mappe un diagramme ZX à un diagramme d'un autre langage appelé ZX-calculus algébrique Wang, 2020. Le calcul ZX algébrique est pratique pour représenter des nombres complexes arbitraires, donc leur procédure de différenciation peut gérer des familles de diagrammes plus générales que $\operatorname{ZX}(\boldsymbol{\beta})$. Cet avantage s'accompagne du coût de l'abandon de l'héritage du calcul ZX vanille qui est de loin le calcul graphique le plus populaire pour le calcul quantique.


[^0]:    ${ }^{1}$ https://www.tesla.com/en_eu/supercharger

[^1]:    ${ }^{2}$ https:/ /pasqal.io/
    ${ }^{3}$ https://www.quandela.com/
    ${ }^{4}$ https://alice-bob.com/

[^2]:    ${ }^{5}$ In theory, Adiabatic evolution under general Hamiltonian is universal Aharonov et al., 2007

[^3]:    ${ }^{6}$ It can be solved with an algorithm in a time that is polynomial in the numeric values of its data $O\left(\operatorname{poly}\left(\sum_{j} w_{j}+\sum_{j} t_{j}\right)\right)$ but not on the length of their binary encodings.

[^4]:    ${ }^{7}$ Extreme points of the polyhedron $D p \leq d$ are integral

[^5]:    ${ }^{8}$ https://tikzit.github.io

